

Sternentstehung - Star Formation

Winter term 2020/2021

Henrik Beuther, Thomas Henning & Sümeyye Suri

<i>03.11 Today: Introduction & Overview</i>	<i>(Beuther)</i>
<i>10.11 Physical processes I</i>	<i>(Beuther)</i>
<i>17.11 Physical processes II</i>	<i>(Beuther)</i>
<i>24.11 Molecular clouds as birth places of stars</i>	<i>(Suri)</i>
<i>01.12 Molecular clouds (cont.), Jeans Analysis</i>	<i>(Suri)</i>
<i>08.12 Collapse models I</i>	<i>(Henning)</i>
<i>15.12 Collapse models II</i>	<i>(Henning)</i>
----- Christmas break -----	
<i>12.01 Protostellar evolution</i>	<i>(Beuther)</i>
19.01 Pre-main sequence evolution & outflows/jets	(Beuther)
26.01 Accretion disks I	(Henning)
02.02 Accretion disks II	(Henning)
09.02 High-mass star formation, clusters and the IMF	(Suri)
16.02 Extragalactic star formation	(Henning)
23.02 Examination week, no star formation lecture	

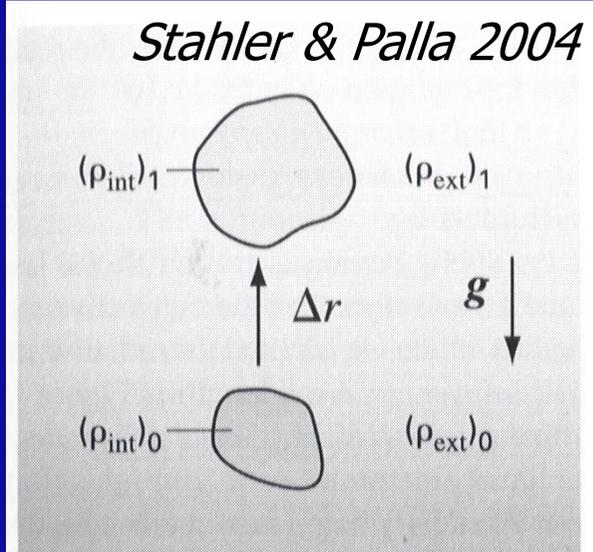
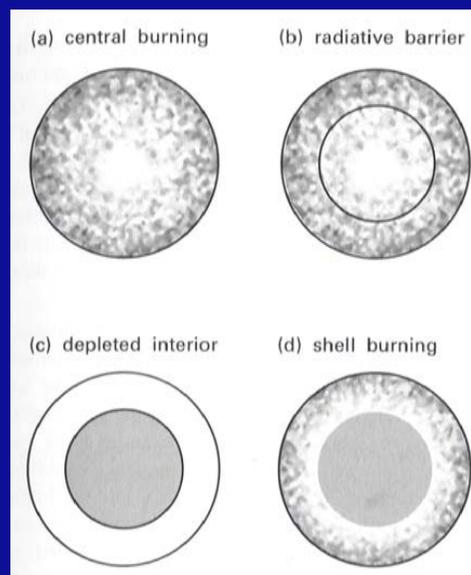
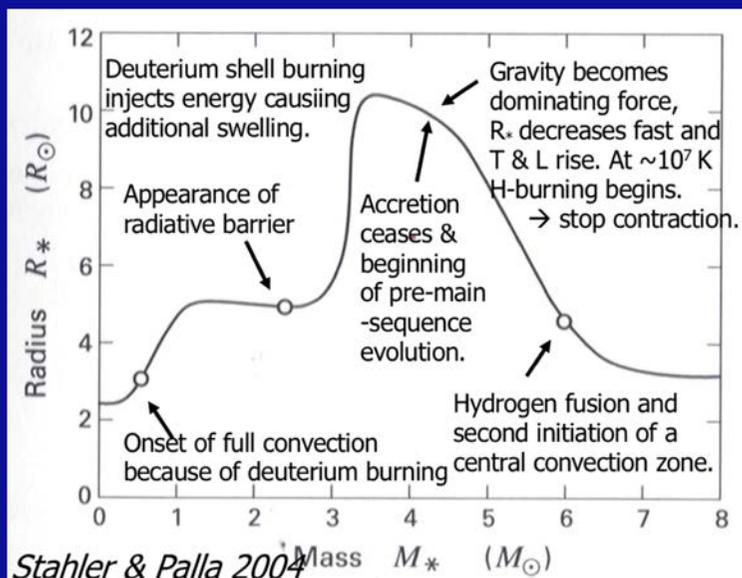
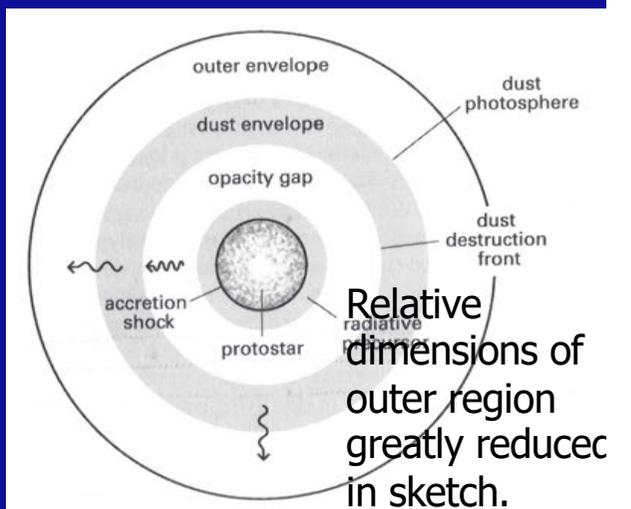
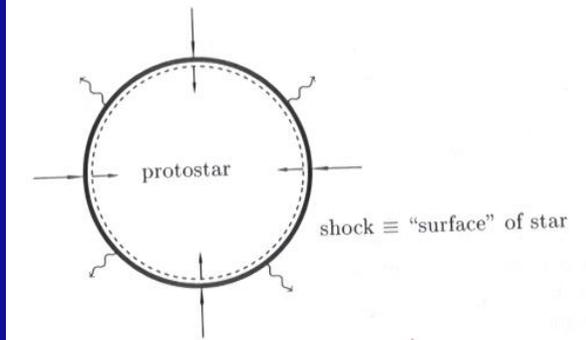
Book: Stahler & Palla: The Formation of Stars, Wileys

More Information and the current lecture files: http://www.mpia.de/homes/beuther/lecture_ws2021.html

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Summary last week

- Protostellar evolution, 1st and 2nd core, accretion luminosity
- Envelope structure
- Convection, entropy profile of protostar
- Structure of protostar
- Definition: protostar vs. pre-main sequence star



Topics today

- Pre-main-sequence evolution



- General outflow properties

- Jet launching

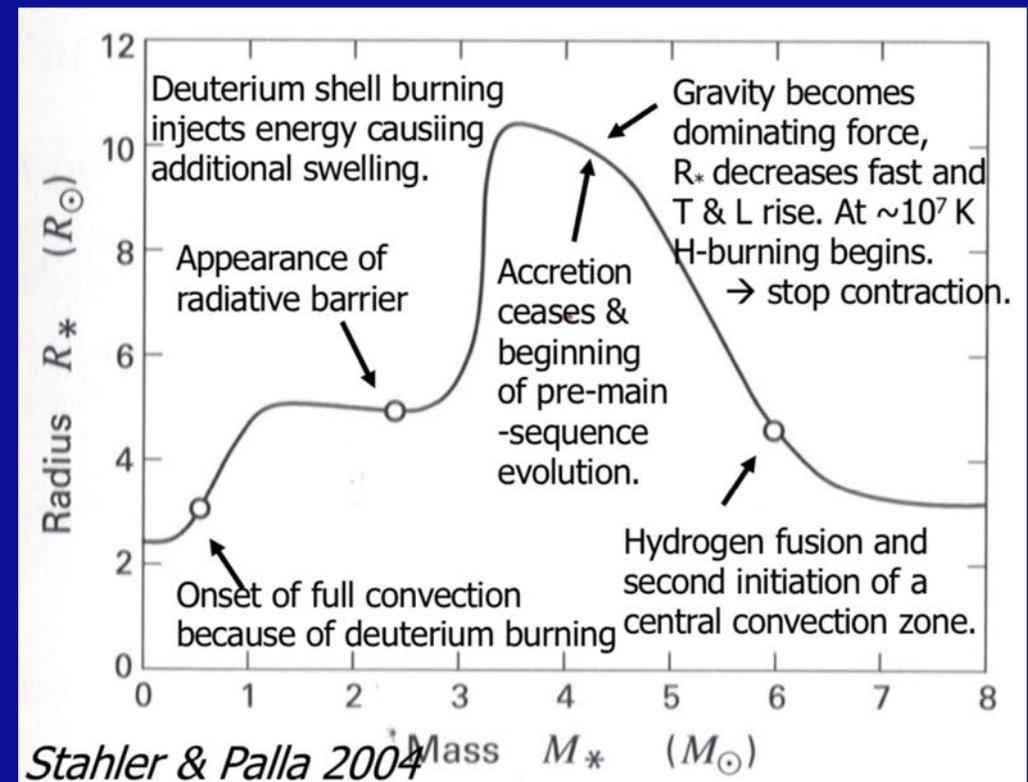
- Outflow driving and entrainment

Protostellar vs. pre-main sequence evolution (mainly for low-mass protostars)

- Accretion ceases → protostar contracts → gain energy by gravity

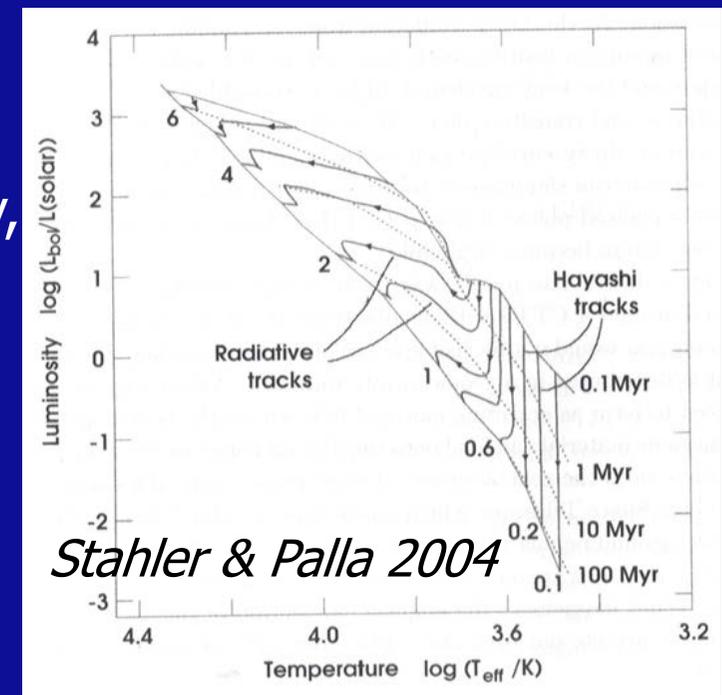
→ Main luminosity not accretion but due to gravitational contraction.

→ Identify this point with end of protostellar and beginning of pre-main sequence phase (low-mass stellar evolution)



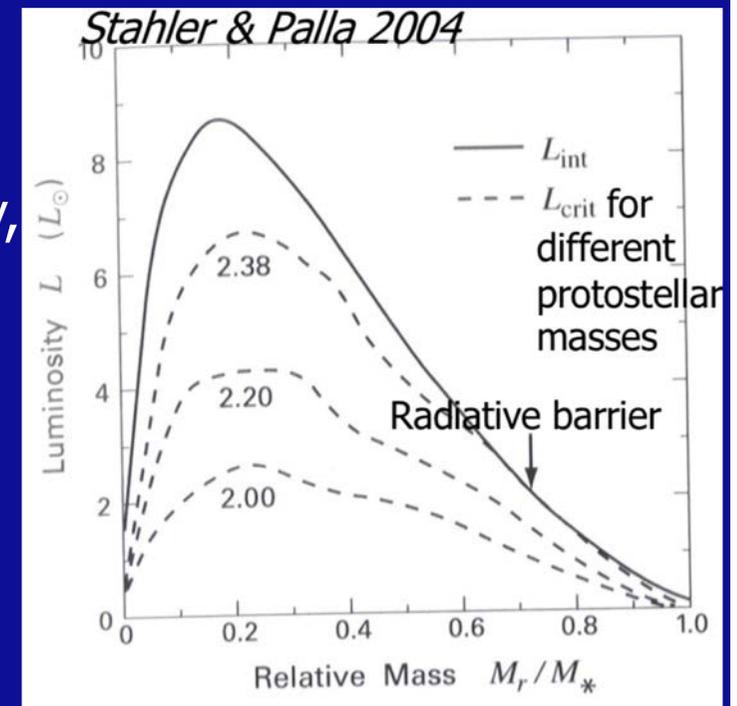
Further contraction until H-burning I

- Different evolution for low- and high-mass pre-main sequence stars:
 - Low-mass:
 - still convective, still remaining deuterium burning
 - Hayashi tracks: Shrinking releases grav. energy & T_{surface} approx. constant
 - $L = 4\pi R^2 \sigma_B T_{\text{eff}}^4 \propto R_*^2 \rightarrow L$ decreases $\rightarrow L$ falls below L_{crit} .
 \rightarrow Radiative core forms again with a shrinking outer convective layer.
 - During further slow contraction internal energy, temperature & luminosity rise again until hydrogen burning starts \rightarrow ZAMS.
 - Stars below $\sim 0.4 M_{\text{sun}}$ reach the ZAMS still fully convective.

