

FIR line emission from high-z galaxies

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1 AIM

By combining high resolution, radiative transfer cosmological simulations of a $z = 6$ galaxy with a sub-grid multi-phase model of its interstellar medium (ISM) we derive the expected intensity of several far infrared (FIR) emission lines.

2 COSMOLOGICAL SIMULATION

- Gadget-2 SPH simulation
- Number of particles: 2×512^3 baryonic + DM
- Simulated volume: $(10 \text{ h}^{-1} \text{ Mpc})^3$ comoving

We select a snapshot at $z=6.6$ and we identify the most massive halo: $M_h = 1.17 \times 10^{11} M_\odot$, $r_{\text{vir}} \approx 20 \text{ kpc}$

We post-processed UV radiative transfer LICORICE.

We assume that stars form in cells characterized by a gas density $\rho > 1 \text{ cm}^{-3}$. The properties are comparable to that of Himiko, the most luminous Lyman Alpha Emitter (LAE) at $z=6.6$.

3 ISM MULTIPHASE MODEL

We adopt a sub-grid scheme based on the model by [1] and [2], in which ISM thermal equilibrium is set by the balance between heating (cosmic rays, X-rays, and photo- electric effect on dust grains) and cooling (H, He, collisional excitation of metal lines, recombination on dust grains) processes [3]:

$$L(n, x_e, T) = n^2 \Lambda - n \Gamma = 0$$

where $n\Gamma$ ($n^2\Lambda$) is the heating (cooling) rate per unit volume and n is the total gas density. The ISM can be described as a two-phase gas in which the cold (CNM) and the warm neutral medium (WNM) are in pressure equilibrium.

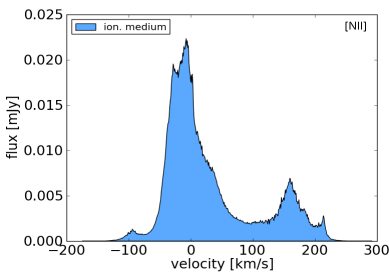
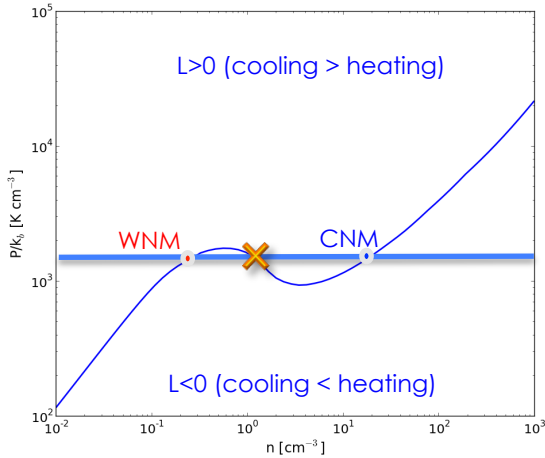


Fig. 3 - Left: Spectrum of [N II] from the ionized medium, binned in 1.0 km/s channels. Right: [N II] maps in mJy km/s with resolution of 0.1 arcsec and integrated over the entire spectral velocity range.

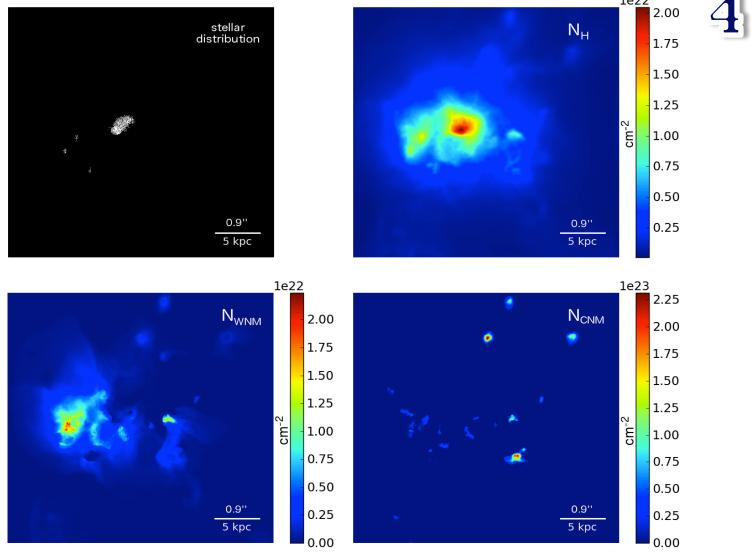


Fig. 1 - Upper panels: Projected stellar distribution (left) and hydrogen column density (right). Lower panels: warm (left), and cold (right) neutral medium column density. The distribution of WNM is more diffuse compared to that of CNM which is predominantly found in small ($D < 2 \text{ kpc}$) clumps far from SF regions.

For each cell we estimate the line luminosities $L_l = \epsilon V_{\text{cell}}$, where the emissivity, is given by:

$$\epsilon(n, T) = \Lambda^H \chi_i n^2 + \Lambda^e \chi_i x_e n^2$$

where n and T are the density and temperature of the WNM/CNM, Λ^H (Λ^e) is the specific cooling rate due to collision with H atoms (free electrons) taken from [3] and χ_i is the abundance of the i -th species.

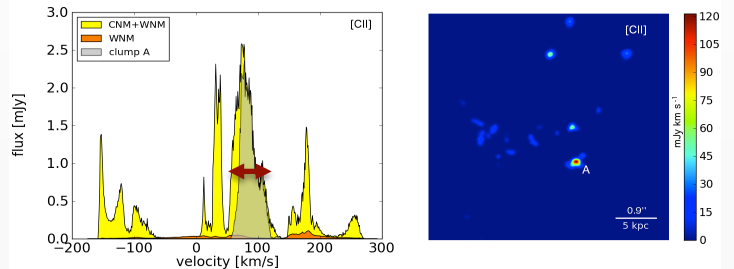


Fig. 2 - Left: Total (CNM+WNM) and WNM only (orange) spectrum of [C II] binned in 1.0 km/s channels. Right: [C II] maps in mJy km/s with resolution of 0.1 arcsec and integrated over the entire spectral velocity range. The contribution of clump A to the [C II] spectrum is plotted in gray. The FWHM (red arrow) of the line is $\sim 50 \text{ km/s}$ consistent with the marginal detection of [C II] in high- z LAEs [5].

6 We evaluate also the [N II] line luminosity which provides a complementary view of the ISM with respect to the [C II] line.

Indeed, [N II] traces only the ionized medium since its ionization potential (14.5 eV) exceeds 1 Ryd.

[N II] cooling rate due to collisions with free electrons is:

$$\epsilon_{\text{NII}}(n, T) = A h \nu (g_u/g_l) [(g_u/g_l + 1) (n_e/n_c)]^{-1} x_e$$

where $A = 7.5 \times 10^{-6} \text{ s}^{-1}$ is the Einstein coefficient, ν is the frequency for the $3^3P_2 \rightarrow 3^3P_1$ transition, h is the Planck constant, g_u/g_l is the ratio of the statistical weights in the upper and lower levels, and $n_c = 300 \text{ cm}^{-3}$ is the [N II] critical density for $T = 10^4 \text{ K}$.

The [C II] emission line is detectable with the ALMA full array in $1.9 < t < 7.7 \text{ hr}$ in star forming, high- z galaxies with $Z_\odot > Z > 0.5 Z_\odot$. We note that the predicted fluxes are sensitive to the actual value of Z , implying that a [C II] line detection can strongly constrain the LAE metallicity.