

Sternentstehung - Star Formation

Winter term 2024/2025

Henrik Beuther, Thomas Henning & Caroline Gieser

15.10 Today: Introduction & Overview	(Beuther)
22.10 Physical processes I	(Beuther)
29.10 --	
05.11 Physical processes II	(Beuther)
12.11 Molecular clouds as birth places of stars	(Beuther)
19.11 Molecular clouds (cont.), Jeans Analysis	(Henning)
26.11 Collapse models I	(Beuther)
03.12 Collapse models II	(Beuther)
10.12 Protostellar evolution	(Gieser)
17.12 Pre-main sequence evolution & outflows/jets	(Henning)
07.01 Accretion disks I	(Henning)
14.01 Accretion disks II	(Henning)
21.01 High-mass star formation, clusters and the IMF	(Gieser)
28.01 Extragalactic star formation	(Henning)
04.02 Planetarium@HdA, outlook, questions	
11.02 Examination week, no star formation lecture	(Beuther, Gieser, Henning)

Book: Stahler & Palla: The Formation of Stars, Wileys

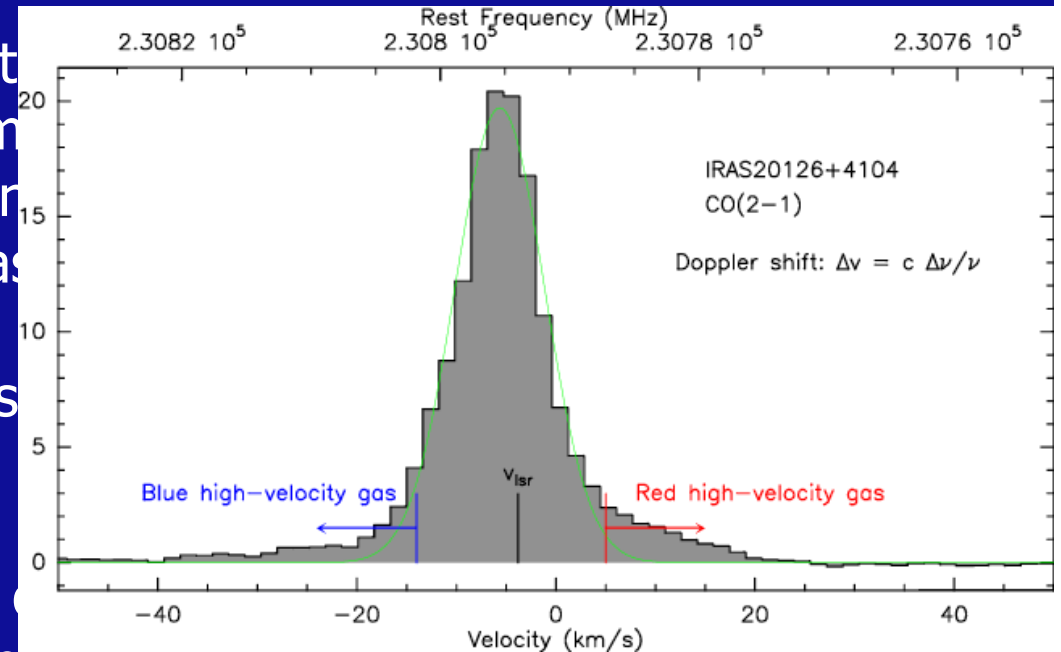
More Information and the current lecture files: http://www.mpia.de/homes/beuther/lecture_ws2425.html
beuther@mpia.de, henning@mpia.de, gieser@mpia.de

Last lecture

- Line profiles (thermal and kinematic broadening) and some applications
- Magnetic fields are very important but difficult to measure:
 - Zeeman effect traces **B** component along line of sight.
 - Dust polarisation traces **B** in plane of the sky.
(Other magnetic field measurements possible.)
- Masers are non-thermal processes. Good for high spatial accuracy and proper motion studies.
- Dust important from many points of view:
 - Traces warm and cold components of ISM.
 - Important coolant at high densities.
 - Traces magnetic field.
 - Chemical catalyst.

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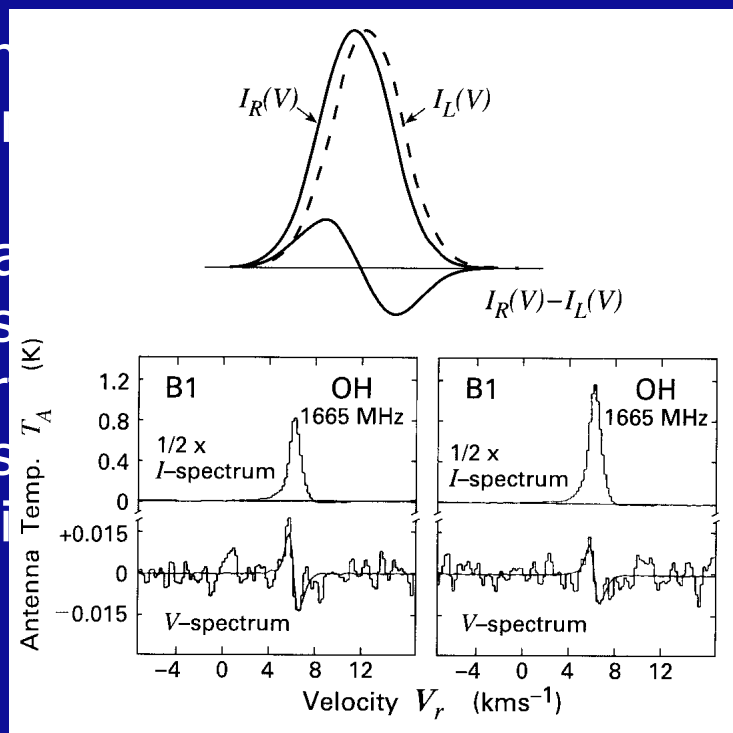
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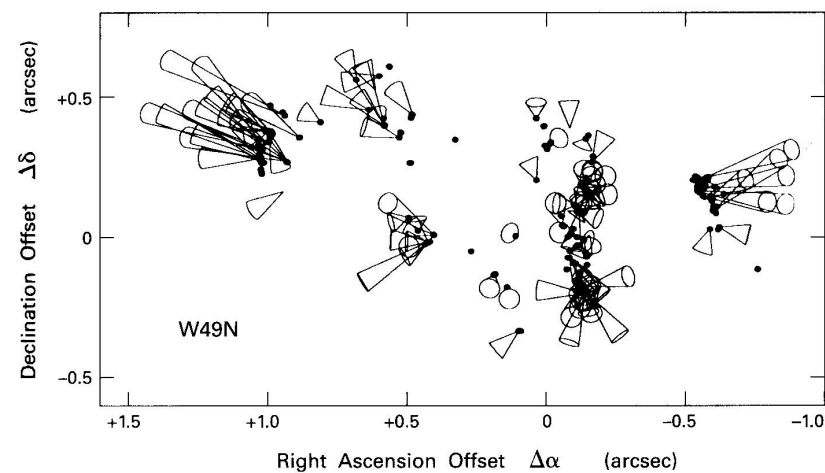
- Dust important for tracing of ISM.

- Traces
- Important
- Traces
- Chemical

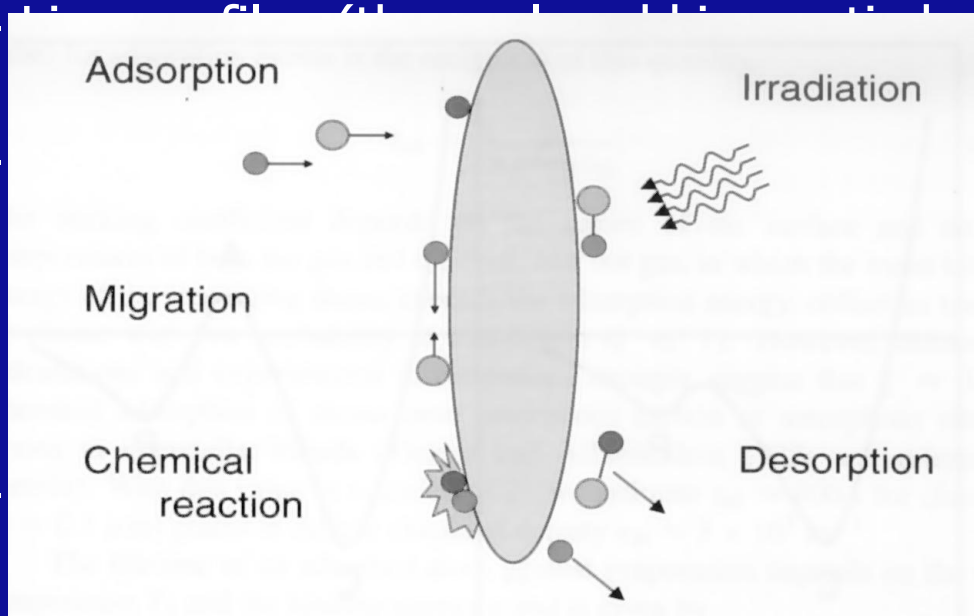


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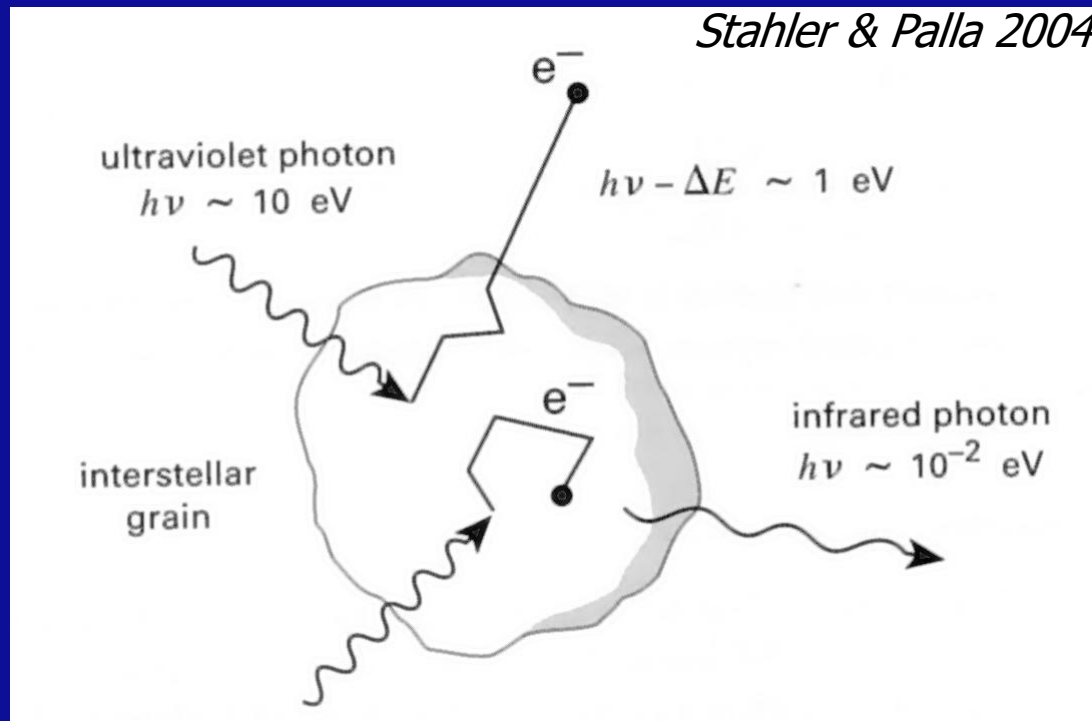
Last lecture



- ... (H₂ formation, grain growth, and some applications)
- ... difficult to measure: (not along line of sight, ... of the sky. ... nents possible.)
- ... and for high spatial accuracy and