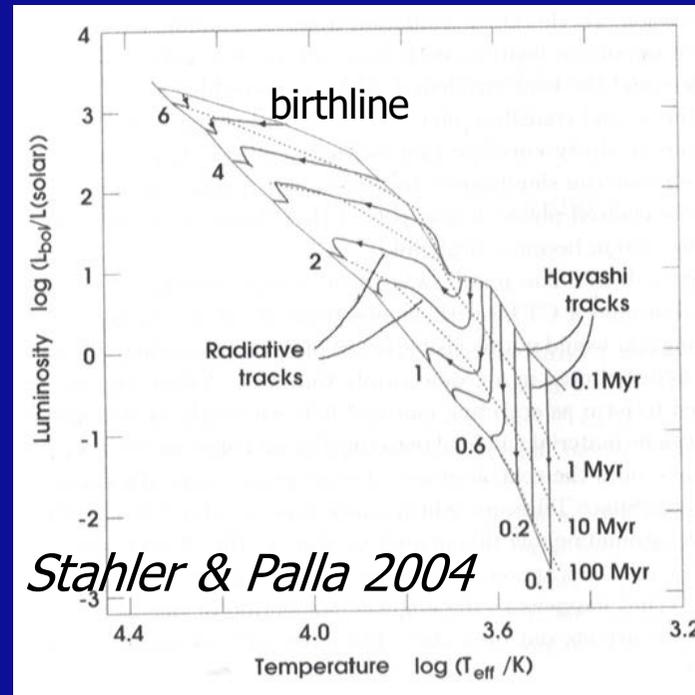


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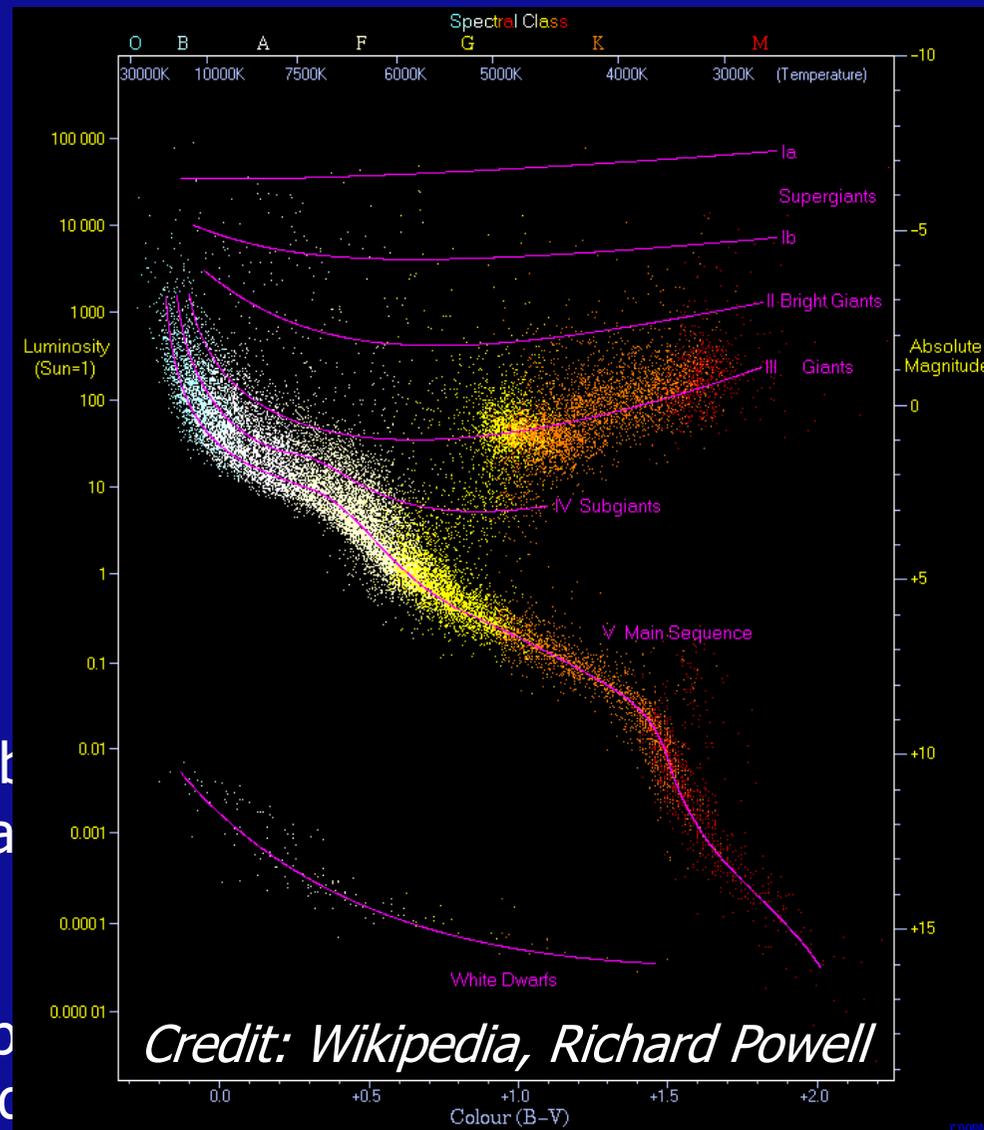


Hertzsprung Russel (HR) diagram I



- Birthline first observationally \rightarrow locus where stars first appear in the HR diagram emanating from their dusty natal envelope.
- Theoretically: birthline the time where the main accretion has stopped \rightarrow pre-main sequence star gains the main luminosity from grav. contraction.
 - \rightarrow Approx. coincides with quasi-static contraction in still convective phase

Hertzsprung Russel (HR) diagram I



- Birthline first of
diagram emanates

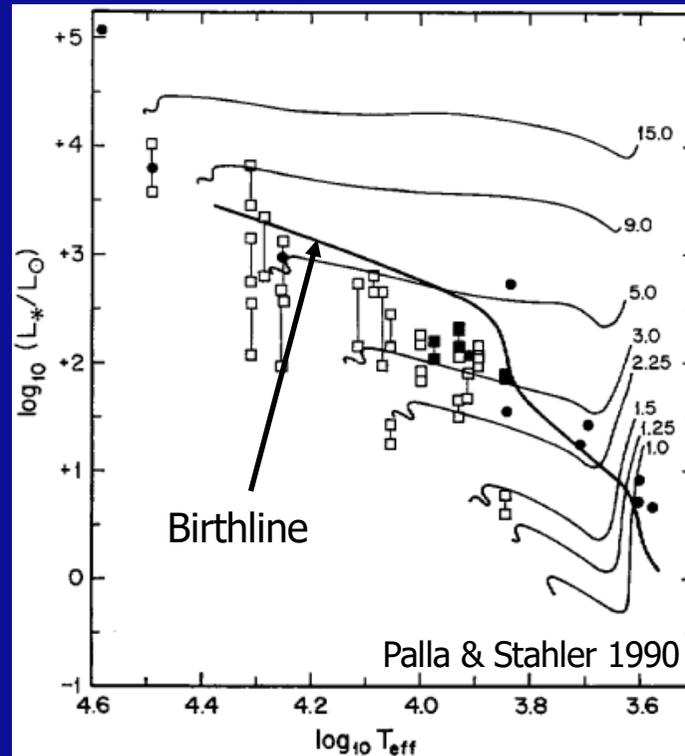
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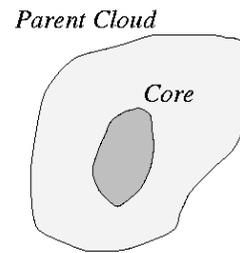
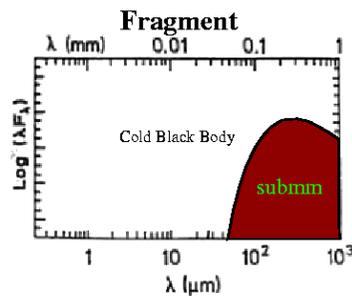
Hertzsprung Russel (HR) diagram II



- Intermediate-mass protostars fully radiative when stopping accretion
→ no vertical Hayashi part but direct horizontal radiative tracks
- High-mass stars: short Kelvin-Helmholtz contraction time-scale
→ start nuclear H-burning - entering the ZAMS – before accretion ends
- no (visible) pre-main sequence evolution since H-burning starts still deeply embedded in their natal cores

Observable spectral energy distributions (SEDs)

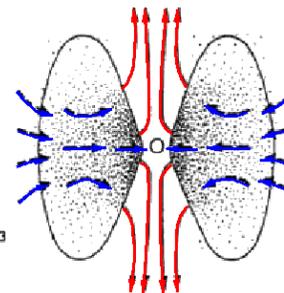
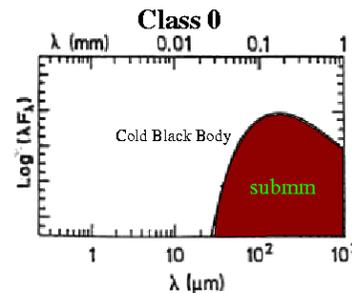
Pre-Stellar Phase



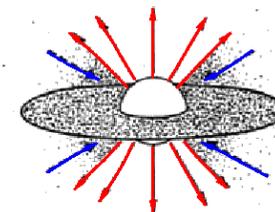
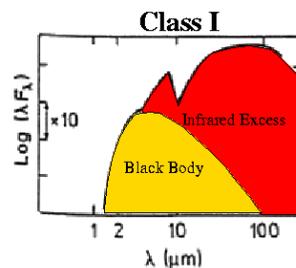
Pre-Stellar Dense Core
 $T_{bol} \sim 10-20 \text{ K}, M_* = 0$
 - 1 000 000 yr

Formation of the central protostellar object

Protostellar Phase



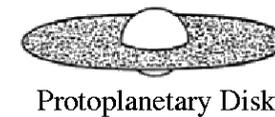
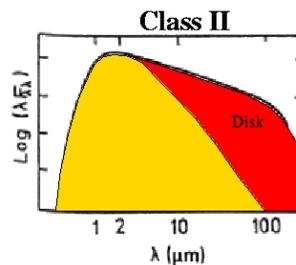
$t \sim 0 \text{ yr}$
 Young Accreting Protostar
 $T_{bol} < 70 \text{ K}, M_* \ll M_{env}$
 < 30 000 yr



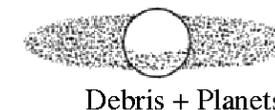
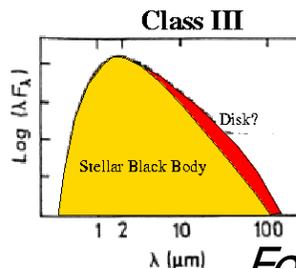
Evolved Accreting Protostar
 $T_{bol} \sim 70-650 \text{ K}, M_* > M_{env}$
 ~ 200 000 yr

Birthline for Pre-main sequence stars

Pre-Main Sequence Phase



Classical T Tauri Star
 $T_{bol} \sim 650-2880 \text{ K}, M_{Disk} \sim 0.01 M_{\odot}$
 ~ 1 000 000 yr



Weak T Tauri Star
 $T_{bol} > 2880 \text{ K}, M_{Disk} < M_{Jupiter}$
 ~ 10 000 000 yr

Followina P. Andre ↓ Time

Topics today

- Pre-main-sequence evolution



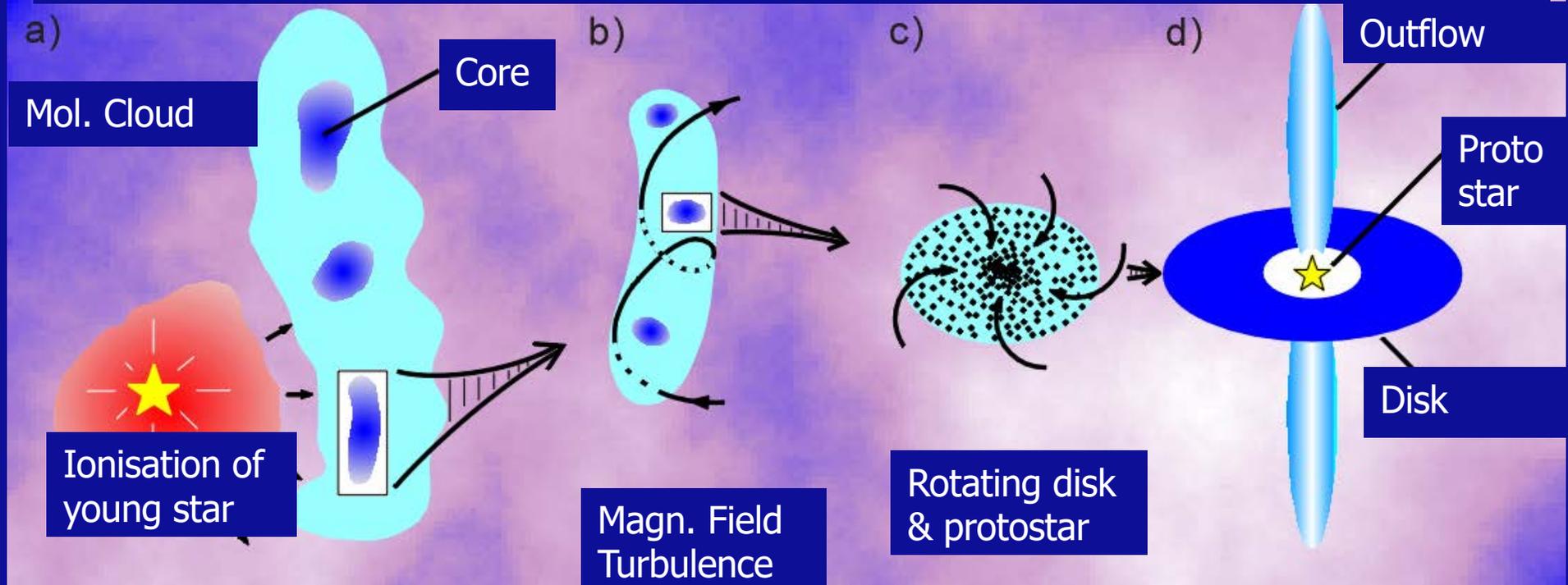
- General outflow properties

- Jet launching

- Outflow driving and entrainment

Star formation paradigm

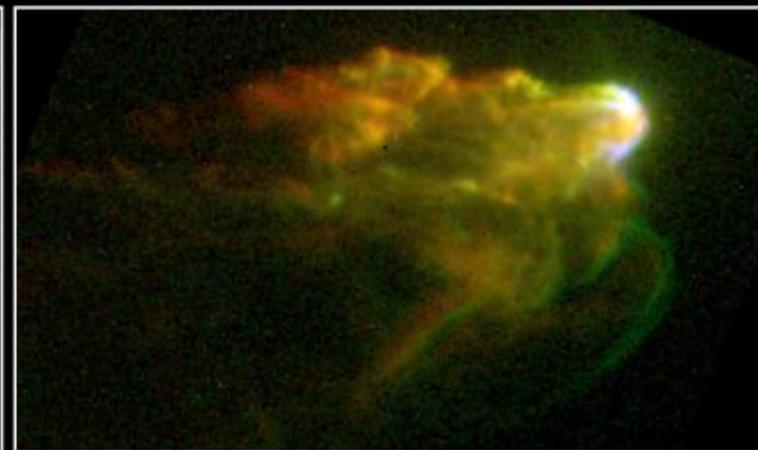
Phases of star formation



<https://www.mpifr-bonn.mpg.de/473576/starform>

Discovery of outflows I

Herbig 1950, 1951; Haro 1952, 1953



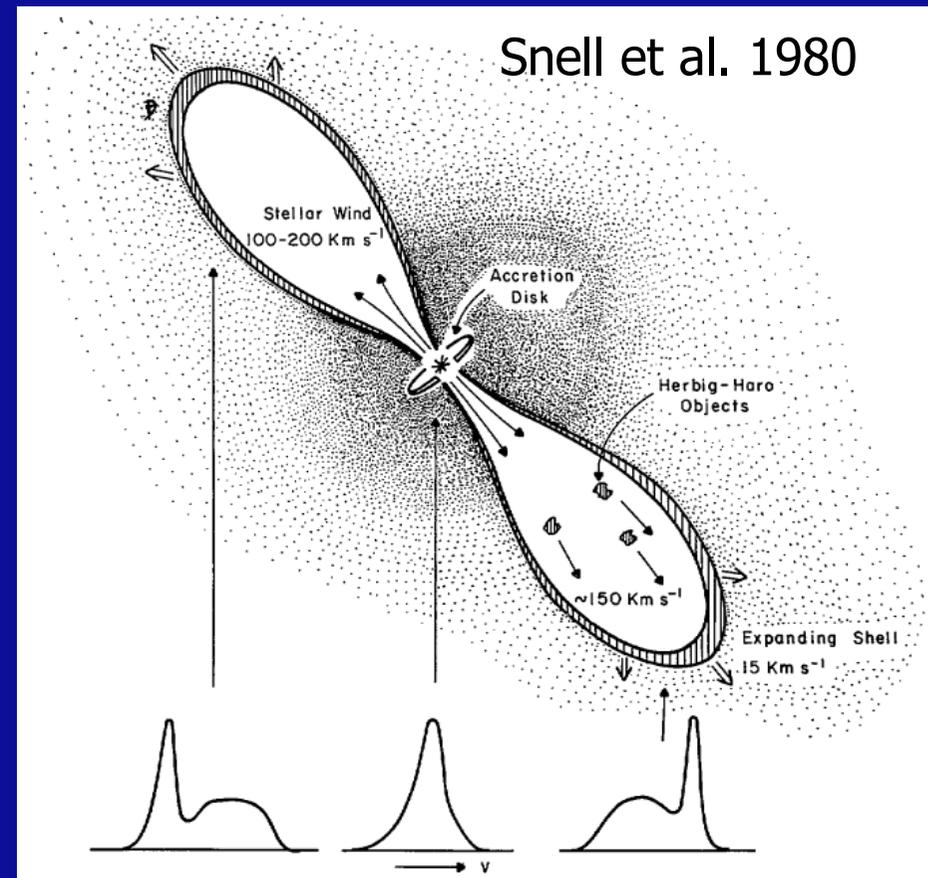
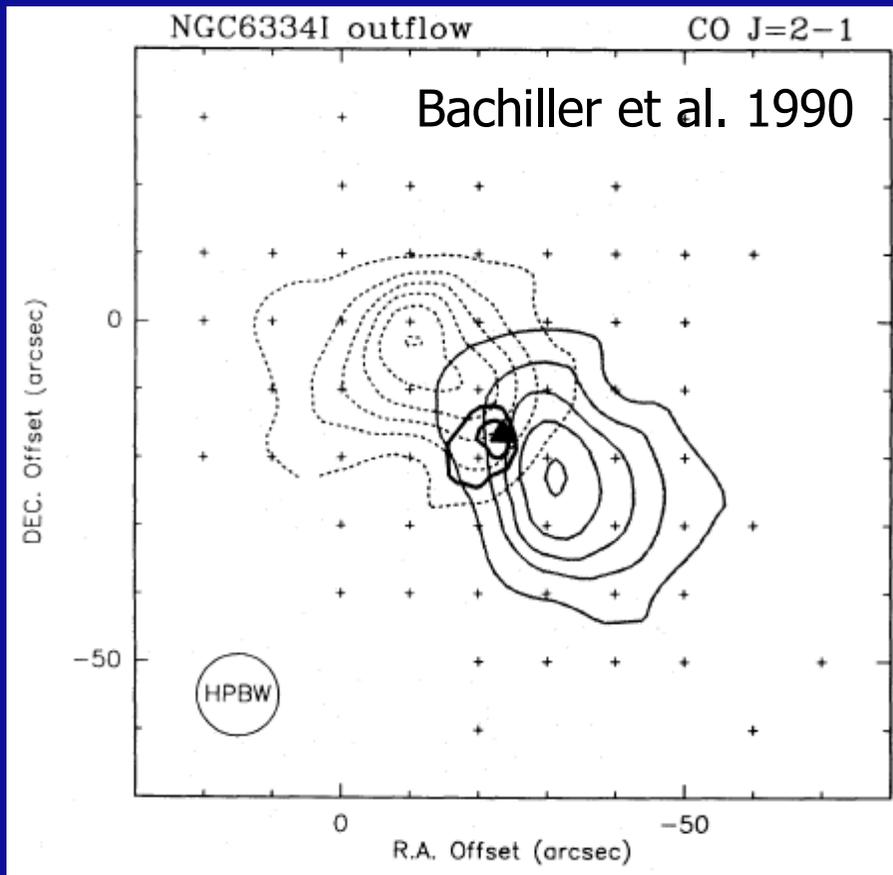
Jets from Young Stars · HH1/HH2

HST · WFPC2

PRC95-24c · ST ScI OPO · June 6, 1995 · J. Hester (AZ State U.), NASA

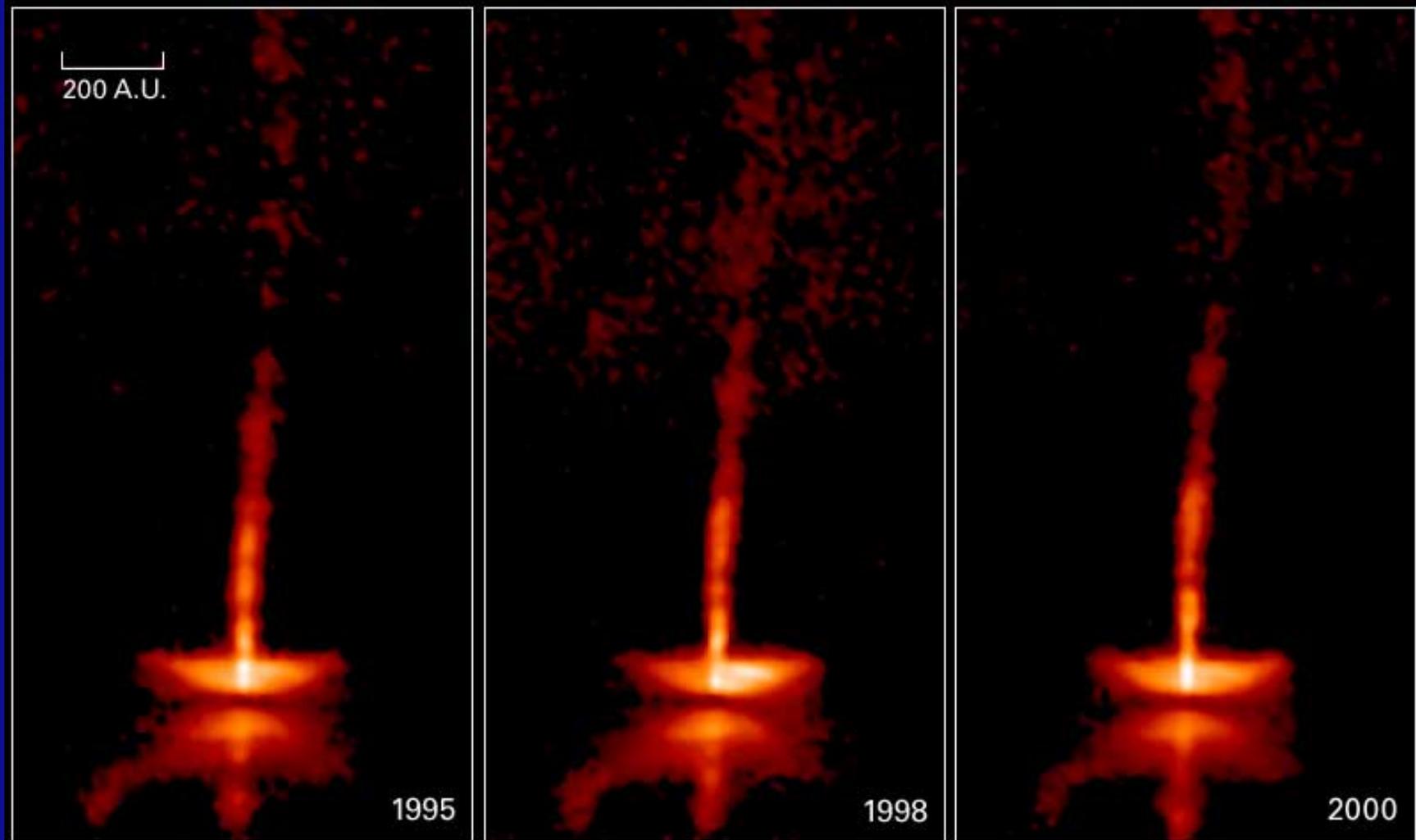
Initially thought to be embedded protostars but soon spectra were recognized as caused by shock waves → jets and outflows indicated

Discovery of outflows II



- In the mid to late 70th, first CO non-Gaussian line wing emission detected (Kwan & Scoville 1976).
- Bipolar structures, extremely energetic, often associated with HH objects

HH30, a disk-outflow system



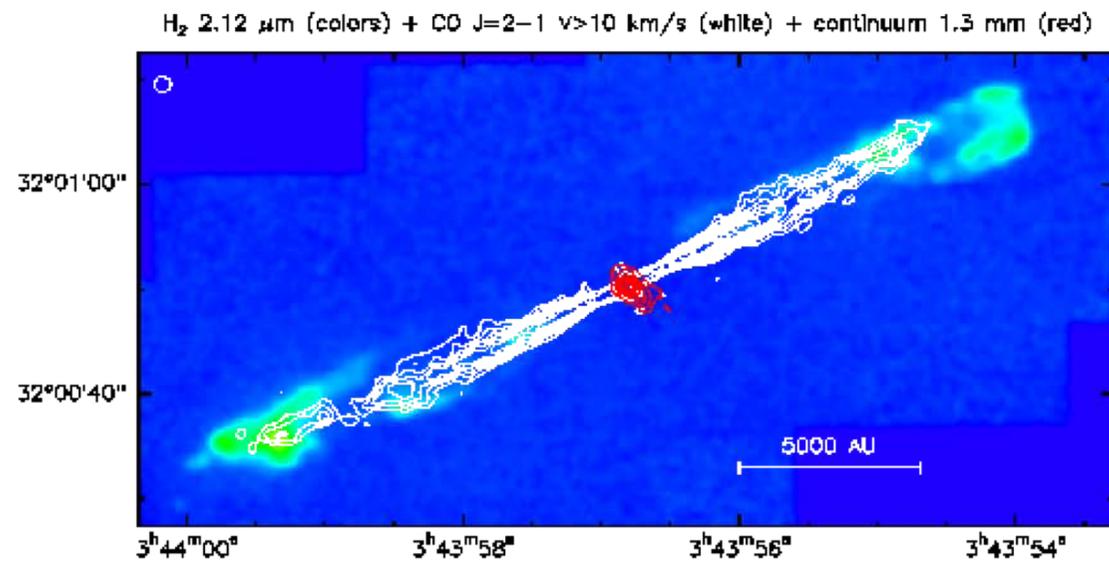
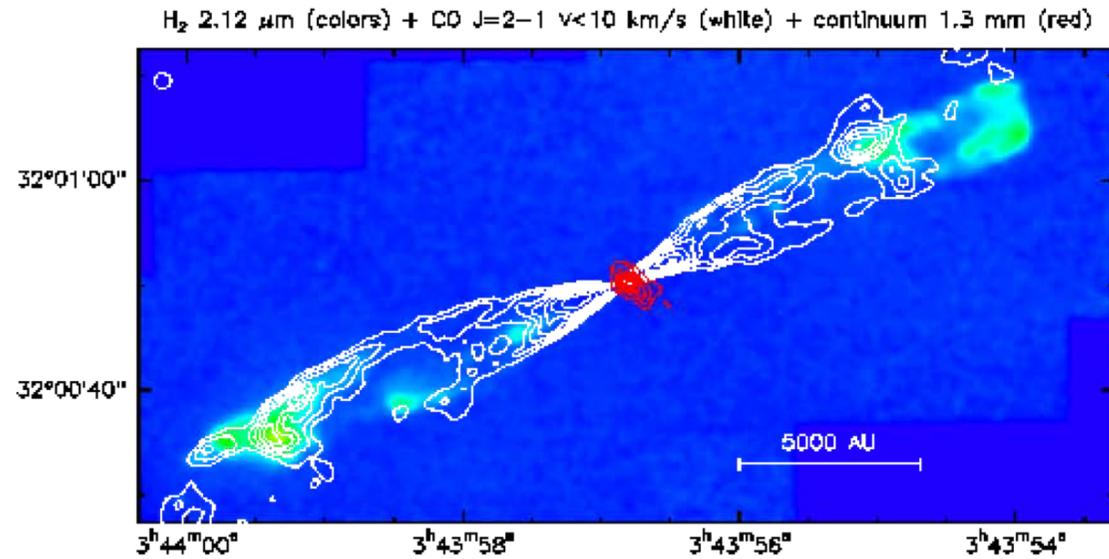
The Dynamic HH 30 Disk and Jet

HST • WFPC2

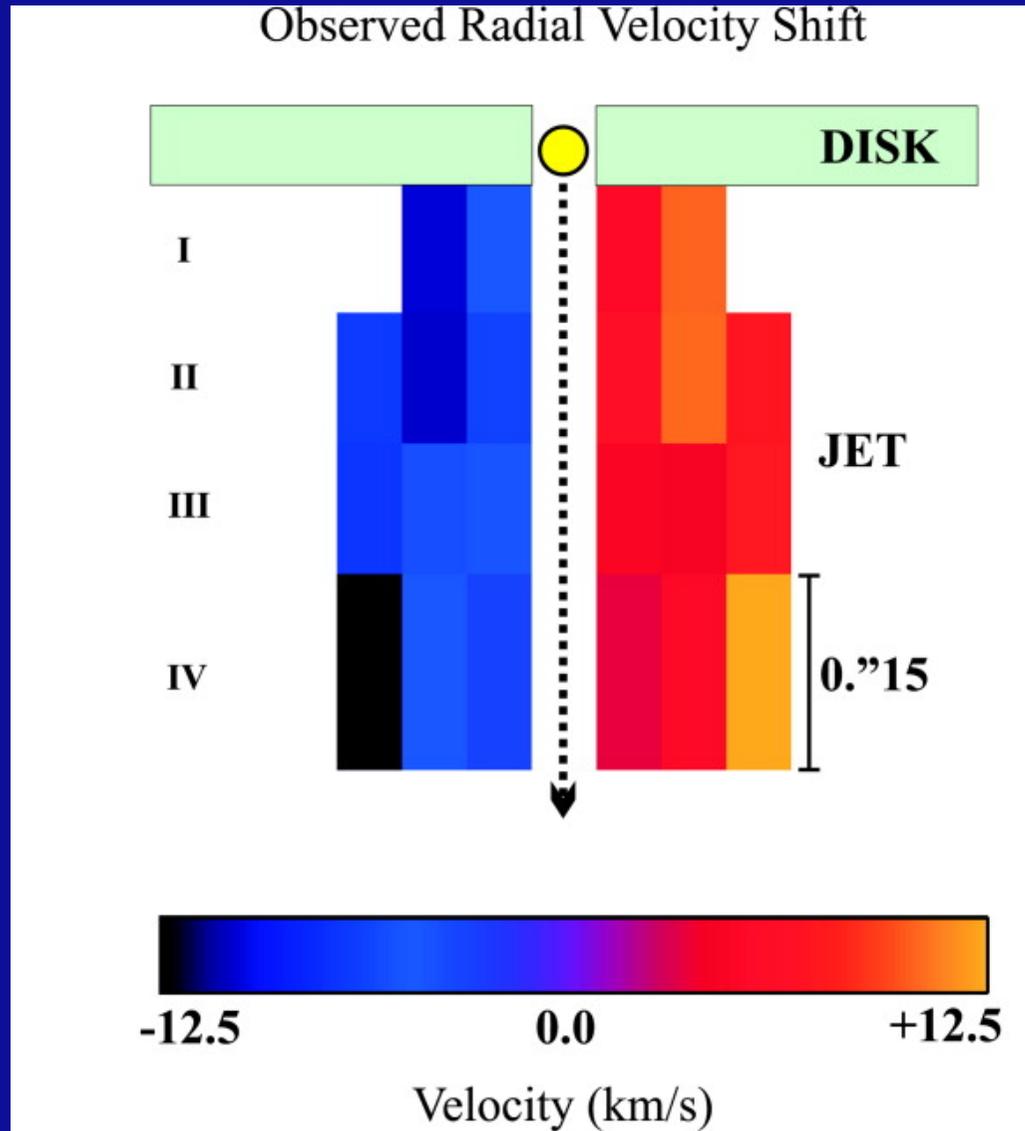
NASA and A. Watson (Instituto de Astronomía, UNAM, Mexico) • STScI-PRC00-32b