- Dust important from many points of view:
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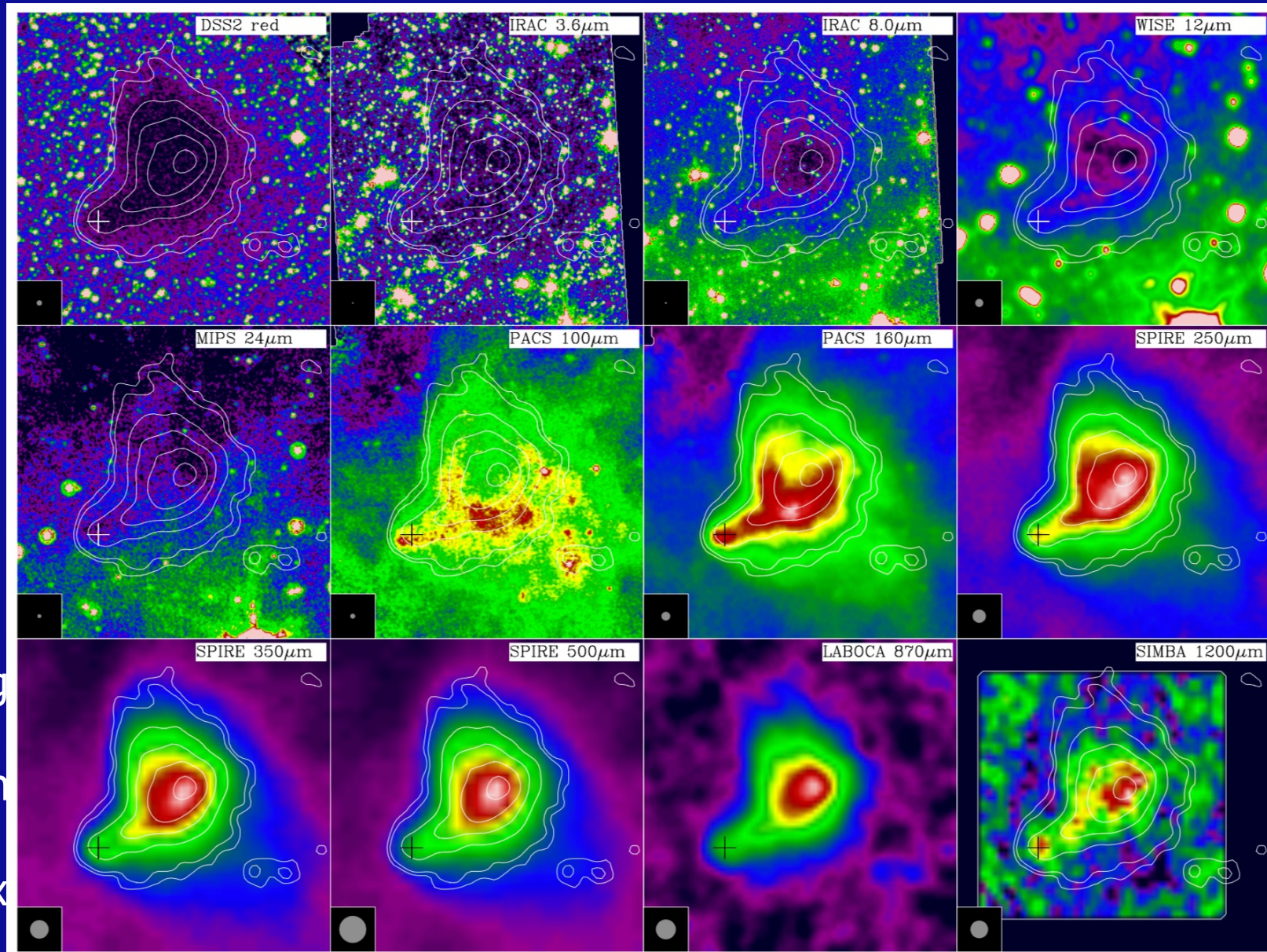
Dust action at longer wavelengths: Re-emission



Dust grain hit by UV photon:

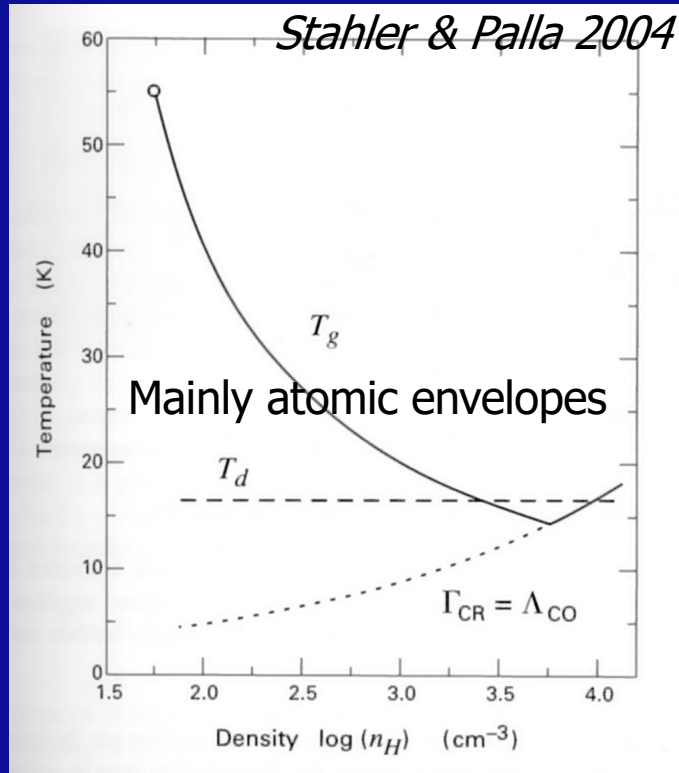
- 1) Photoelectrical effect \rightarrow give energy to e^- \rightarrow leaves grain and heats gas.
- 2) Excites lattice vibrations \rightarrow transformed to (far)-IR photons and re-emitted.

Dust action at longer wavelengths: Re-emission

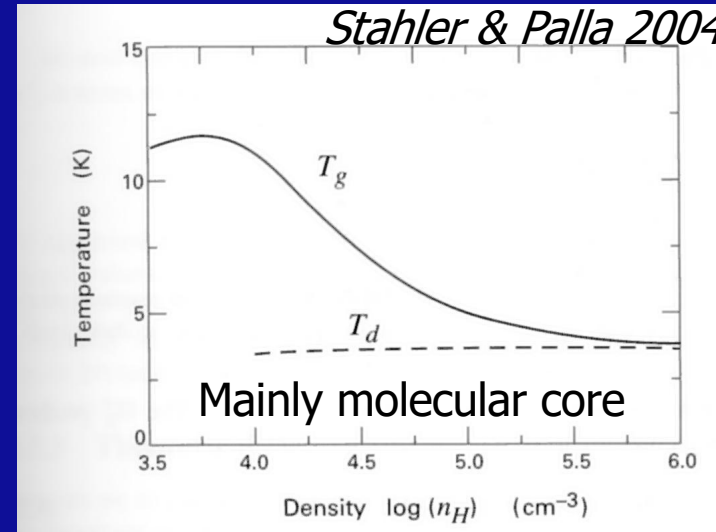


Nielbock et al. 2012, Contours 870 μ m

Dust and gas coupling



Cosmic ray heating more sensitive to density than CO cooling. Hence T_g rises again.

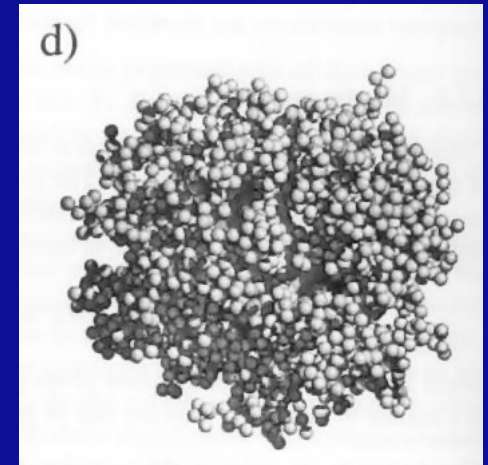
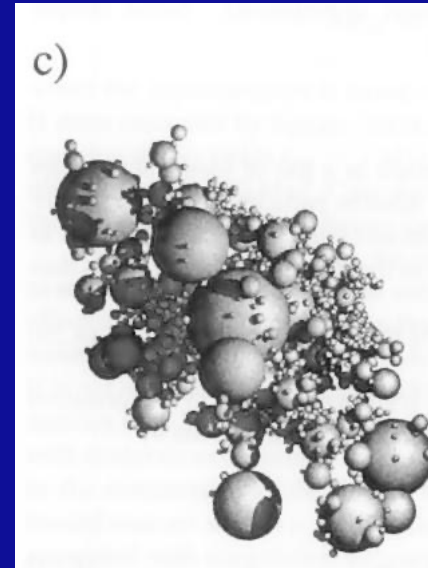
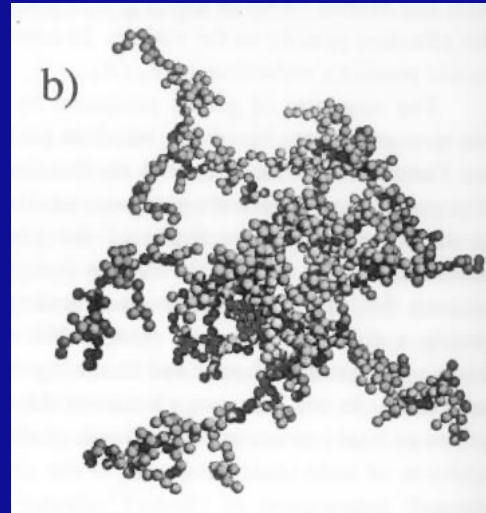
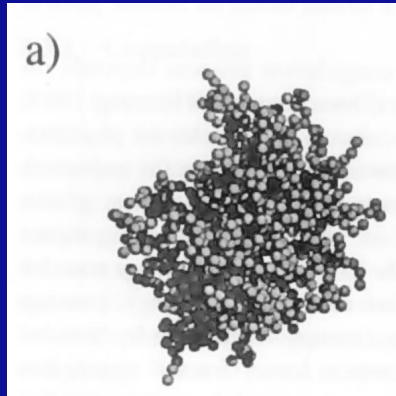


simplified model \rightarrow lowest obs. $T_d \sim 7-10\text{K}$

Γ – heating rate Λ – Cooling rate

- Low densities: gas and dust de-coupled; at high densities coupled.
- Low densities gas cooling mainly CO; high densities via CO & dust.
- At very high densities gas and dust temperatures approach each other
 \rightarrow CO cooling becomes insignificant then!

Dust incarnations



Dust can grow and coagulate in very dense environments, e.g., disks.

Figures: Simulations of dust grain cluster growth for different initial parameters (gas and dust density, temperature, stickiness, grain charge, coagulation time ...).
(From Dorschner & Henning 1995)

Topics today

- Physical distributions (cont.)
- Components of the interstellar medium
- General characteristics of molecular clouds
- Important cloud relations
- Cloud fragmentation

Physical conditions : Micro-Level

A medium in thermodynamic eq. can be described by 4 distribution laws:

1.) MAXWELL distribution of the particle velocity contributions (kinetic energy):

$$N(v;T) = 4\pi \left(\frac{m}{2kT} \right)^{3/2} v^2 \exp\left(-\frac{mv^2}{2kT} \right) \quad \nu: \text{particle velocities}$$

2.) BOLTZMANN distribution of the population numbers of the particle energy levels:

$$\frac{N_o}{N_u} = \frac{g_o}{g_u} \exp\left(-\frac{E_o - E_u}{kT} \right) \quad \begin{array}{l} E_{o/u} \longrightarrow \text{Energies of the upper (o) and lower (u) levels} \\ g_{o/u} \longrightarrow \text{Corresponding statistical weights} \end{array}$$

3.) PLANCK radiation law (distribution of the photon energies):

$$B_\nu = \frac{2h\nu^3 / c^2}{\exp(h\nu / kT) - 1} \quad \nu: \text{photon frequencies}$$

4.) SAHA equation (distribution of the ionisation levels in plasma):

$$\frac{N_{j+1} N_e}{N_j} = \frac{2 U_{j+1}(T)}{U_j(T)} \left(\frac{2\pi m_e kT}{h^2} \right)^{3/2} \exp(-\chi_{j,j+1} / kT)$$

N_{j+1}, N_j - Number densities of (j+1)-fold and j-fold ionised particles

N_e - electron density

$\chi_{j,j+1}$ - ionisation energy needed to get from ionisation level j to $j+1$

U_{j+1}, U_j - partition function for both states

Physical conditions : Micro-Level

Are these distribution functions valid in the ISM?

General rule: time scale for processes leading to equilibrium short compared to time scales of disturbing processes

1. Example : Collisions between H-atoms:

Consider:

$T = 100 \text{ K} \rightarrow \text{mean } v \sim 1 \text{ km/s}$; cross section $\sigma = \pi R_H^2 \sim \pi (0.1 \text{ nm})^2$

Probability for collision: $P = v \sigma n_H \tau_s$

\rightarrow average time τ_s between two collision for $P=1$: $\tau_s = (v \sigma n_H)^{-1}$

\rightarrow with HI density of $1 \text{ cm}^{-3} \rightarrow \tau_s \sim 1000 \text{ yrs}$

\rightarrow short compared to most interstellar processes (except shock fronts)

\rightarrow Maxwell distribution valid, introduction of kinetic temp. T_{kin} reasonable

Physical conditions : Micro-Level

2. Example: Balance for energy level population numbers for ISM:

Correction factor to Boltzmann:
$$\frac{1}{1 + (A_{21} / (n Q_{21}))}$$

(Pure Boltzmann only if $(n Q_{21}) \gg A_{21}$)

A_{21} [s^{-1}] Einstein coefficient for
spontaneous radiative decay
 Q_{21} [$m^3 s^{-1}$] collision rate
 n [m^{-3}] number density

- In thin ISM collision rate small (Example 1) \rightarrow sub-thermal

- For dense cores: E.g. CO(1-0) at density $10^5 cm^{-3}$: $A_{21} = 7.2 \times 10^{-8} s^{-1}$,
 $Q_{21} = 3.3 \times 10^{-11} cm^3 s^{-1}$

$\rightarrow A_{21} / (n Q_{21}) \sim 0.02$

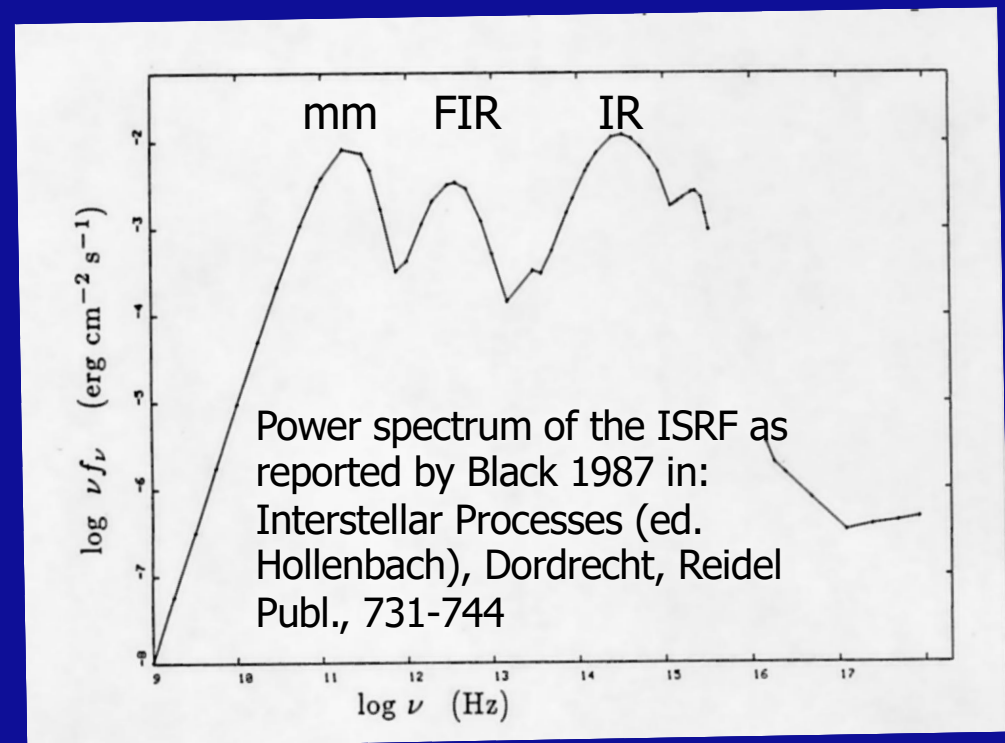
\rightarrow Boltzmann distribution valid in dense cores!

Physical conditions : Micro-Level

3. Example : Interstellar radiation field (ISRF) :

Sum of emission contributions from all emitting objects (stars, dust, gas) in the nearer and further vicinity of the gas cloud

ISRF cannot be approximated by a black body (i.e., Planck function not applicable)
ISRF hence far from thermodynamic equilibrium ...



However: Dense cores and stars can be fitted relatively well with single or multiple black body functions.

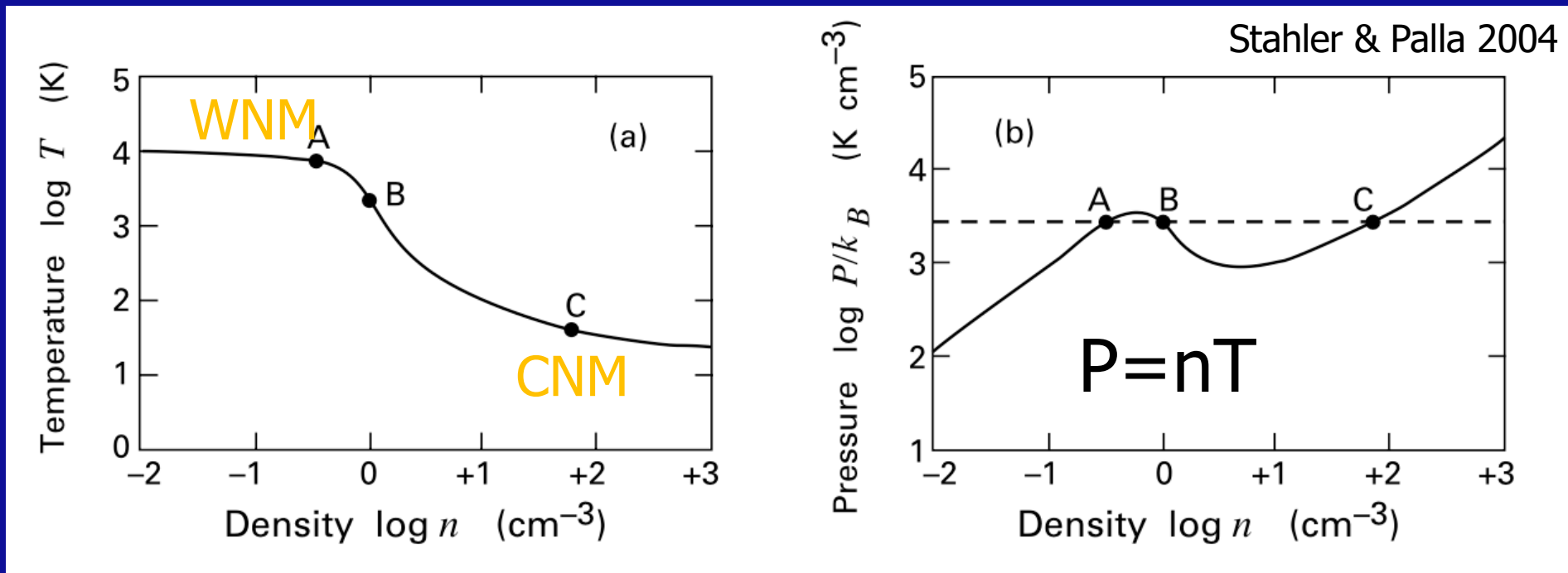
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- **Components of the interstellar medium**
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Historical Models of the ISM (I)

1.) Simple Two-Phase Model (Field et al. 1969)

- Considering only the atomic gas
- Underlying idea: pressure equilibrium with environment

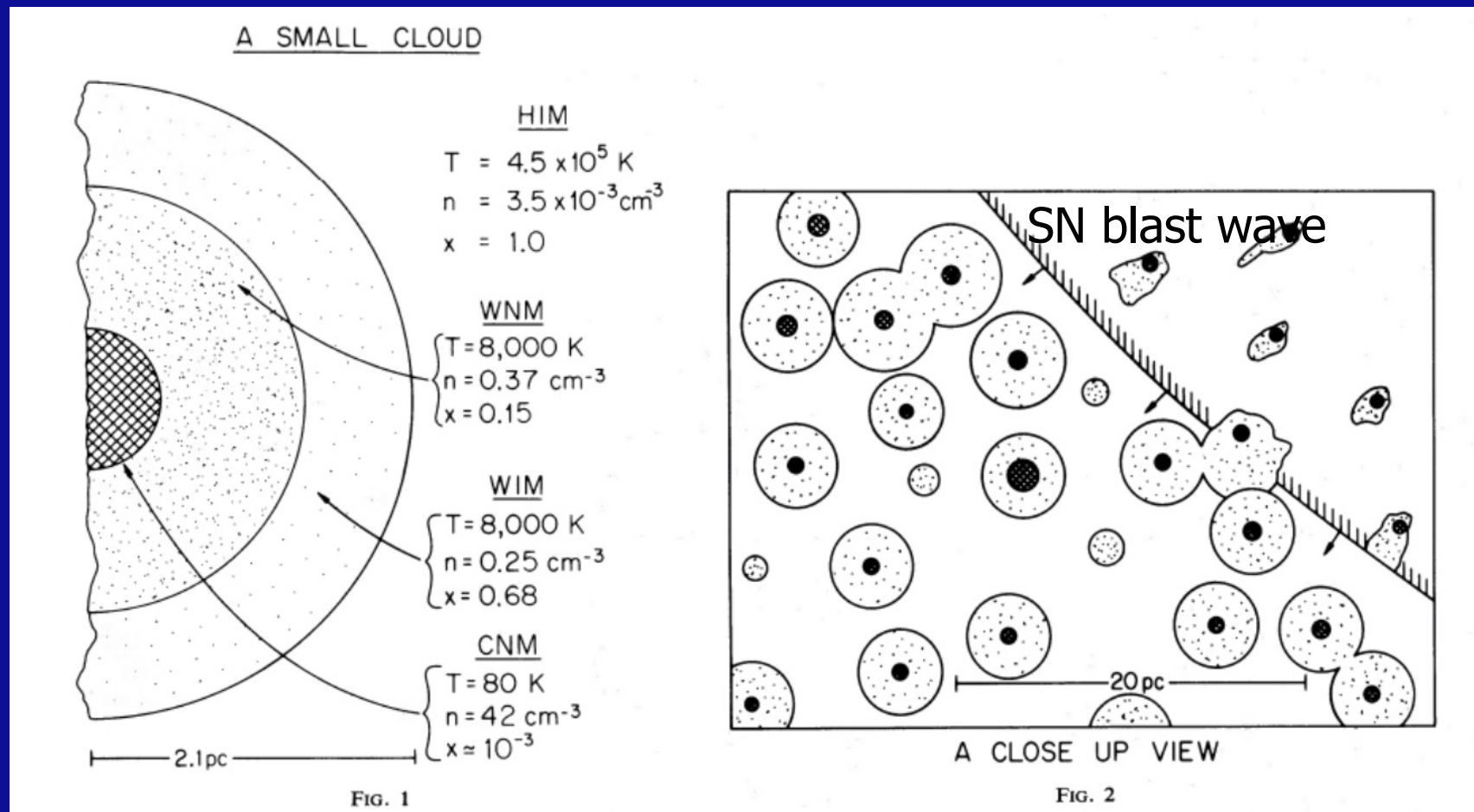


Shortcomings: does not account for hot ionized nor molecular medium

Historical Models of the ISM (II)

2.) Three-Phase Model (McKee & Ostriker 1977)

- Takes into account hot component of ISM and supernova blast waves.
- More dynamical and coupled to the formation (and death) of massive stars



Historical Models of the ISM (III)

2.) Three-Phase Model (McKee & Ostriker 1977)

Shortcomings in the original model:

- SN rate and SN “luminosity” overestimated, SNe not arbitrarily distributed
- Observations indicate considerable amount of evenly distributed (i.e., not bound to clouds) warm HI gas
- Model still assumes global pressure equilibrium between the phases
- One very important component still missing: **molecular clouds !!**
($T \approx 10 \text{ K}$, $n > 300 \text{ cm}^{-3}$)

Phase transitions are possible (e.g., by heating and cooling)

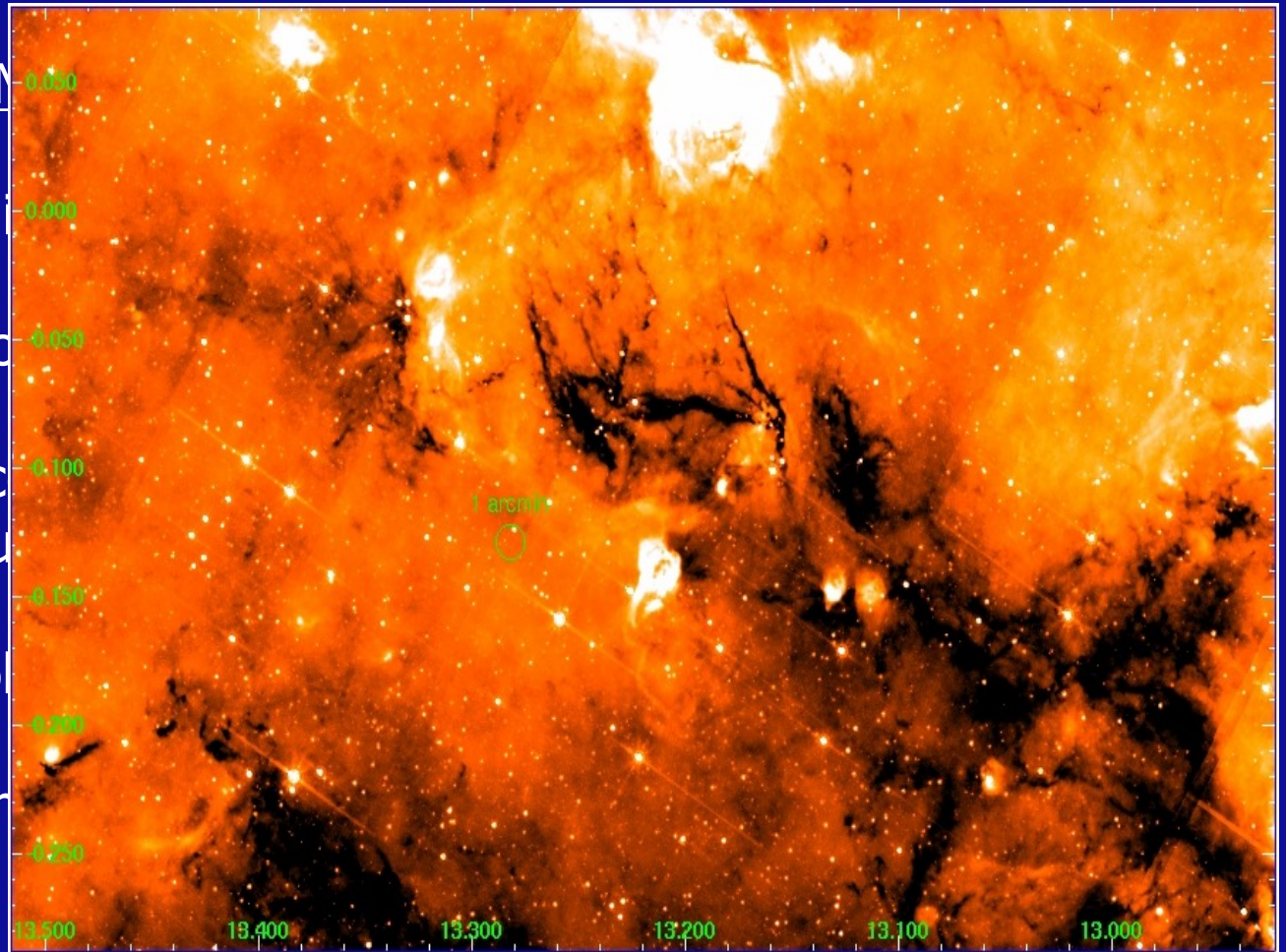
Diffuse clouds \longrightarrow molecular clouds \longrightarrow stars