The prototypical molecular outflow HH211

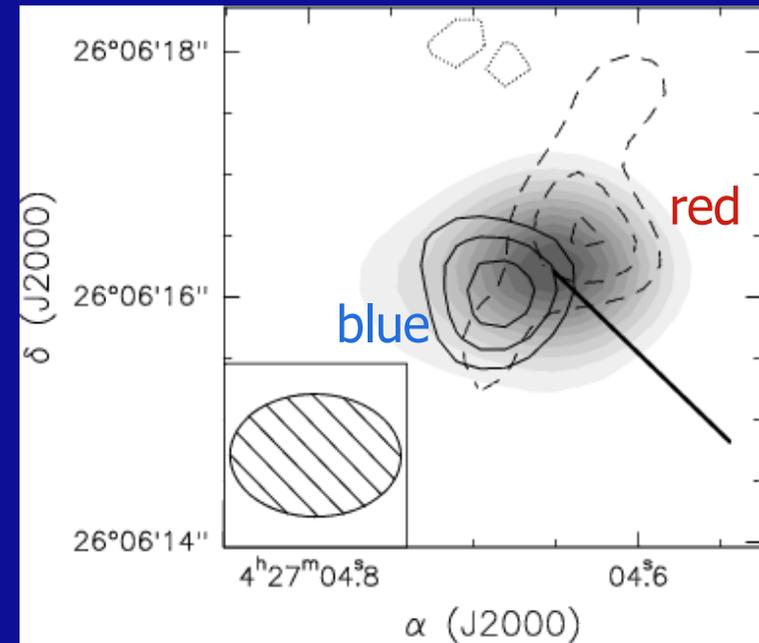
HH211, Gueth et al. 1999



Jet rotation in DG Tau



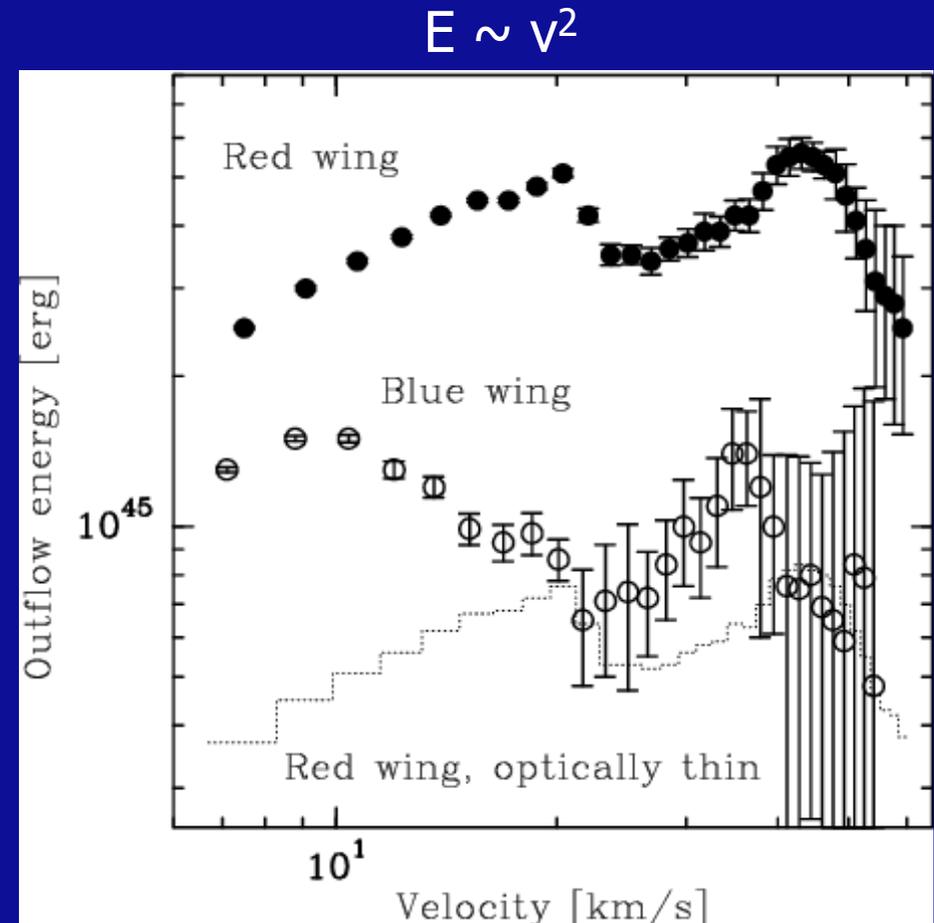
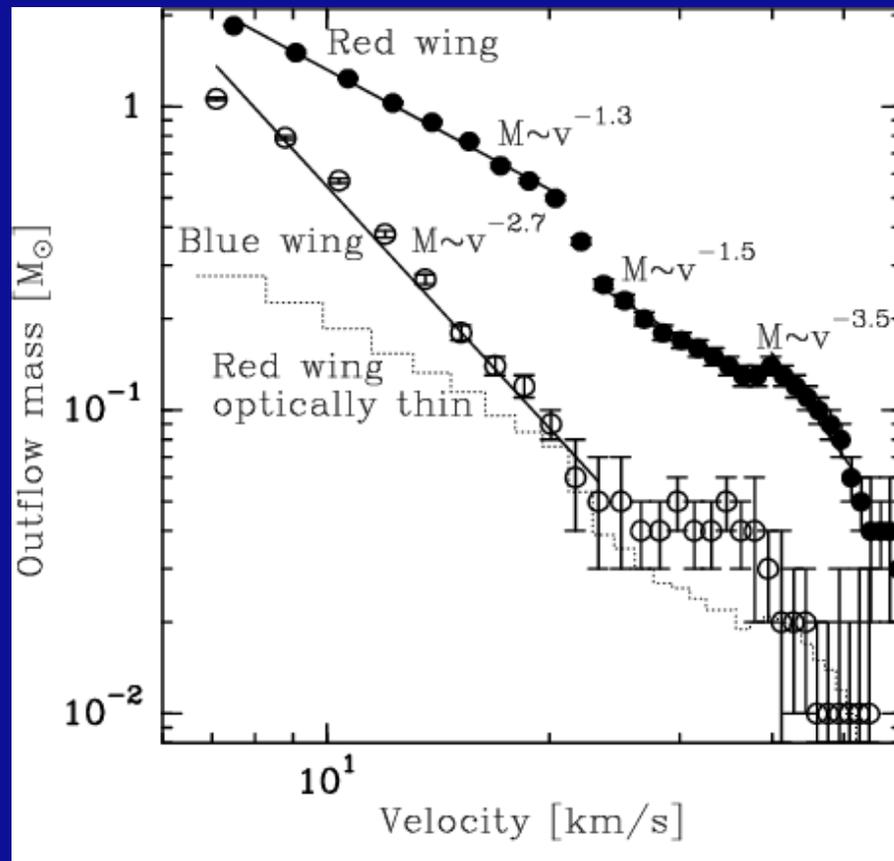
Bacciotti et al. 2002



Testi et al. 2002

Corotation of disk and jet

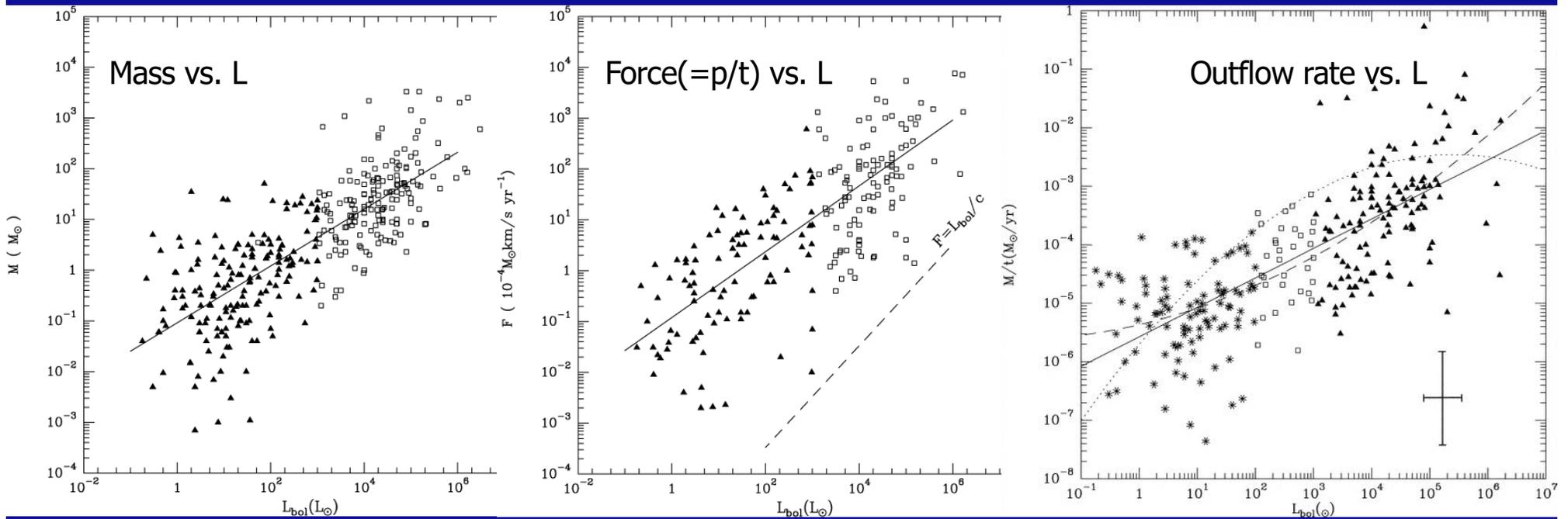
Mass vs. velocity, energy vs. velocity



- Mass-velocity relation exhibits broken power-law, steeper further out.
- Energy at high velocities of the same magnitude than at low velocities.

General outflow properties

- Jet velocities 100-500 km/s \Leftrightarrow Outflow velocities 10-50 km/s
- Estimated dynamical ages between 10^3 and 10^5 years
- Size between 0.1 and 1 pc
- Force provided by stellar radiation too low (middle panel)
→ non-radiative processes necessary!



Wu et al. 2004, 2005

Impact on surrounding cloud

- Entrain large amounts of cloud mass with high energies.
- Partly responsible to maintain turbulence in cloud.
- Can disrupt the cores to stop any further accretion.
- May trigger collapse in neighboring cores.
- Via shock interactions heat the cloud.
- Alter the chemical properties.

Topics today

- Pre-main-sequence evolution



- General outflow properties

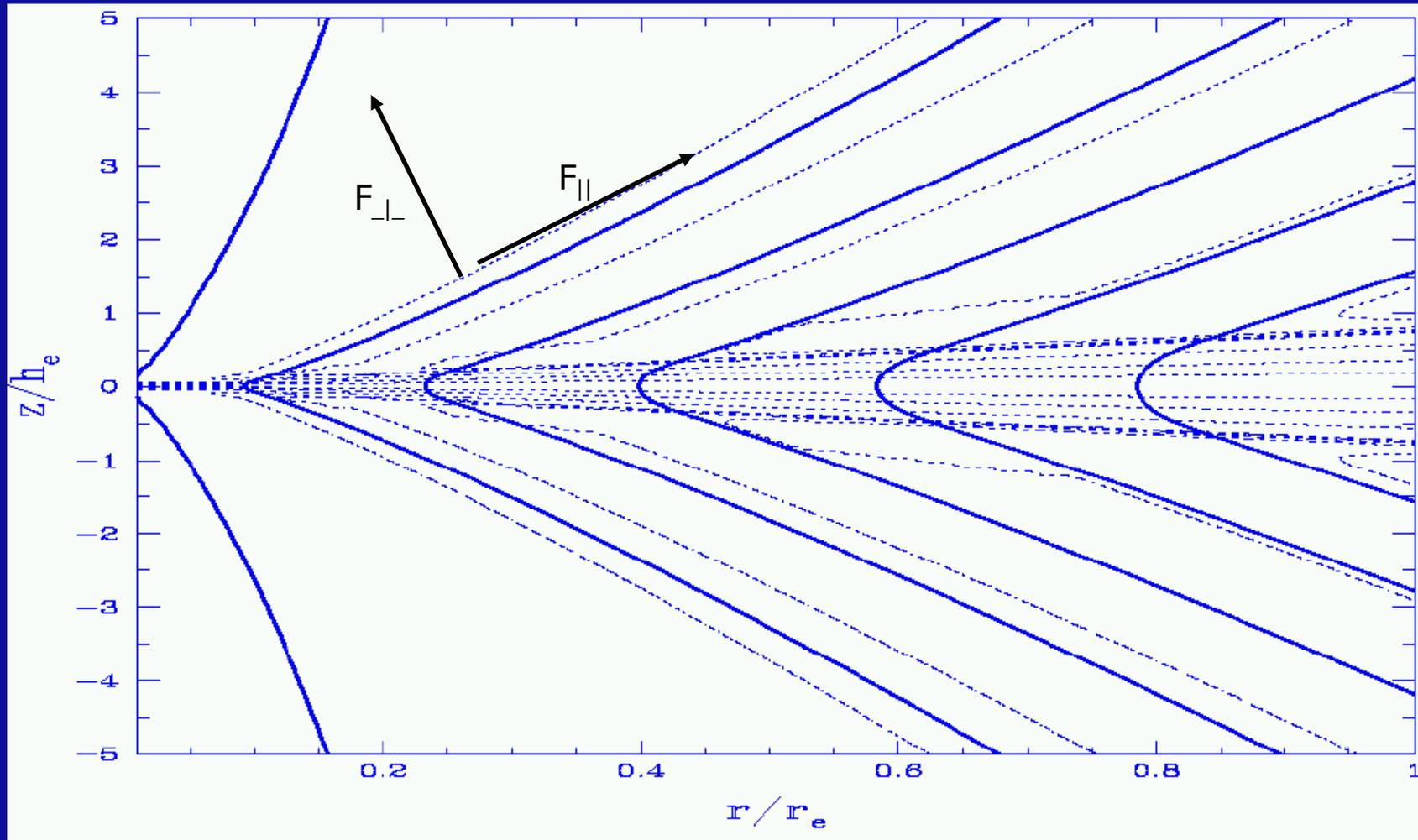
- **Jet launching**

- Outflow driving and entrainment

Jet launching from accretion disks

“magnetic accretion-ejection structures” (Ferreira et al 1995-1997):

- 1) disk material diffuses across magnetic field lines,
- 2) is lifted upwards by MHD forces, then
- 3) couples to the field and 4) becomes accelerated magnetocentrifugally and 5) collimated



Magnetic field lines (thick)
and streamlines (dashed)