Historical Models of the ISM (III)

2.) Three-Phase Model (M

Shortcomings in the original

- SN rate and SN "luminosity"
- Observations indicate clouds are not bound to clouds (i.e., not bound to clouds)
- Model still assumes global equilibrium
- One very important component



Phase transitions are possible (e.g., by heating and cooling)

Diffuse clouds → molecular clouds → stars

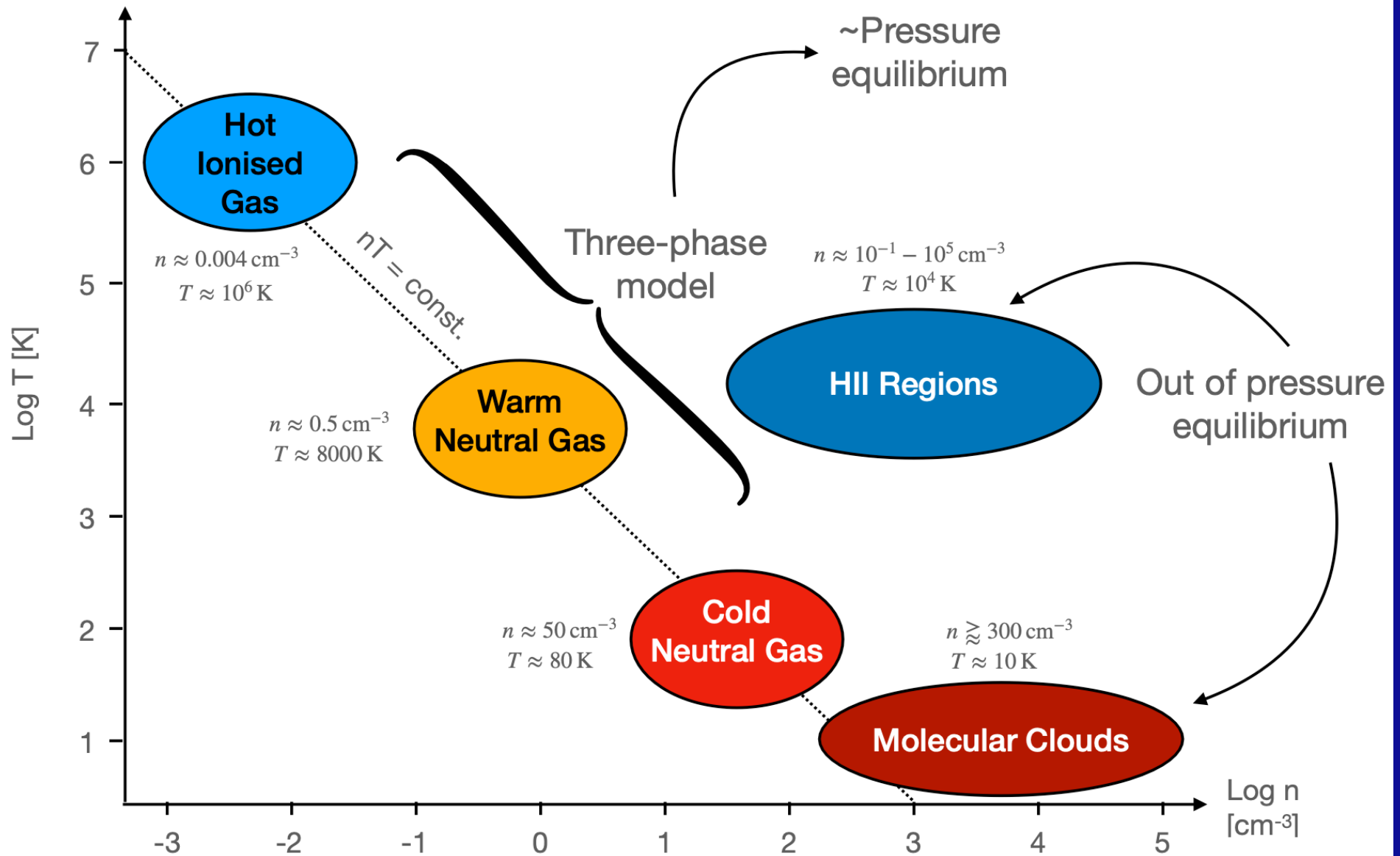
3.) Overview of the components

Phase	n [cm^{-3}]	T [K]	f	M [$10^9 M_{\odot}$]
Hot ionised medium	0.003	10^6	0.5	0.1
Warm ionised medium	0.3	8000	0.1	1.0
Warm neutral medium	0.5	8000	0.3...	1.4
Diffuse HI clouds	50	80	-	2.5
Molecular clouds	>300	10	-	2.5
HII regions	$1 - 10^5$	10^4	-	0.05

f as volume filling factor regarding the Galactic disk

By mass in Milky Way: ~20 ionized, 60% neutral, 20% molecular

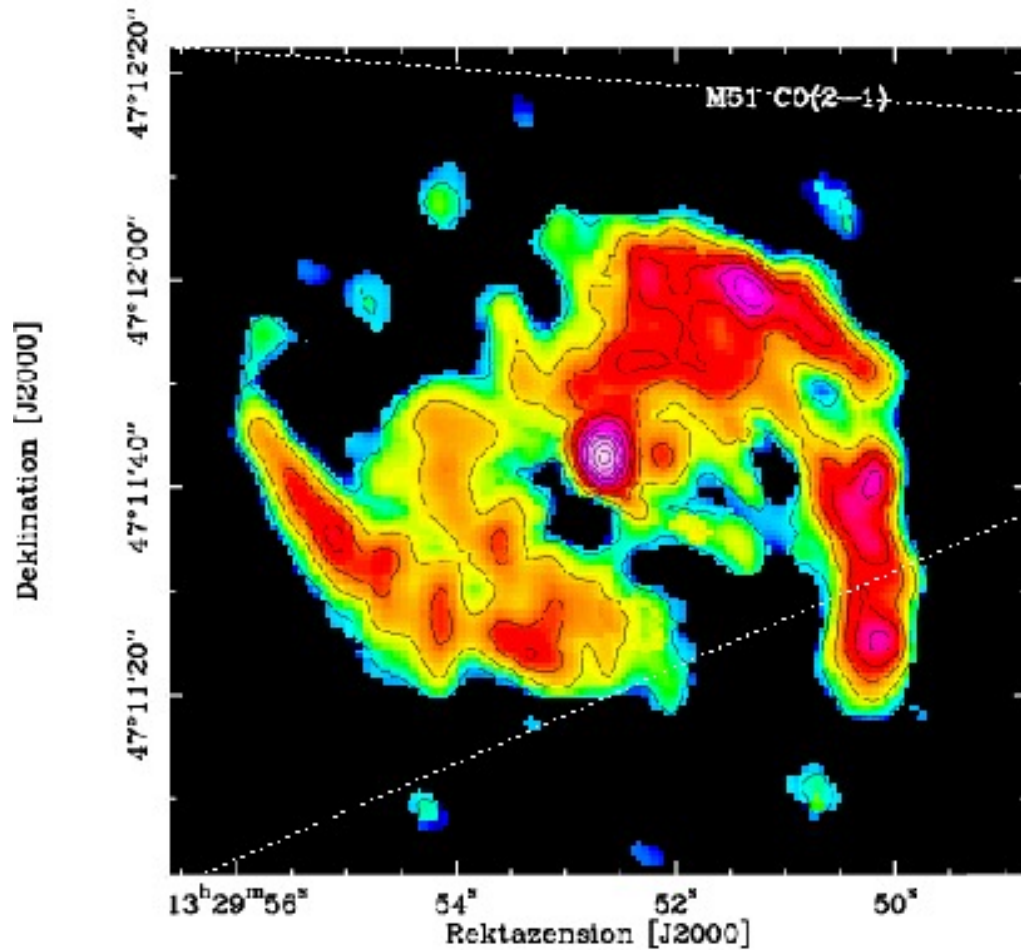
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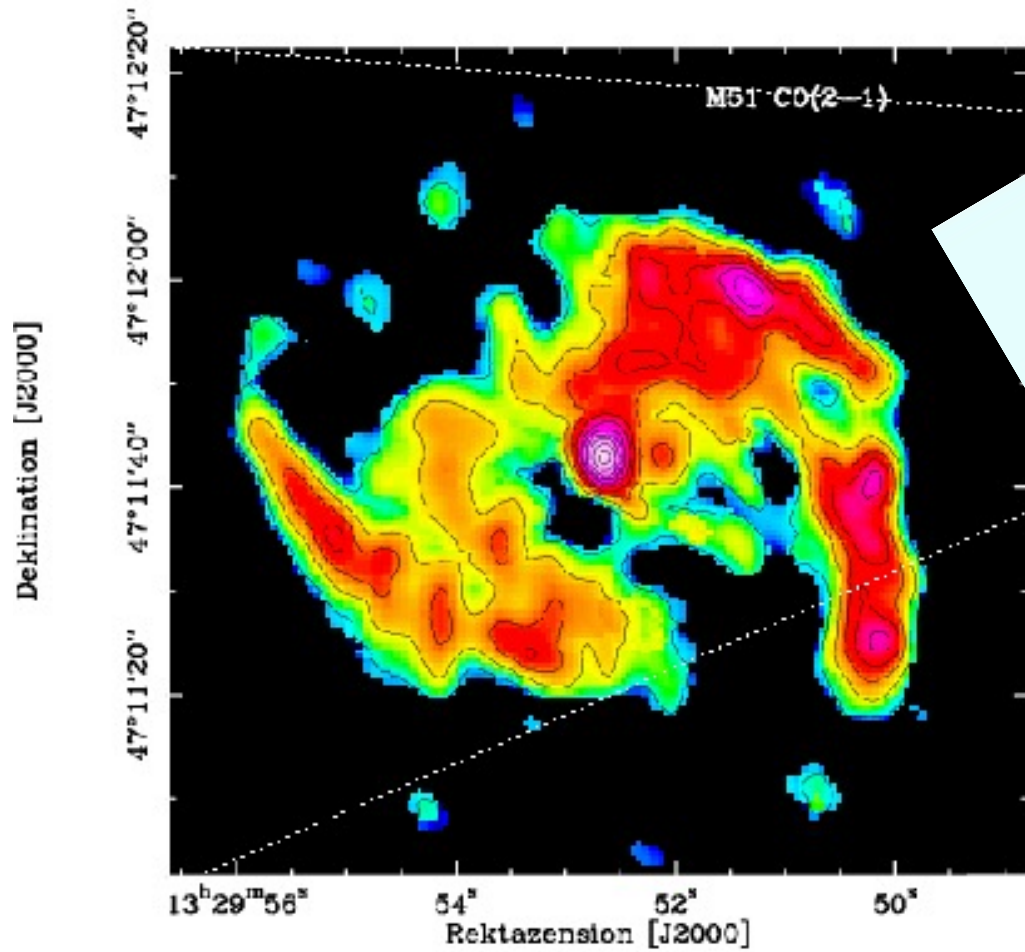
M51: The Whirlpool Galaxy



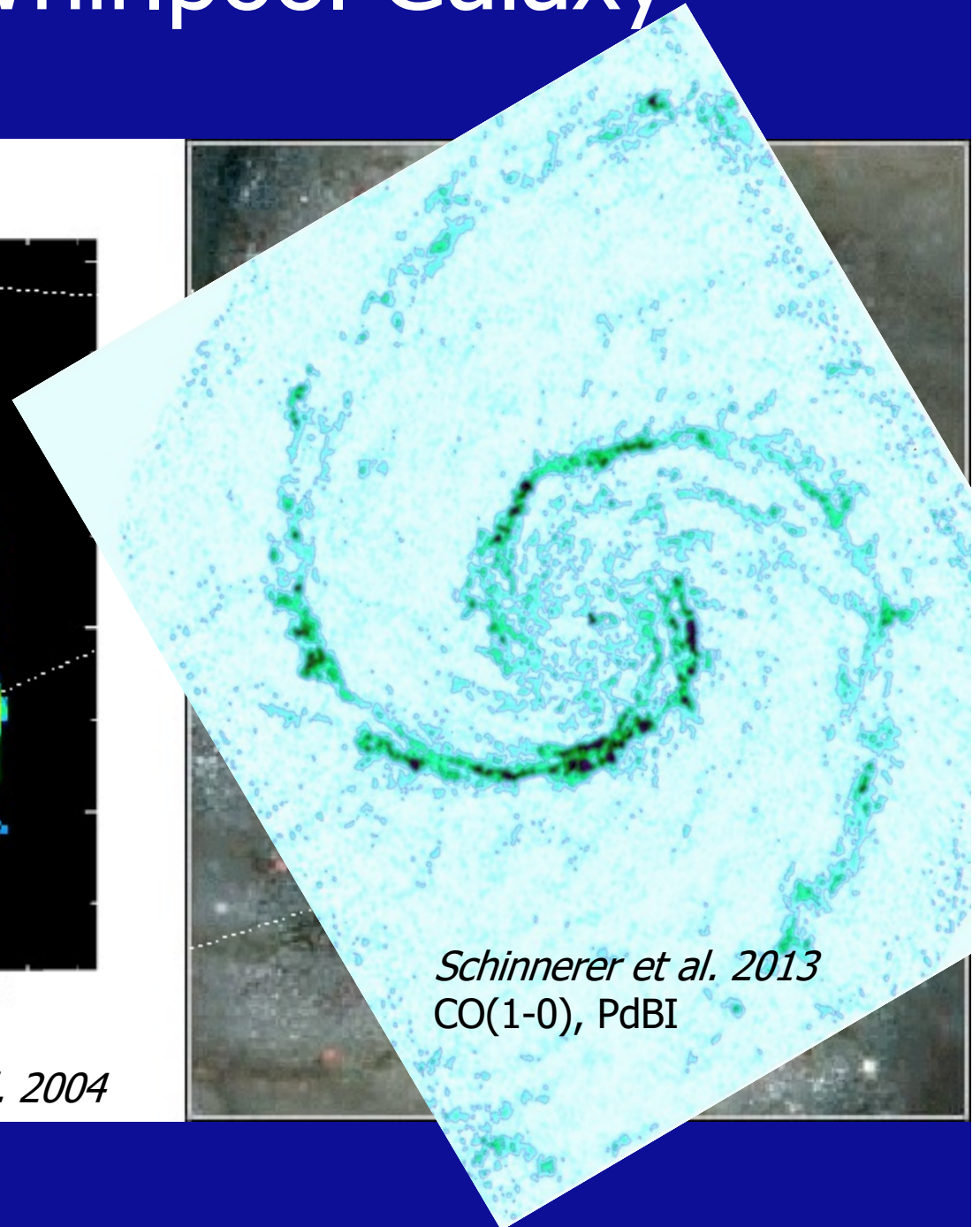
Matsushita et al. 2004



M51: The Whirlpool Galaxy

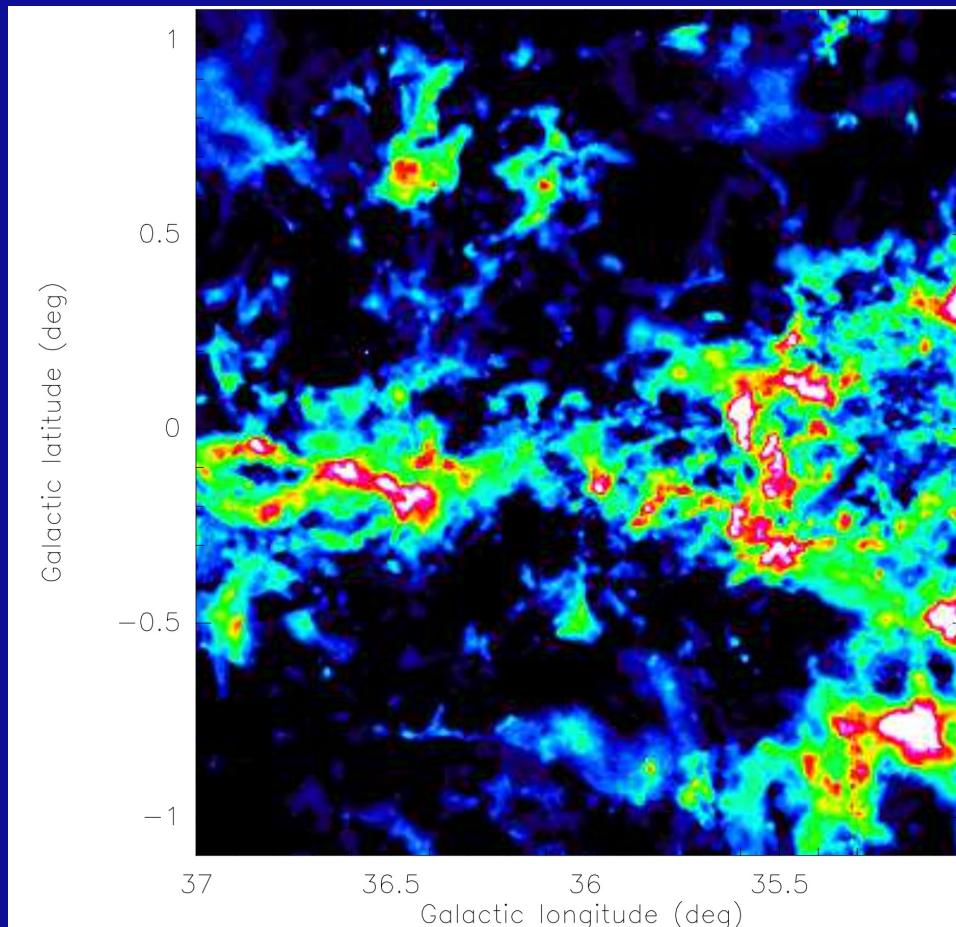


Matsushita et al. 2004



Schinnerer et al. 2013
CO(1-0), PdBI

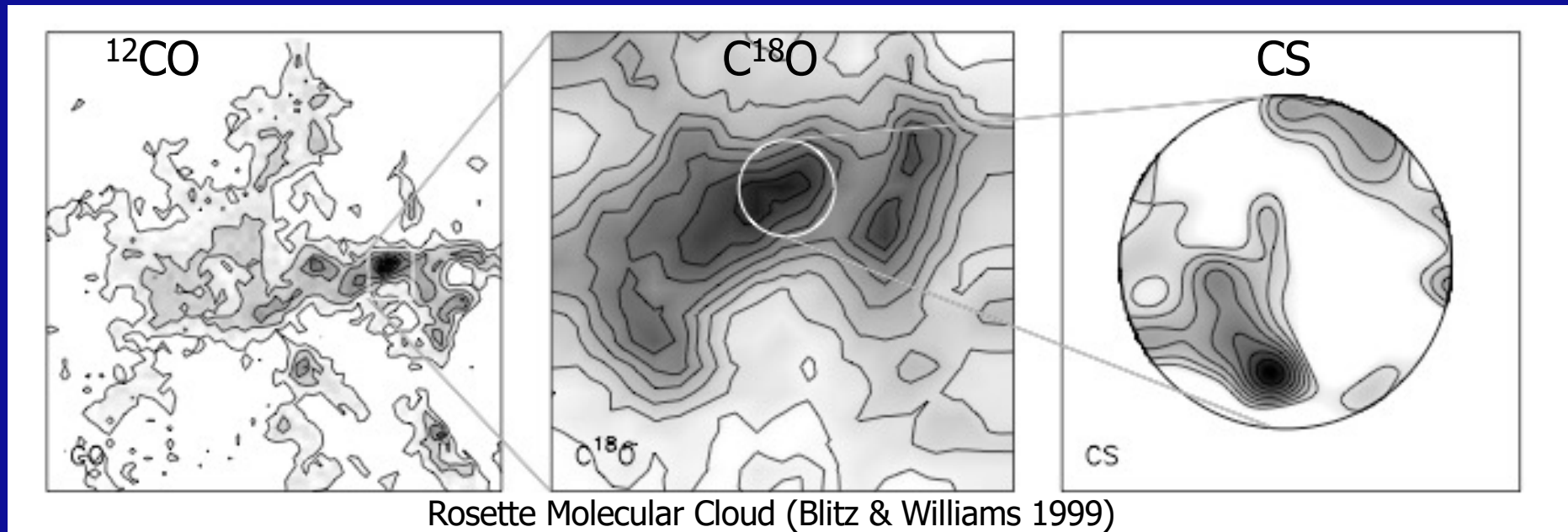
Giant Molecular Clouds



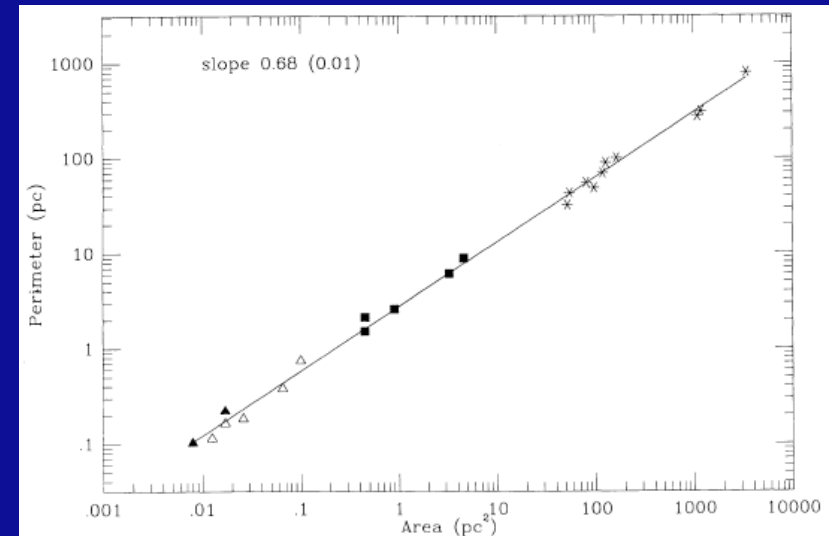
Galactic Ring survey
 $^{13}\text{CO}(2-1)$
Jackson et al. 2006

Sizes: 20 to 100pc; Masses: 10^4 to $10^6 M_{\text{sun}}$; Temperatures: 10 to 20K
Supersonic velocity dispersion $\sim 2-3$ km/s mainly due to turbulence
Magnetic field strengths on the order of $10\mu\text{G}$
Average local densities $\sim 10^4\text{cm}^{-3}$; Volume-averaged densities $\sim 10^2\text{cm}^{-3}$
--> highly clumped material

Hierarchical cloud structure

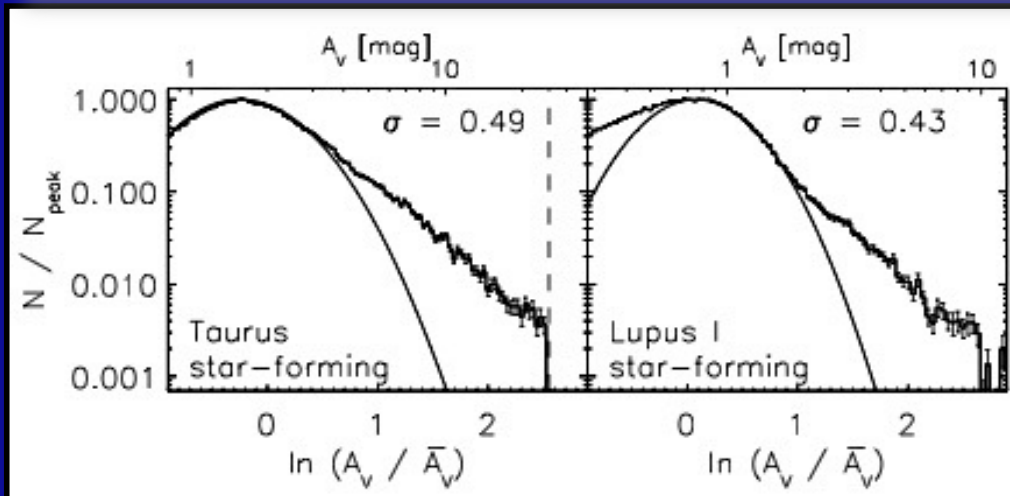
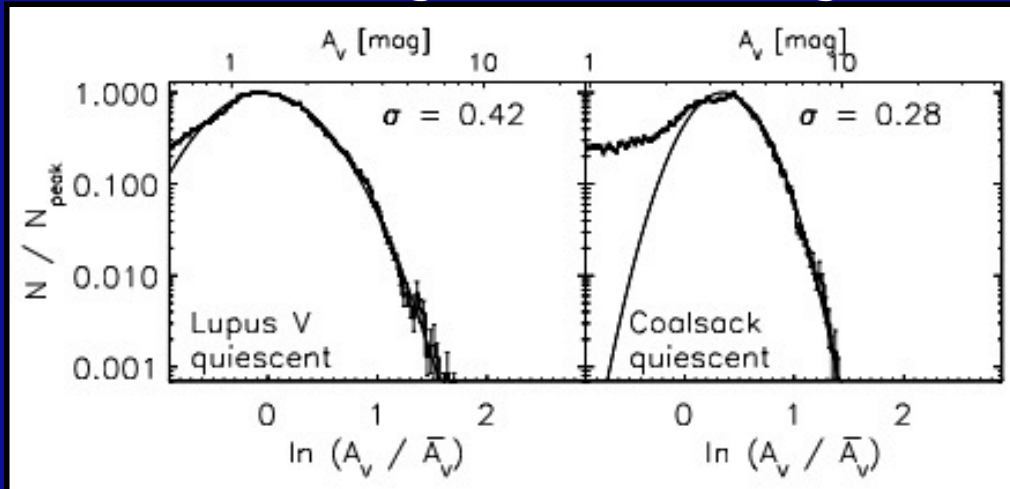