Jet launching

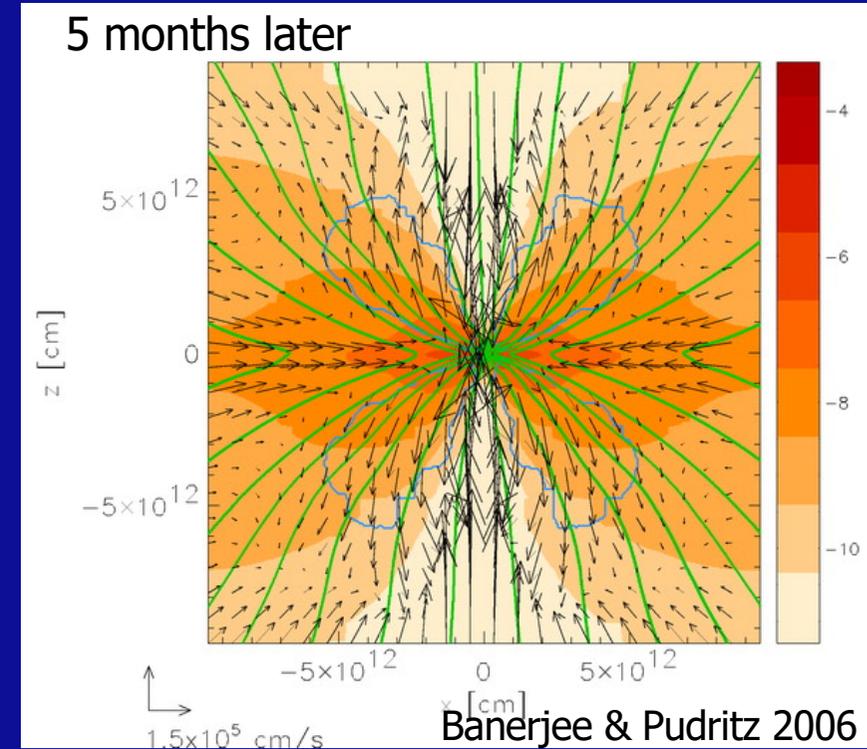
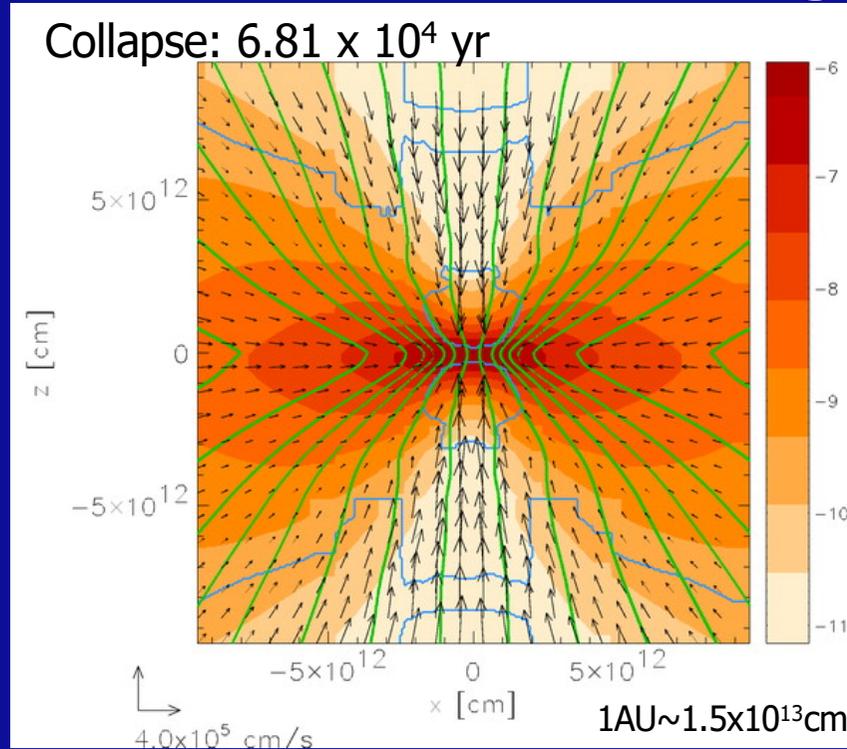
- Consensus: Jets are driven by magnetocentrifugal winds from magnetic field lines anchored in rotating circumstellar disks.

Disk winds $\leftarrow \rightarrow$ X-winds

Launching over larger
disk area?

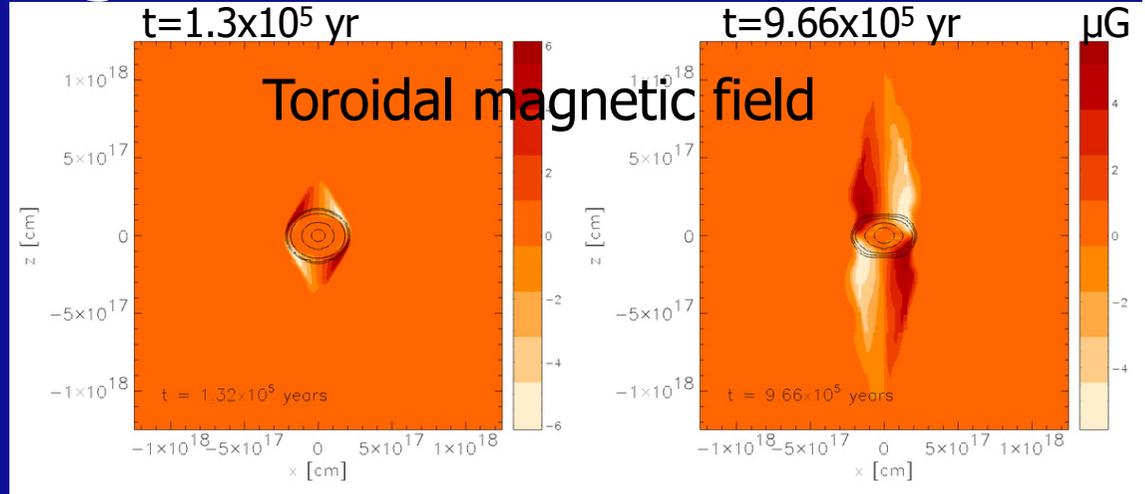
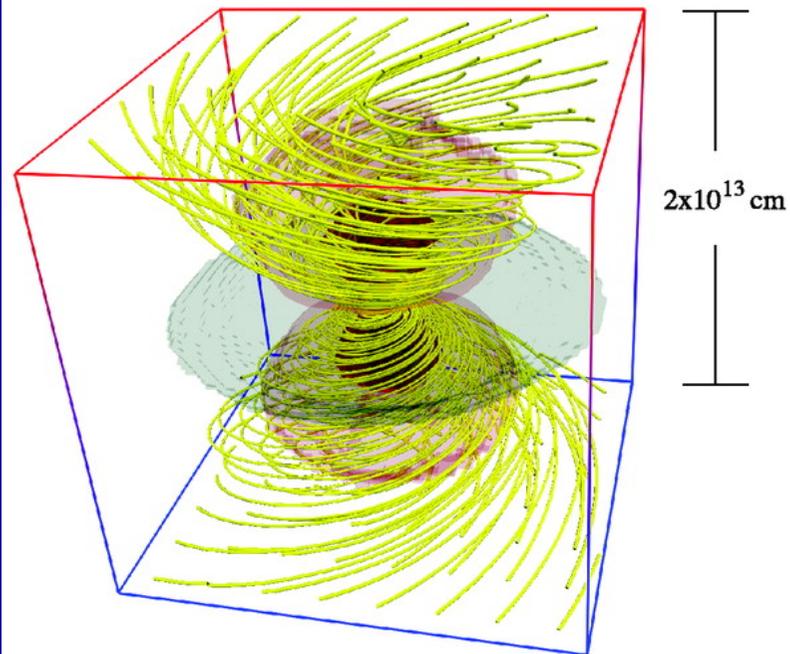
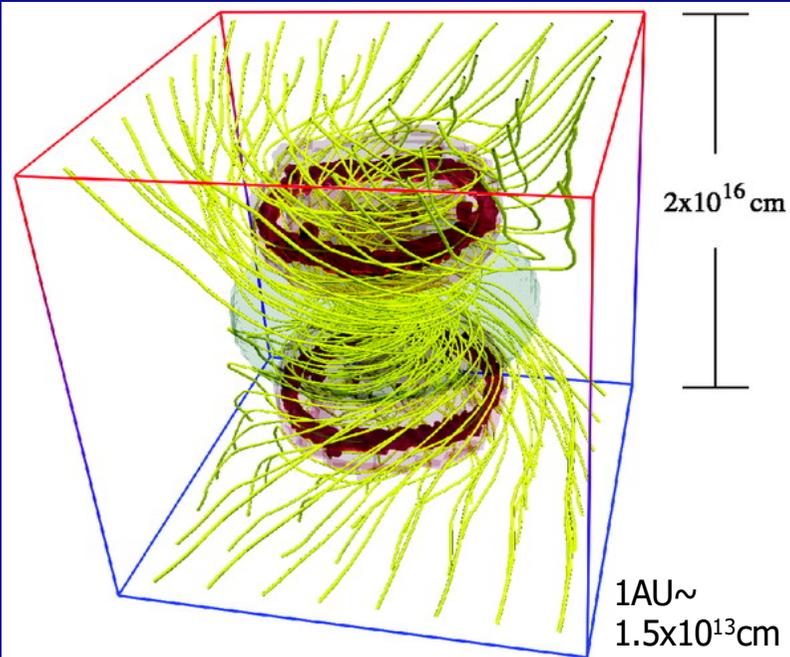
$\leftarrow \rightarrow$ Launching from a small area
close to disk truncation?

Jet-launching: Disk winds I



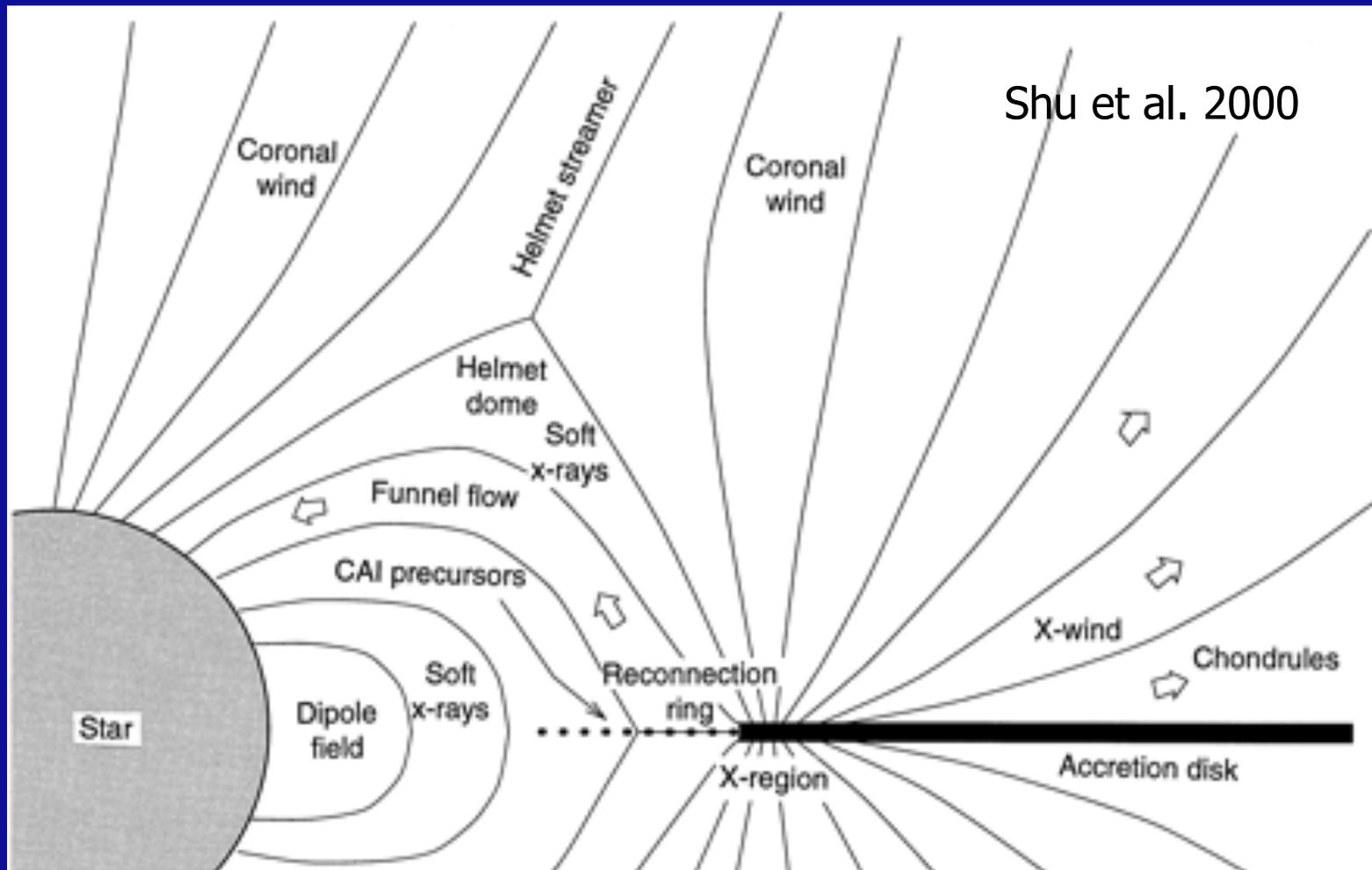
- Infalling core pinches magnetic field.
- If poloidal magnetic field component B_p has angle larger 30° from vertical
→ centrifugal forces launch matter-loaded wind along field from disk
- Wind transports away from 60 to 100% of disk angular momentum.