- Clouds fractal and self-similar from 100pc to 0.1pc
- Independent of star-forming or non-star-forming clouds
- Fractal dimension of perimeter P and area A : $P \sim A^{D/2} \rightarrow D \sim 1.4$



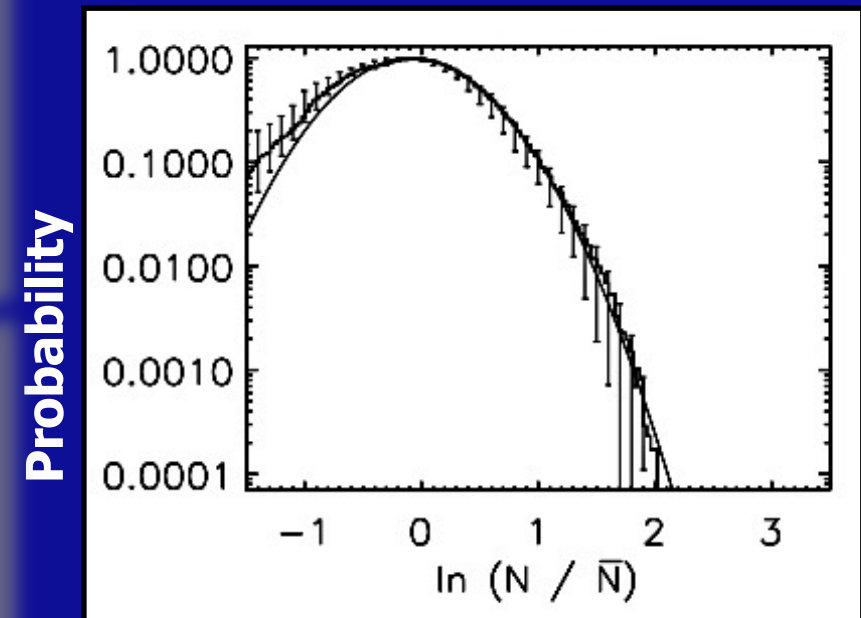
Probability density distributions (PDFs)

Non-star-forming vs. star-forming clouds



Observed distributions

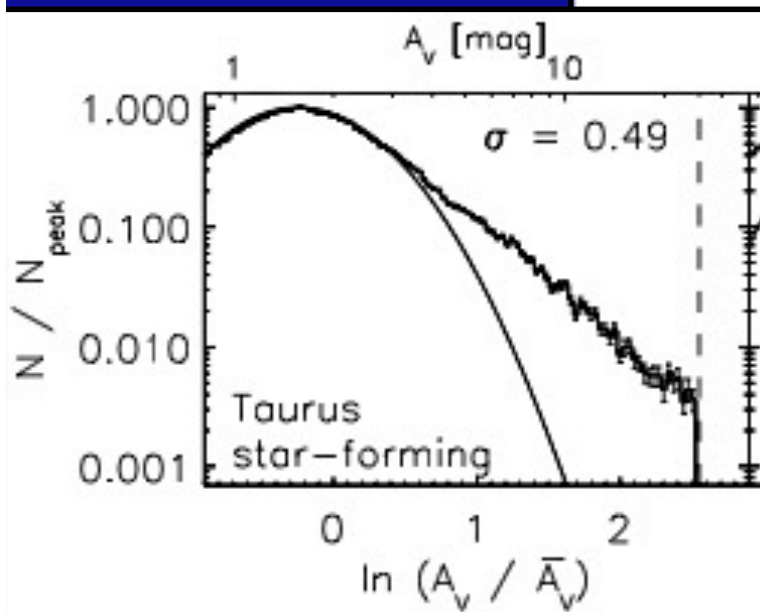
Typical prediction for turbulent media



- non-star-forming: log-normal
- star-forming: high column density excess

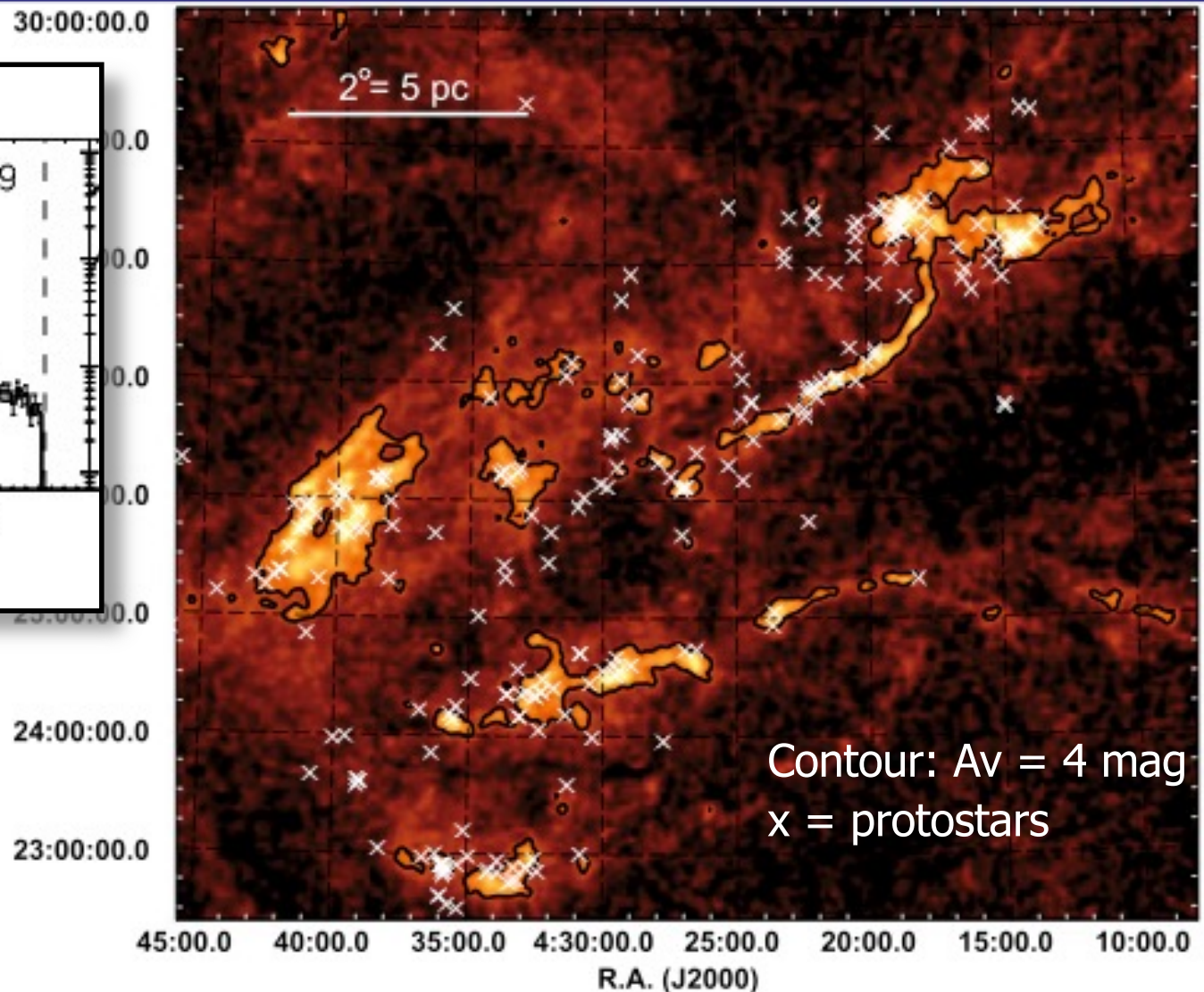
Correlation of "tail" with star formation

Dust column density in Taurus, logarithmic color scale



Excess tail can contain up to 50% of the mass

Kainulainen et al. 2009

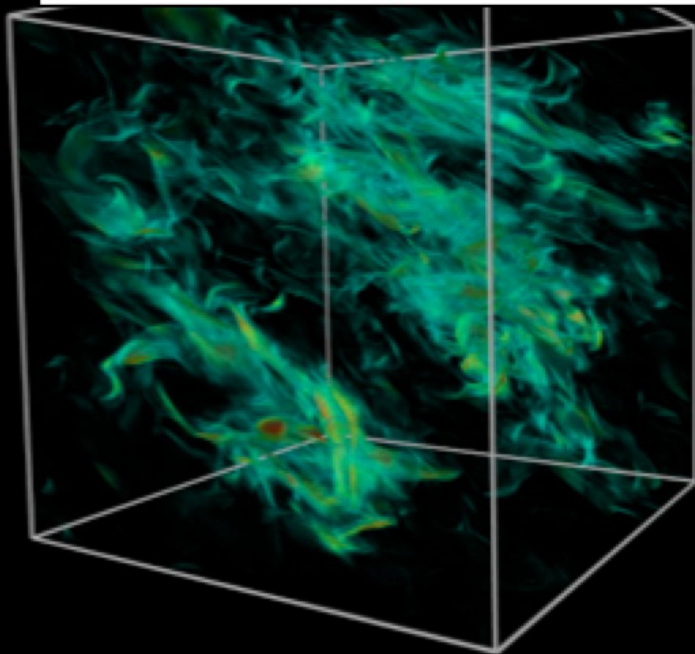


Contour: $A_V = 4 \text{ mag}$
x = protostars

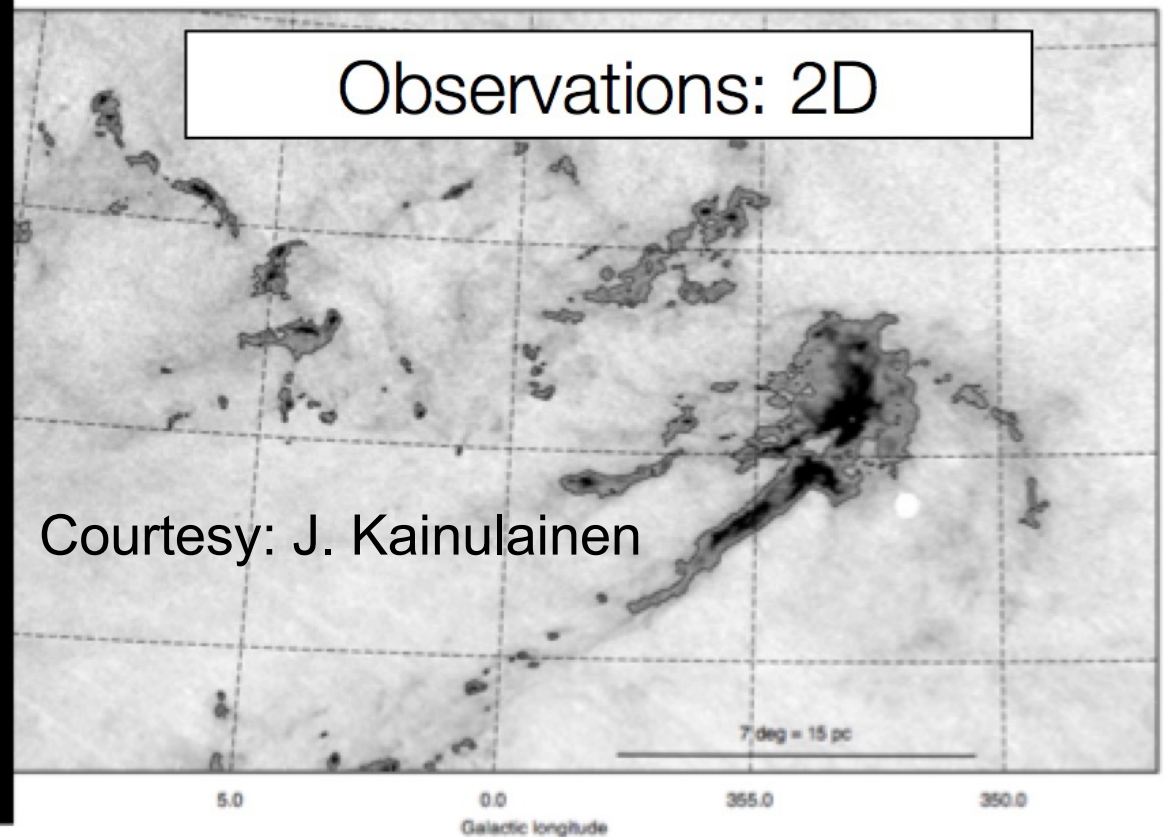
The 3D structure of molecular clouds?

- Observations probe column densities, but theories deal with volume densities.
- How to estimate the 3-dimensional structure?