Jet-launching: Disk winds II



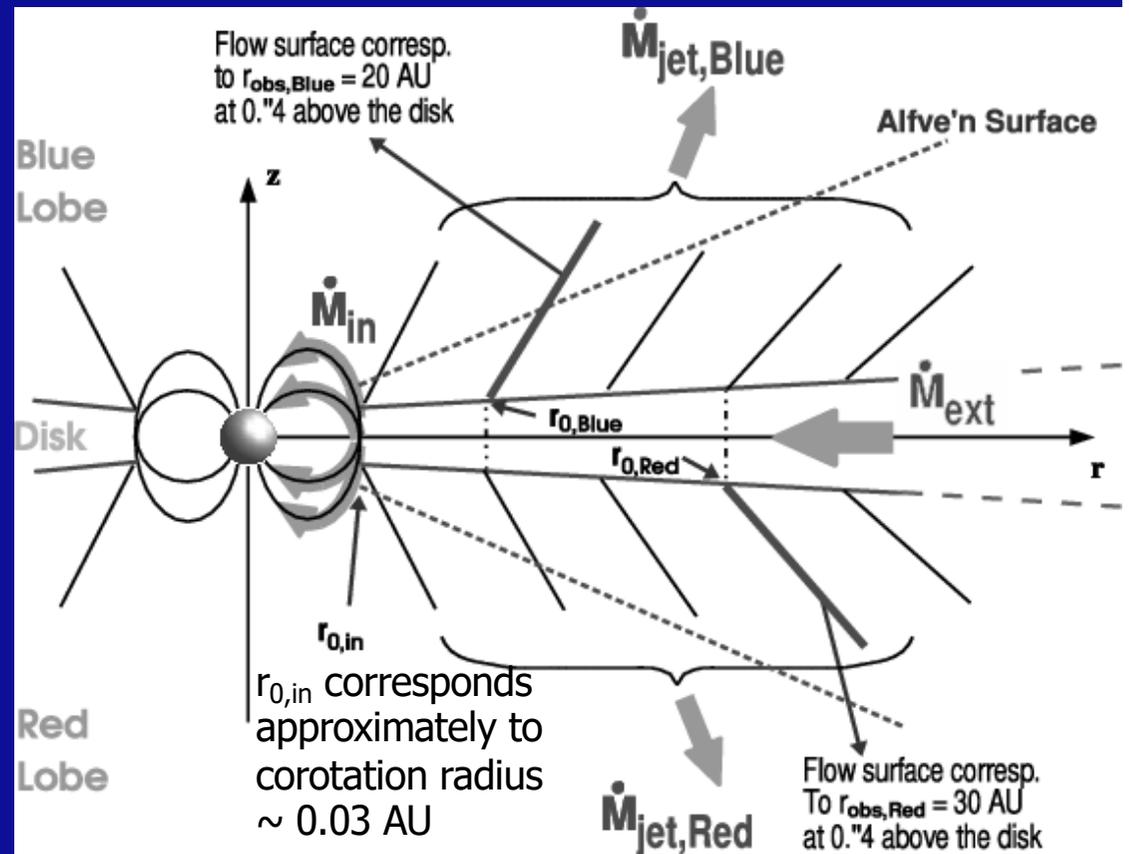
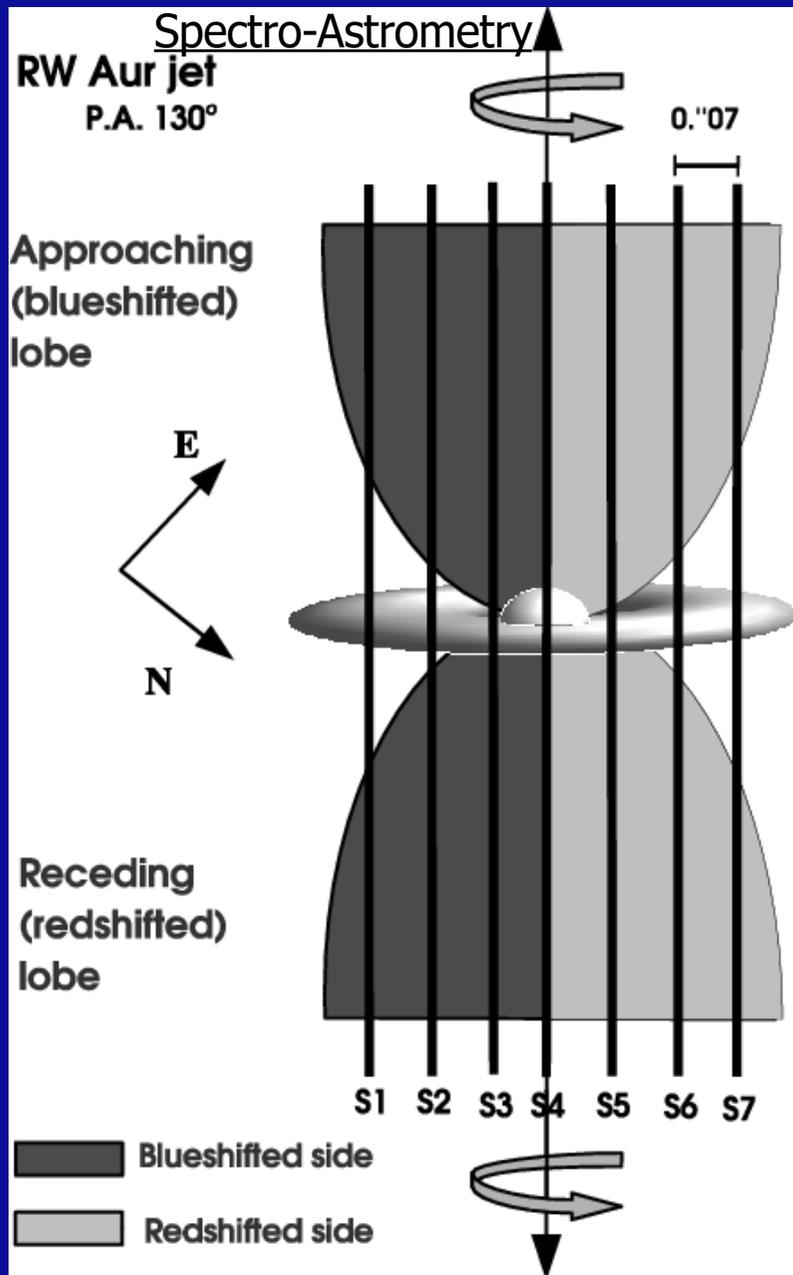
- On larger scales, a strong toroidal magnetic field B_ϕ builds up during collapse.
- At large radii (outside Alfvén radius r_A , the radius where kin. energy equals magn. energy) B_ϕ/B_p much larger than 1
 \rightarrow collimation via Lorentz-force $F_L \sim j_z B_\phi$

X-winds



- The wind is launched magneto-centrifugally from the inner co-rotation radius of the accretion disk ($\sim 0.03\text{AU}$)

Jet-launching points and angular momenta



- From toroidal and poloidal velocities
 → footpoints r_0 , where gas comes from
 → outer r_0 for the blue and red wing are about 0.4 and 1.6 AU (lower limits)
 → consistent with disk winds
- About 2/3 of the disk angular momentum may be carried away by jet.

Woitas et al. 2005

Topics today

- Pre-main-sequence evolution



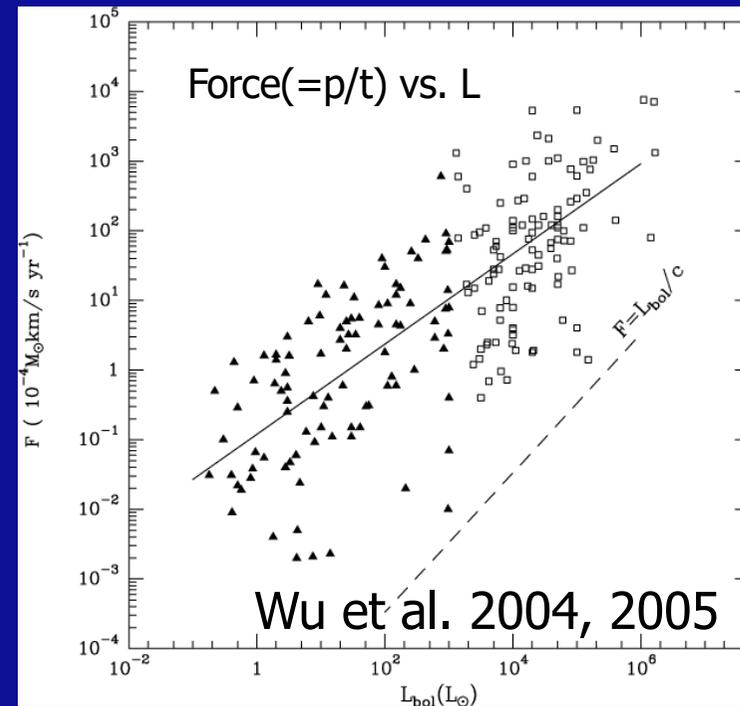
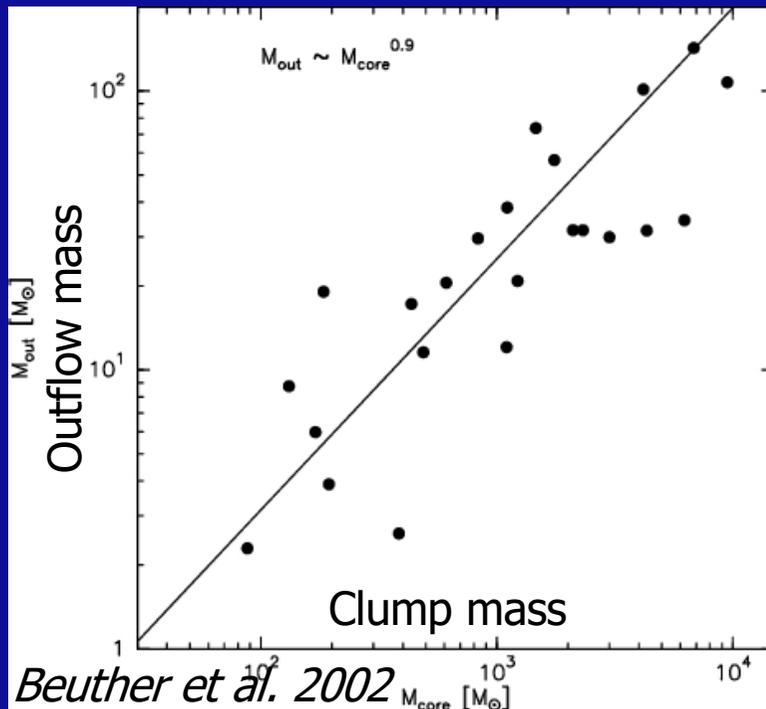
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- Outflow driving and entrainment

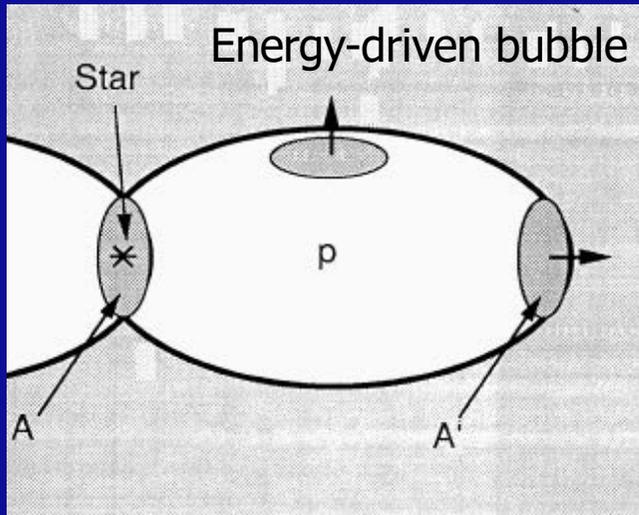
Outflow driving I

- Molecular outflow masses usually much larger than stellar masses
→ outflow-mass not directly from star-disk but swept-up entrained gas.
- Force observed in outflow cannot be explained just by force exerted from central object → other outflow driving and entrainment processes required.



Outflow driving II

Momentum-driven vs. energy-driven molecular outflows



- Energy-driven: jet-energy conserved in pressurized bubble that gets released adiabatically as the bubble expands.
- large transverse velocities which are not observed
- momentum conservation better!

Completely radiative shock → only dense plug at front

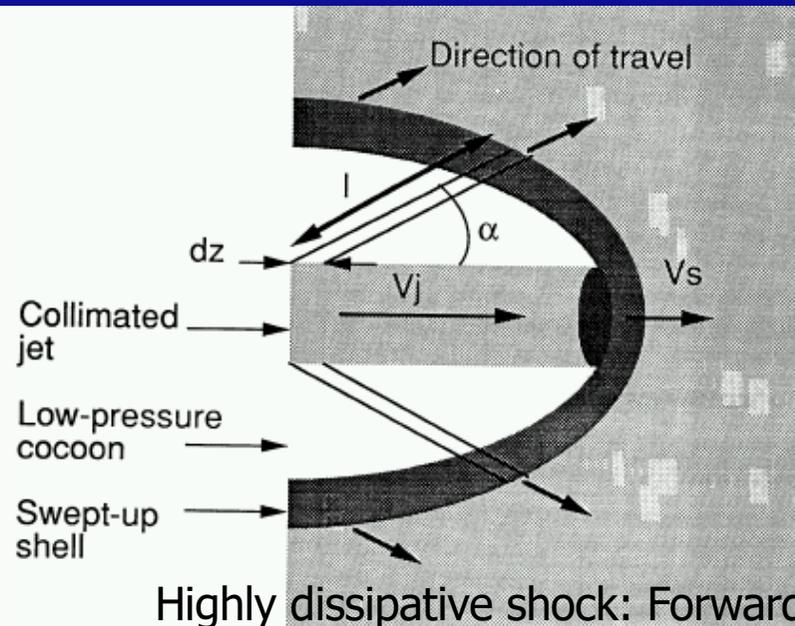
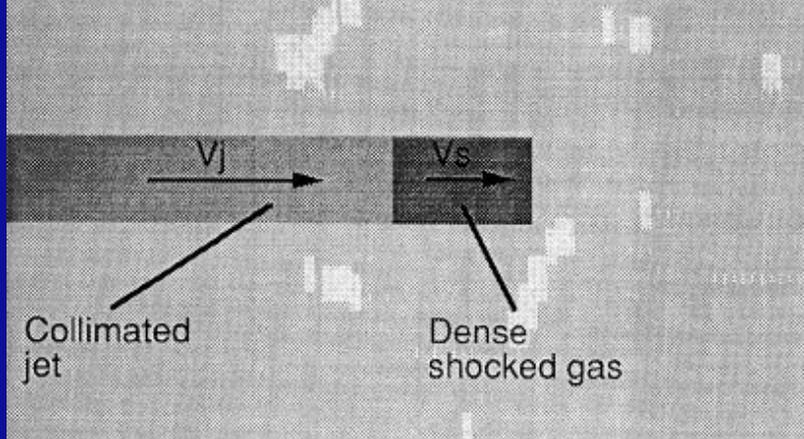
Completely adiabatic shock → large bow shocks with mainly transverse motions

Both wrong → Hence intermediate solution with highly dissipative shock required → forward motion and bow shock!

→ This can accelerate the ambient gas!

Completely radiative shock

→ No bow shock forms, just dense plug at head of shock



Highly dissipative shock: Forward motion AND bow-shock for gas entrainment.

Outflow entrainment models I

Basically 4 outflow entrainment models are discussed in the literature:

Turbulent jet entrainment model

- Working surfaces at the jet boundary layer caused by Kelvin-Helmholtz instabilities form viscous mixing layer entraining molecular gas.
 - The mixing layer grows with time and whole outflow gets turbulent.
- Broken power-law of mass-velocity relation is reproduced, but velocity decreases with distance from source → opposite to observations

Jet-bow shock model

- Jet impacts on ambient gas → bow shocks are formed at head of jet.
 - High pressure gas is ejected sideways
 - broader bow shock entraining the ambient gas.
 - Episodic ejection produces chains of knots and shocks.
- Numerical modeling reproduces many observables, e.g. Hubble-law (outflow velocity increases with distance).

Outflow entrainment models I

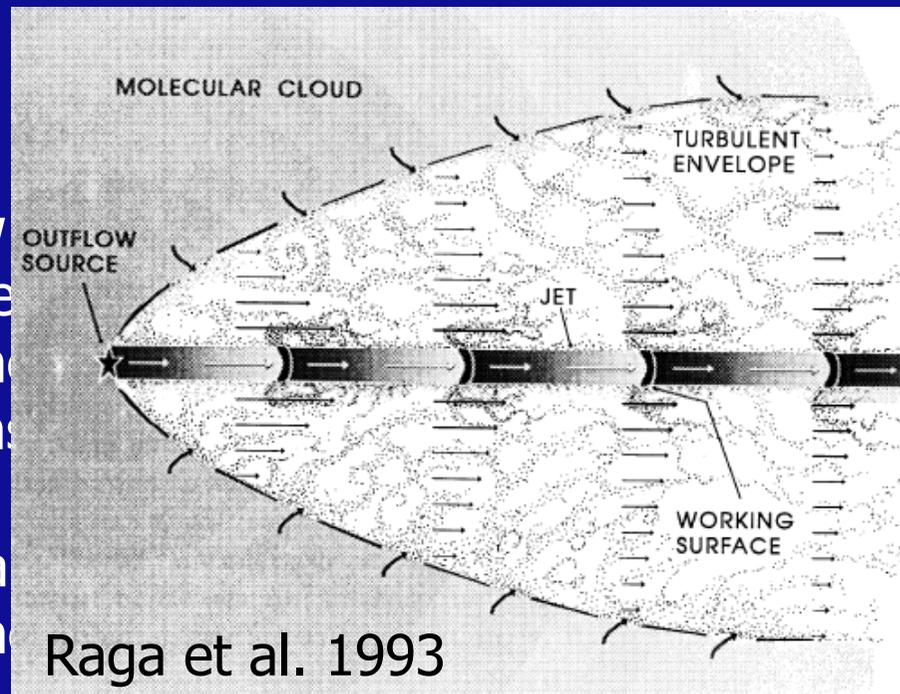
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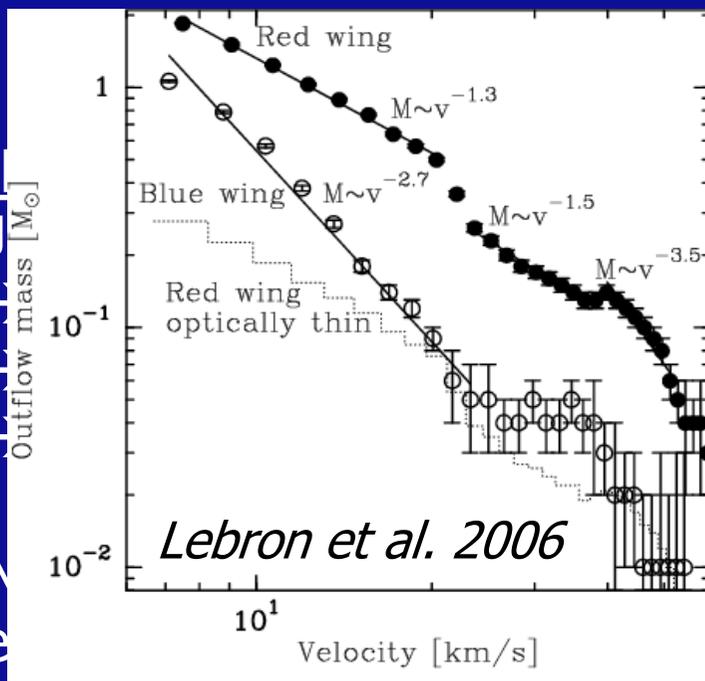
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Jet-

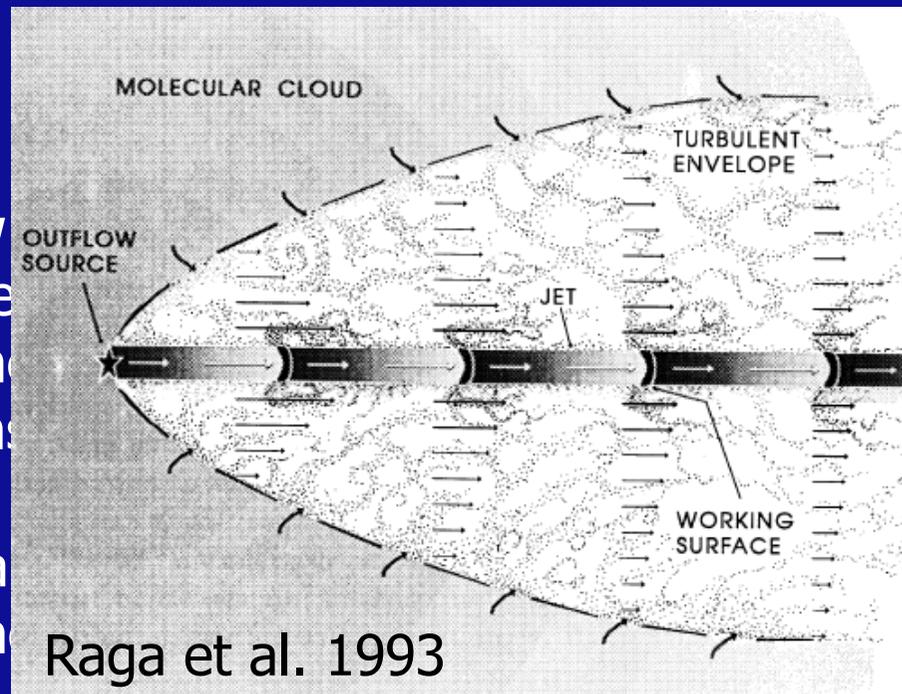
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Outflow entrainment models I

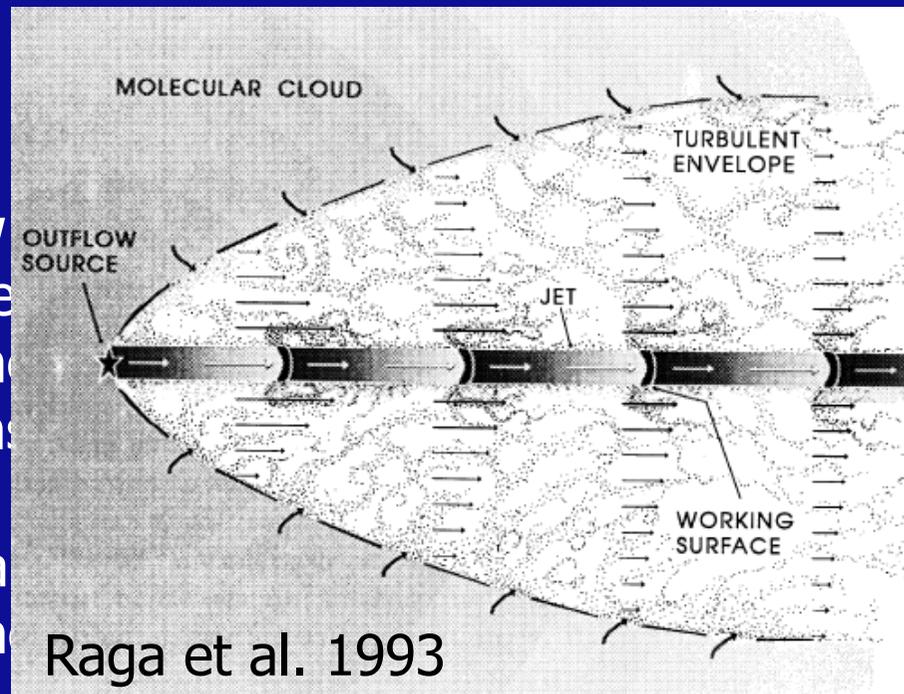
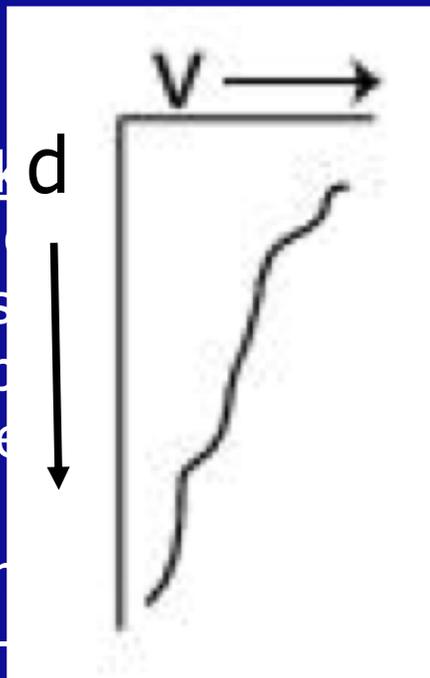
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Jet-bow shock

- Jet impacts → bow shock
→ High pressure and side lobes
→ broader beam
→ Episodic ejection
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e.g. Hubble-
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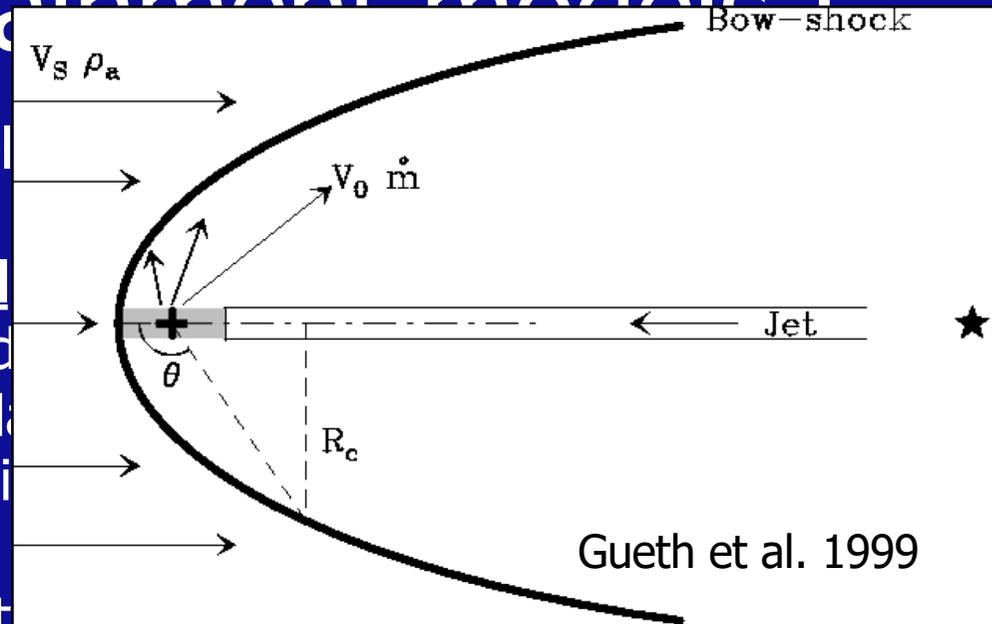


Outflow entrainment models I

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Turbulent jet entrainment model

- Working surfaces at the jet boundaries instabilities form viscous mixing layer
→ The mixing layer grows with time

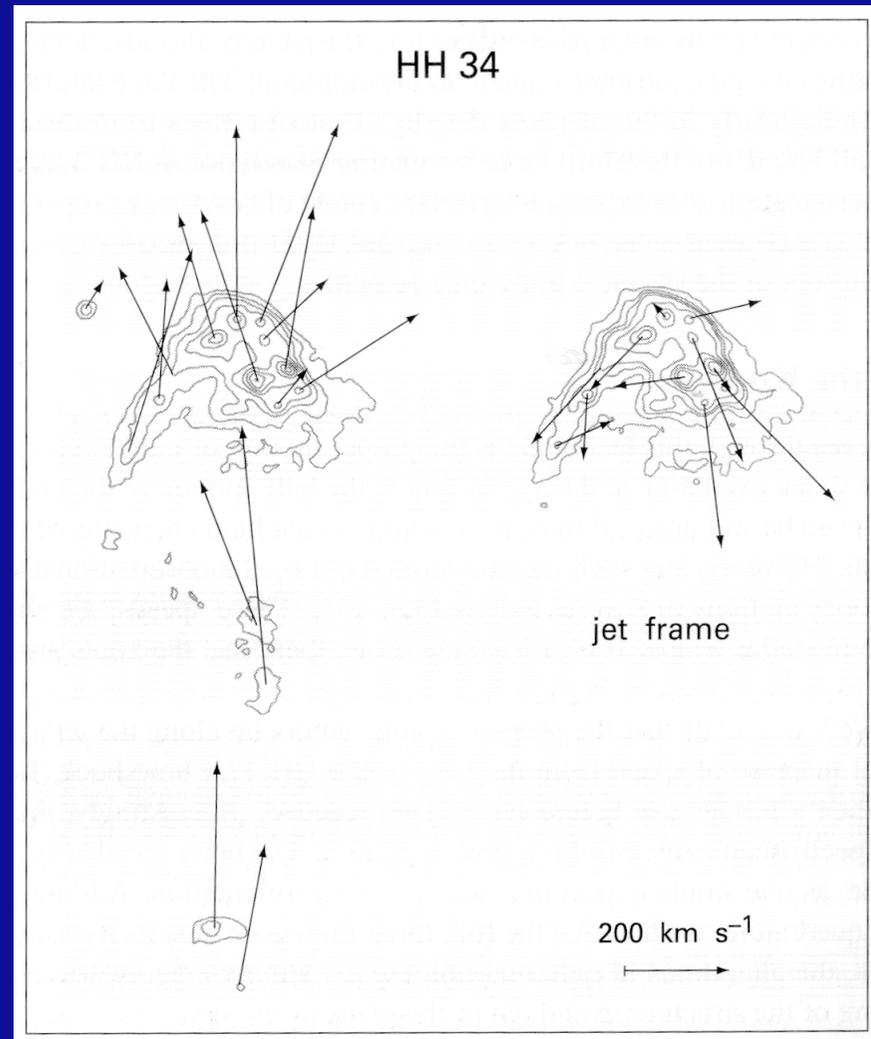


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- Jet impacts on ambient gas → bow shocks are formed at head of jet.
 - High pressure gas is ejected sideways
 - broader bow shock entraining the ambient gas.
 - Episodic ejection produces chains of knots and shocks.
- Numerical modeling reproduces many observables, e.g. Hubble-law (outflow velocity increases with distance).

The case of the HH34 bow shock



In the jet-frame, after subtracting the velocity of the mean axial flow, the knots are following the sides of the bow shock.

Reipurth et al. 2002

Jet simulations I

H_2 1→0 S(1) t = 0 yr

3-dimensional hydrodynamic simulations, including H, C and O chemistry and cooling of the gas, this is a pulsed jet.

CO 0→0 R(1) t = 0 yr

Jet simulations II: small precession

P5 H₂ 1→0 S(1) t = 0 yr

P5 CO 0→0 R(1) t = 0 yr

Jet simulations III, large precession

P20 H₂ 1→0 S(1) t = 0 yr

P20 CO 0→0 R(1) t = 0 yr

Outflow entrainment models II

Wide-angle wind model

- Wide-angle wind blows into ambient gas forming a thin swept-up shell.
- Different degrees of collimation can be explained by different density structures of the ambient gas.
- Attractive models for older and low collimated outflows.

Circulation model

- Molecular gas is not entrained by underlying jet or wind, but it is rather infalling gas that was deflected from the central protostar in a region of high MHD pressure.
- This model was proposed to explain also massive outflows because it was originally considered difficult to entrain that large amounts of gas. Maybe not necessary today anymore ...