Theory: 3D

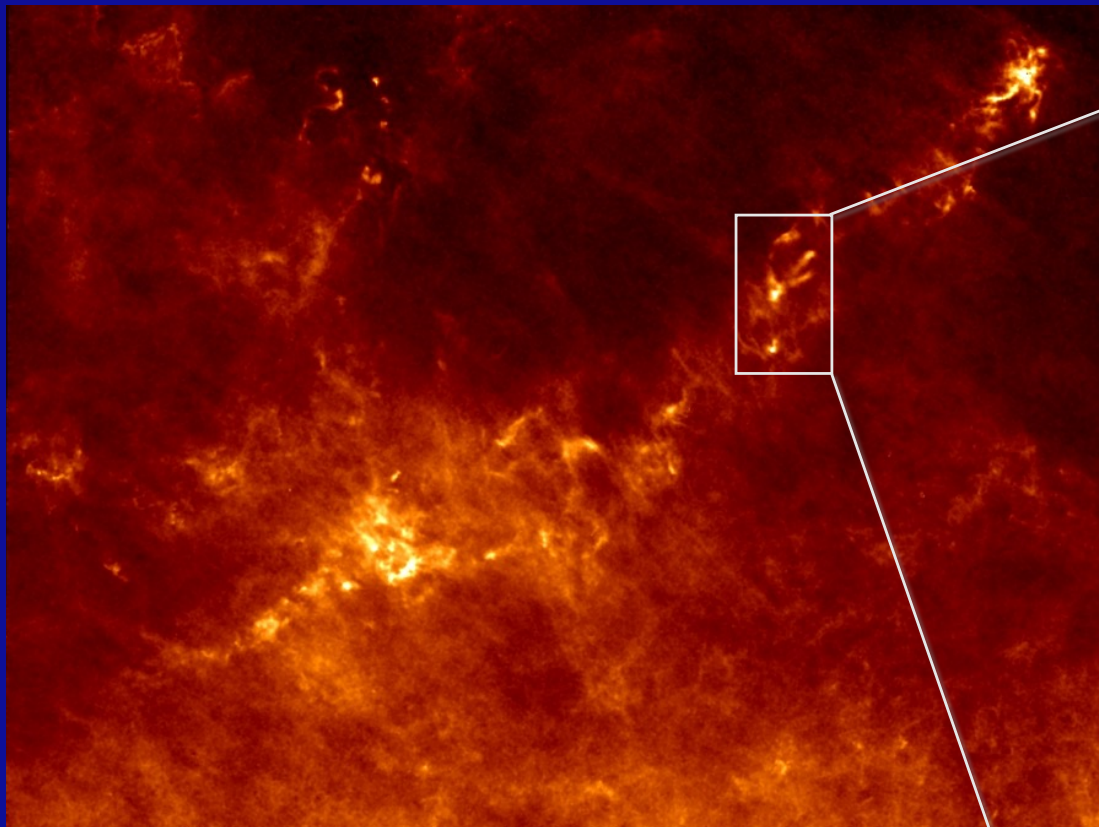


Observations: 2D



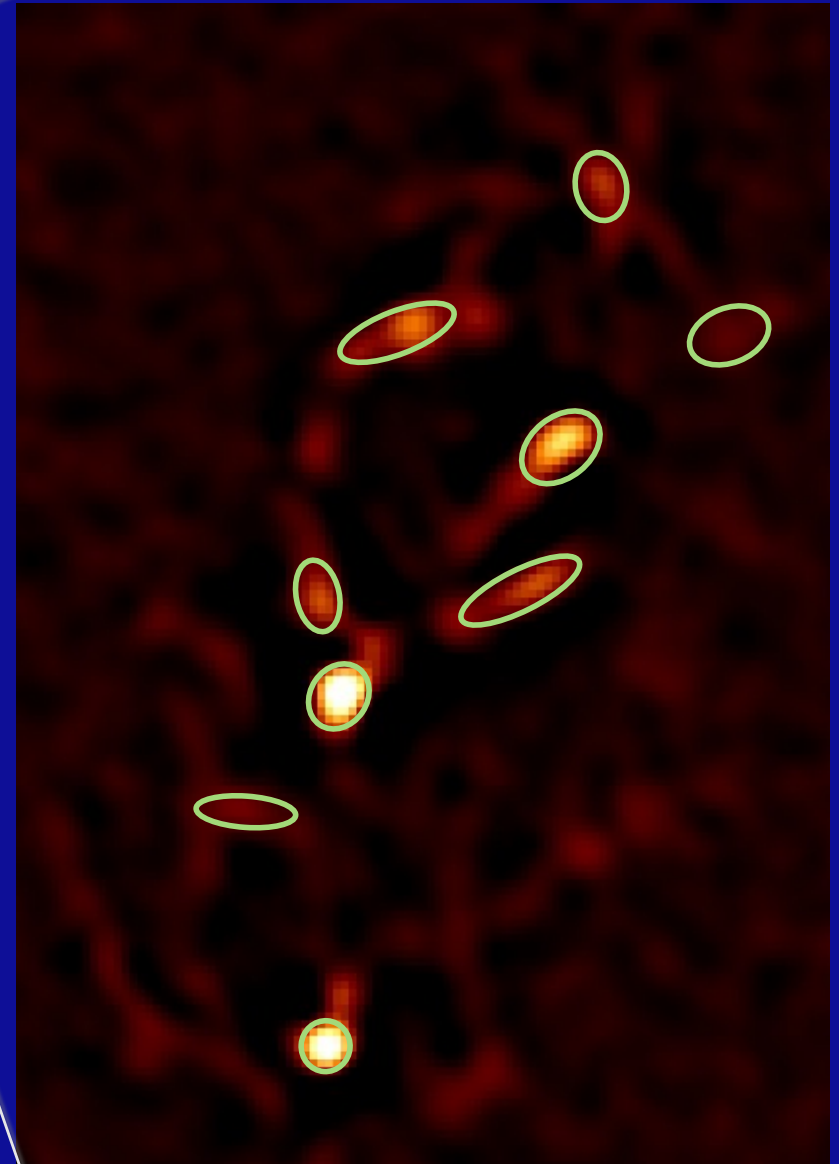
Courtesy: J. Kainulainen

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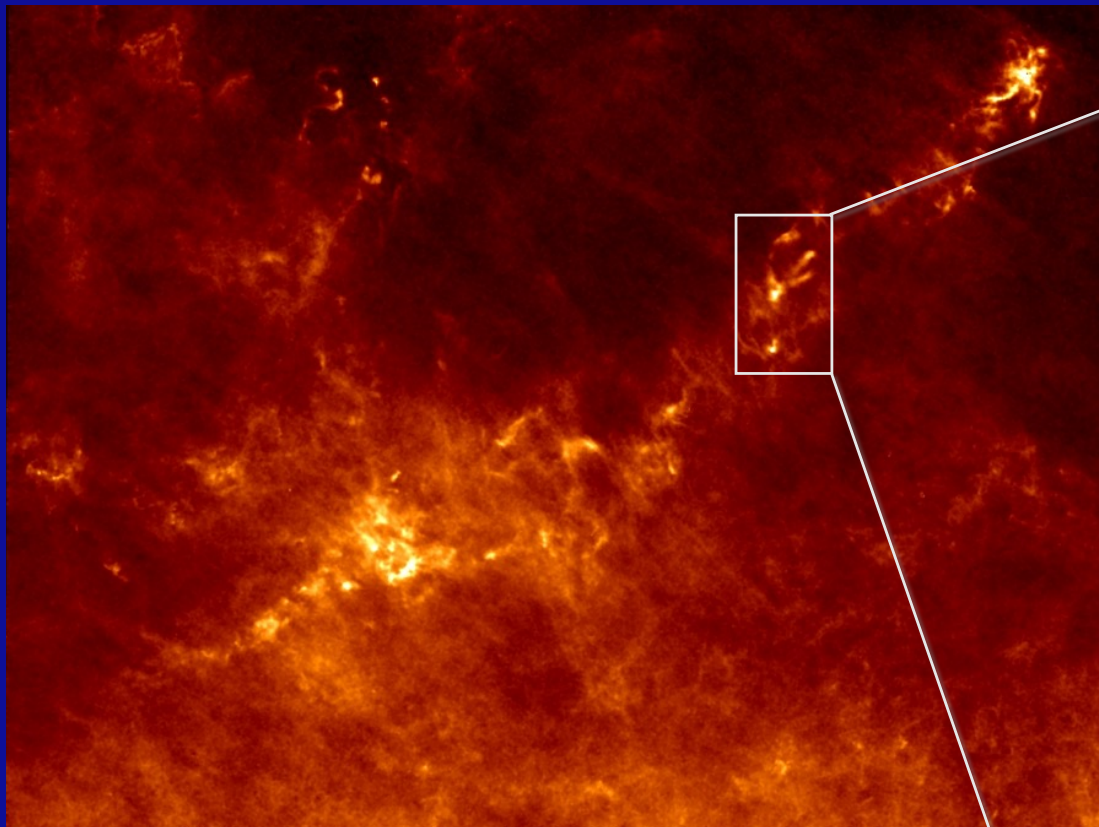


Column density map

Scale decomposition
and object recognition

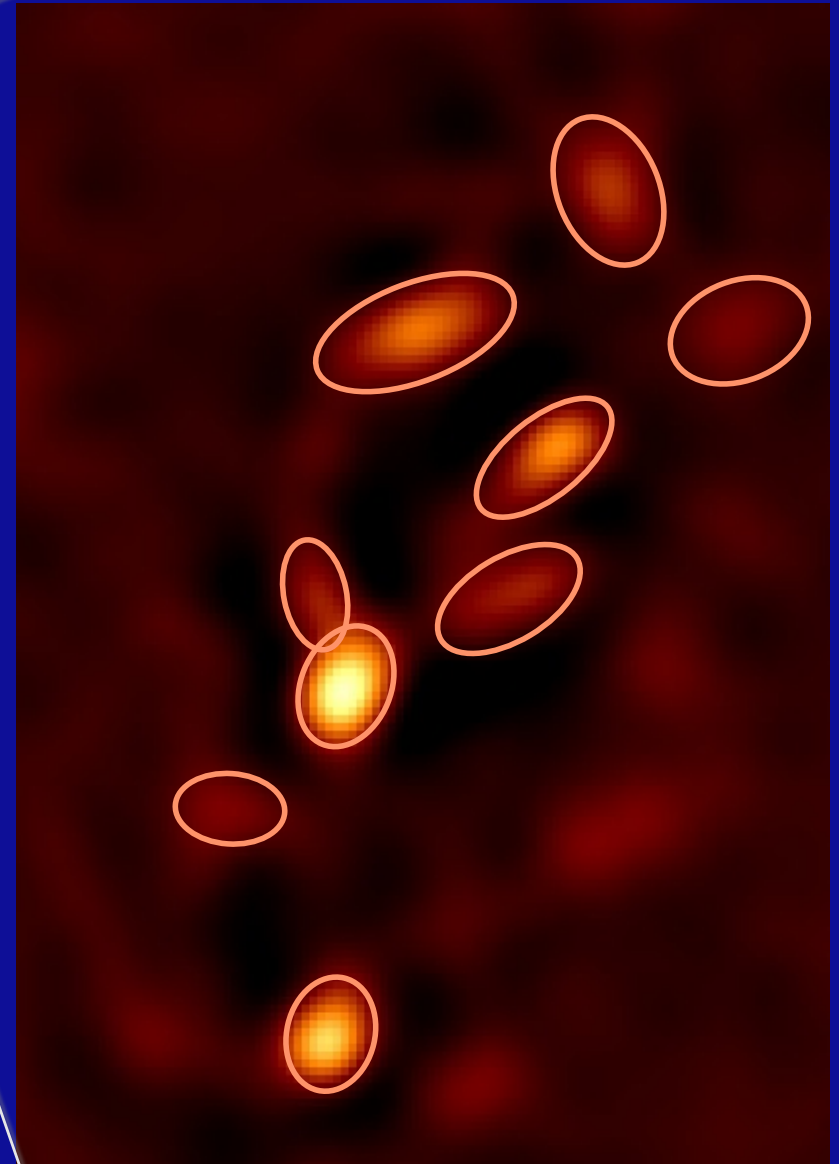


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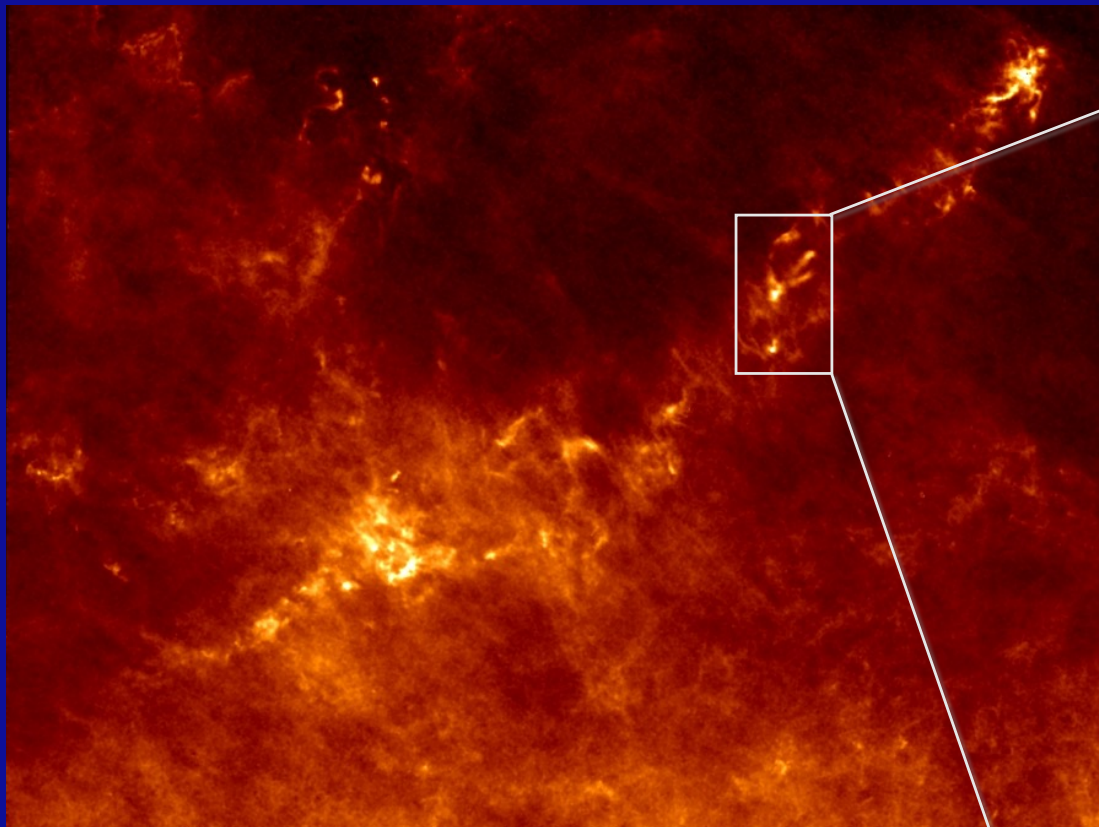