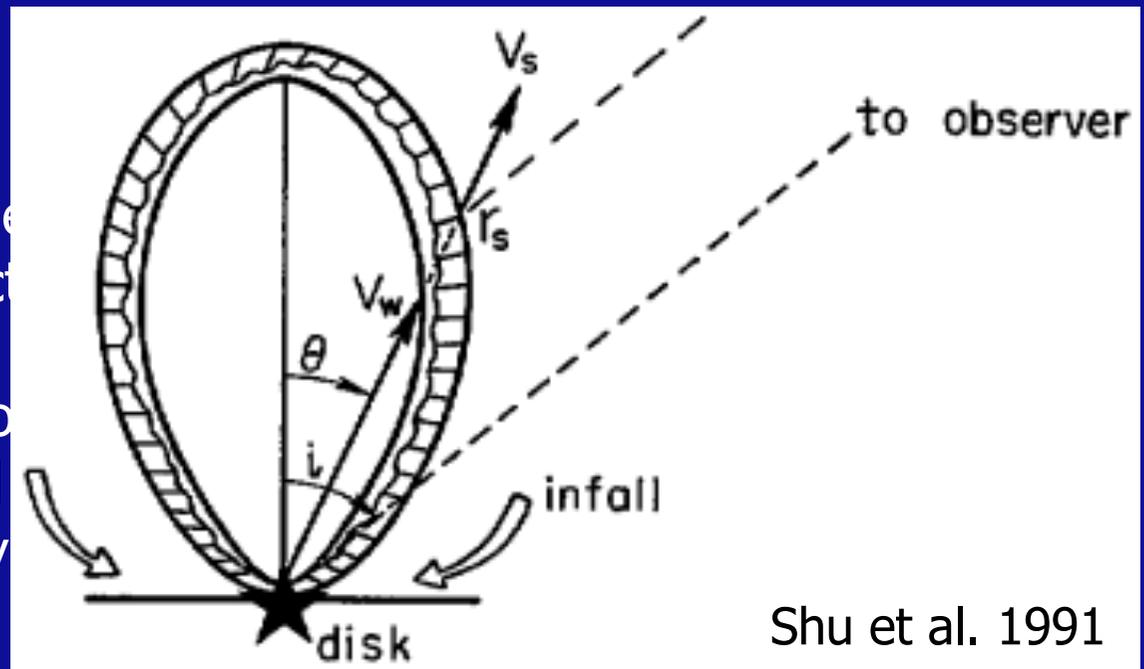
Outflow entrainment models II

Wide-angle wind model

- Wide-angle wind blows into ambient gas forming a thin swept-up shell.
- Different degrees of collimation can be explained by different density structures of the ambient gas.
- Attractive models for older and low collimated outflows.

Circulation model

- Molecular gas is not entrained by the outflow; instead, it is infalling gas that was deflected by the outflow of high MHD pressure.
- This model was proposed to explain the collimation of outflows; it was originally considered for the case of a star with a disk. Maybe not necessary today.



Shu et al. 1991

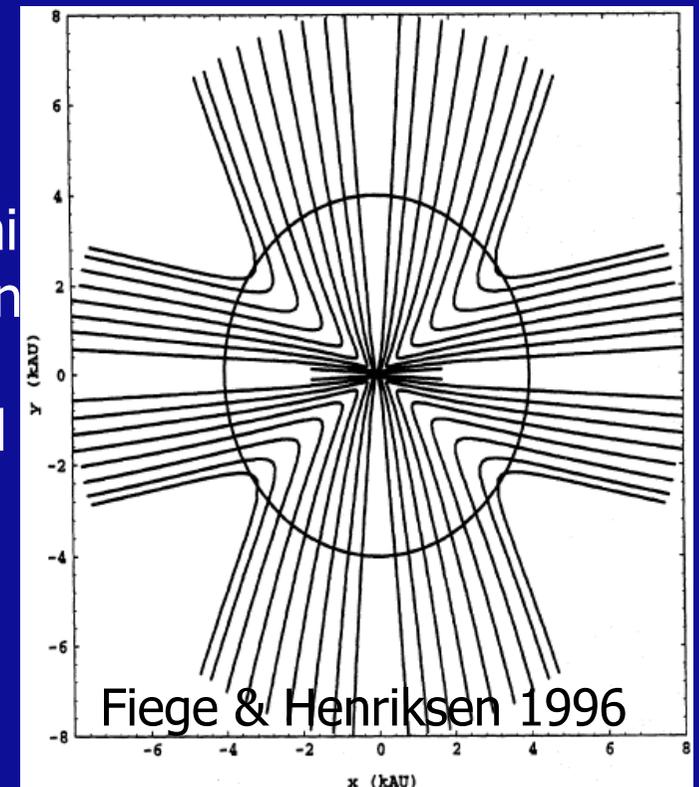
Outflow entrainment models II

Wide-angle wind model

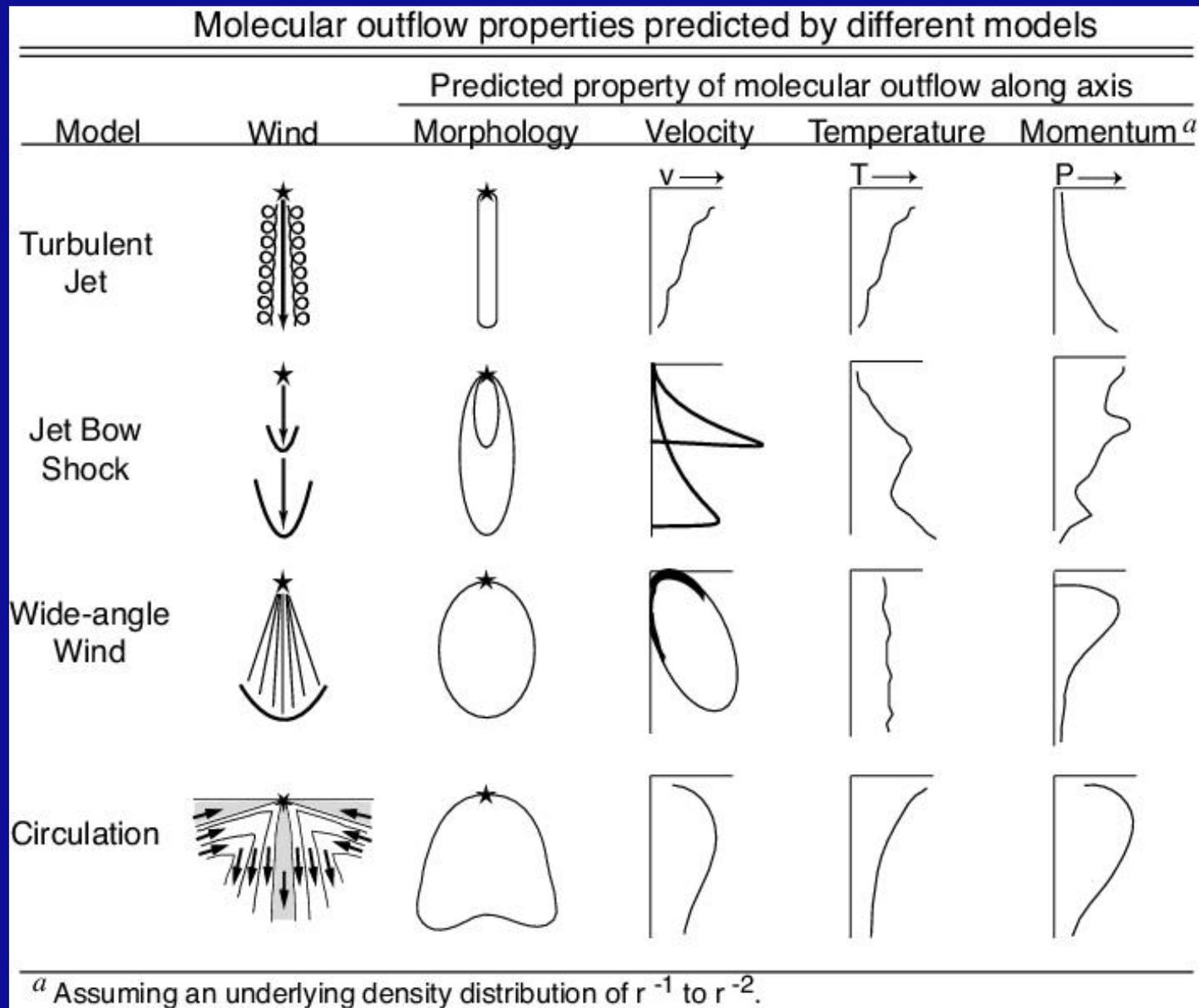
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Circulation model

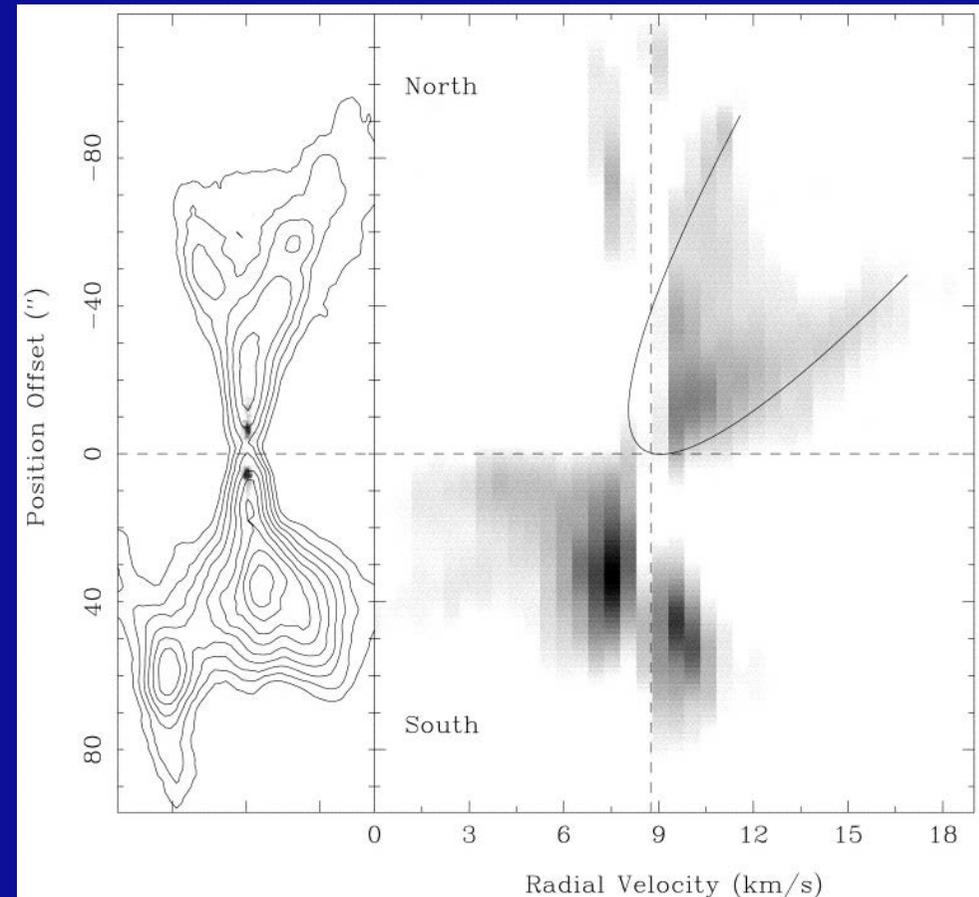
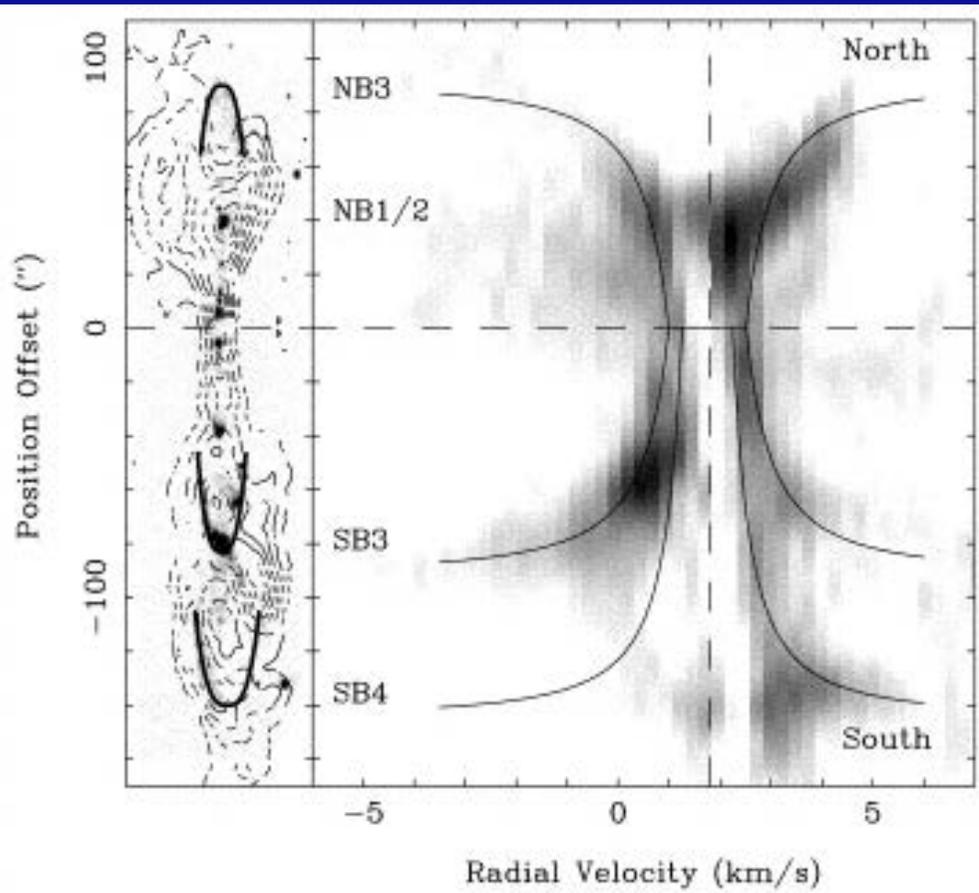
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Outflow entrainment models III



Collimation and pv-structure



HH212: consistent with jet-driving

VLA0548: consistent with wind-driving

- pv-structure of jet- and wind-driven models very different
- Often Hubble-law observed → increasing velocity with increasing distance from the protostar

Lee et al. 2001

Summary

- End of protostellar/beginning or pre-main sequence evolution → birthline.
 - Pre-main sequence evolution in the Hertzsprung-Russel (HR) diagram.
 - Connection of HR diagram with protostellar and pre-main sequence classes.
-
- Outflows and jets are ubiquitous and necessary phenomena in star formation.
 - Transport angular momentum away from protostar.
 - They are formed by magneto-centrifugal disk-winds.
 - Collimation is caused by Lorentz forces.
 - Gas entrainment can be due to various processes: turbulent entrainment, bow-shocks, wide-angle winds, circulation ...
 - They inject significant amounts of energy in the ISM, may be important to maintain turbulence and disrupt their maternal clouds.

Sternentstehung - Star Formation

Winter term 2020/2021

Henrik Beuther, Thomas Henning & Sümeyye Suri

<i>03.11 Today: Introduction & Overview</i>	<i>(Beuther)</i>
<i>10.11 Physical processes I</i>	<i>(Beuther)</i>
<i>17.11 Physical processes II</i>	<i>(Beuther)</i>
<i>24.11 Molecular clouds as birth places of stars</i>	<i>(Suri)</i>
<i>01.12 Molecular clouds (cont.), Jeans Analysis</i>	<i>(Suri)</i>
<i>08.12 Collapse models I</i>	<i>(Henning)</i>
<i>15.12 Collapse models II</i>	<i>(Henning)</i>
----- Christmas break -----	
<i>12.01 Protostellar evolution</i>	<i>(Beuther)</i>
<i>19.01 Pre-main sequence evolution & outflows/jets</i>	<i>(Beuther)</i>
26.01 Accretion disks I	(Henning)
02.02 Accretion disks II	(Henning)
09.02 High-mass star formation, clusters and the IMF	(Suri)
16.02 Extragalactic star formation	(Henning)
23.02 Examination week, no star formation lecture	

Book: Stahler & Palla: The Formation of Stars, Wileys

More Information and the current lecture files: http://www.mpia.de/homes/beuther/lecture_ws2021.html

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