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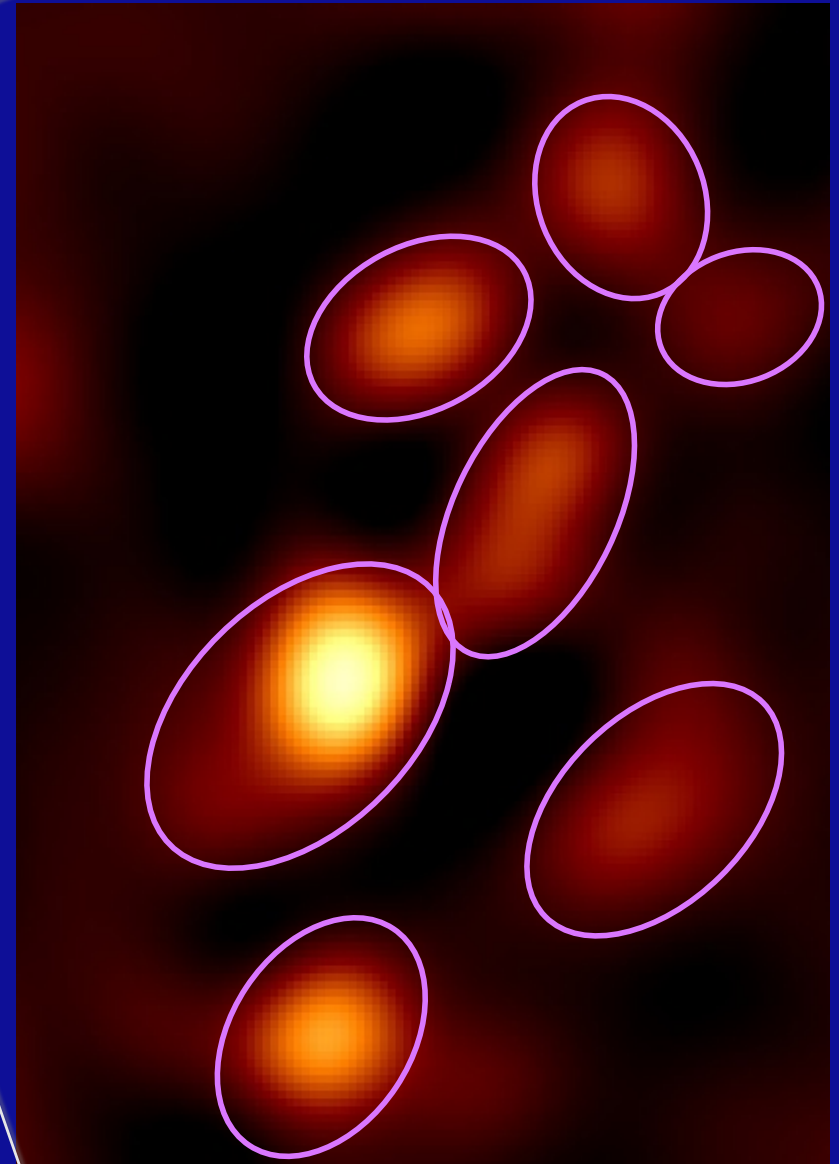


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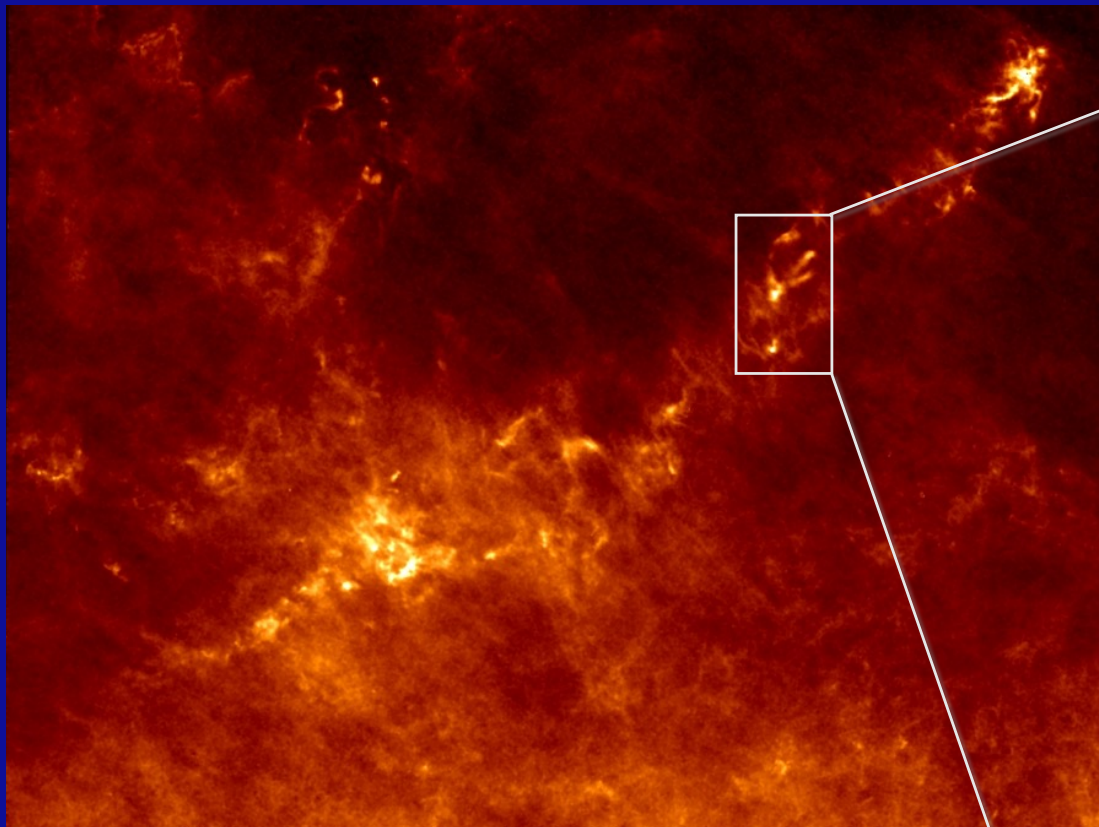


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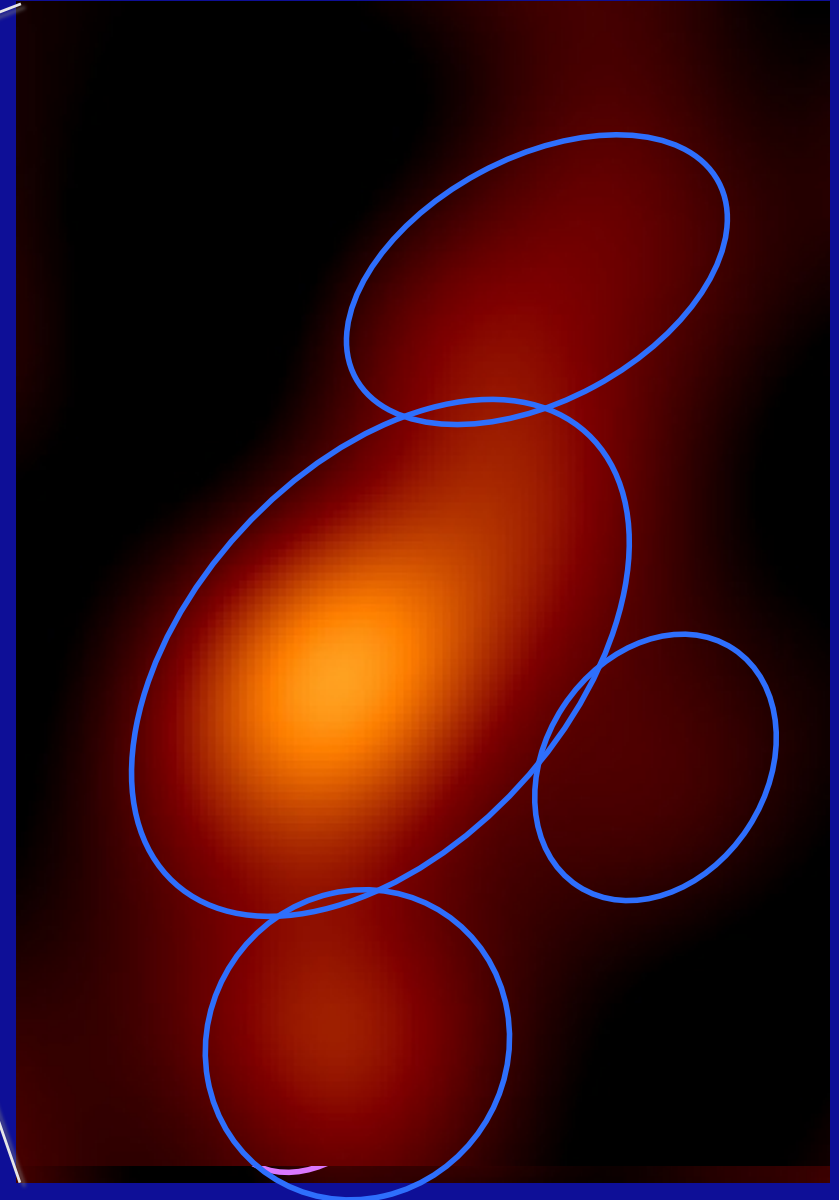


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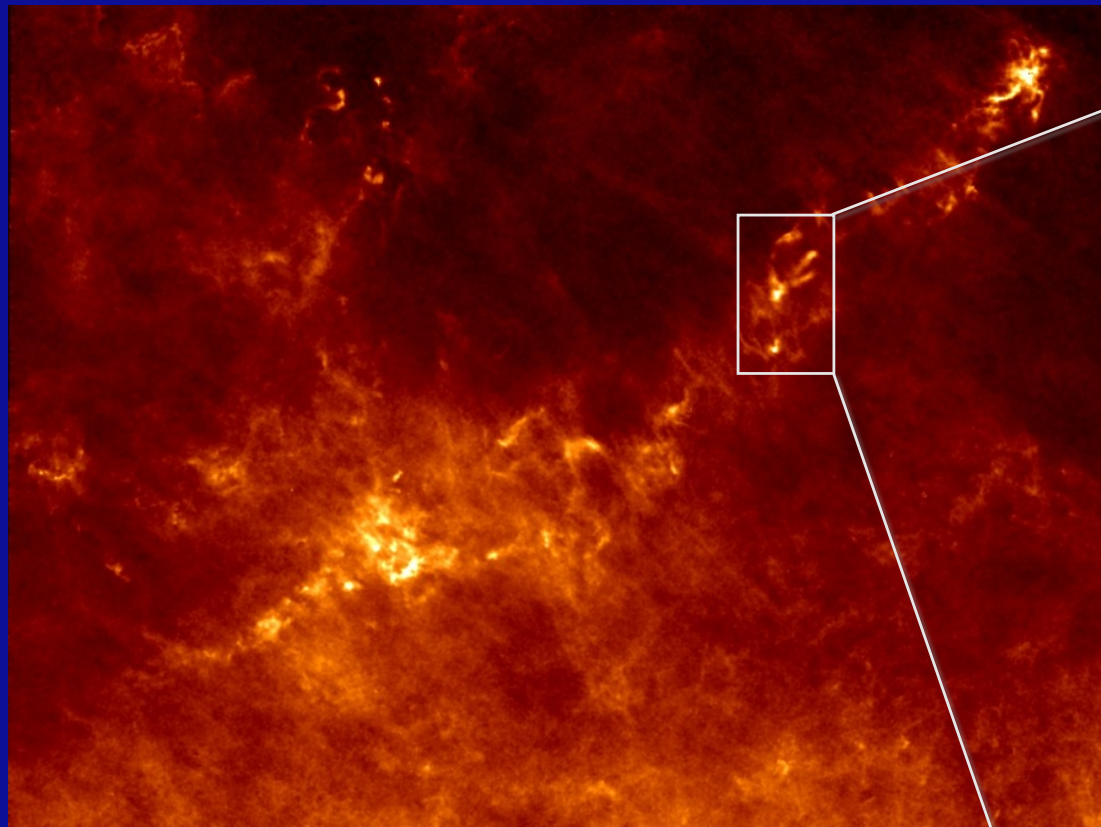


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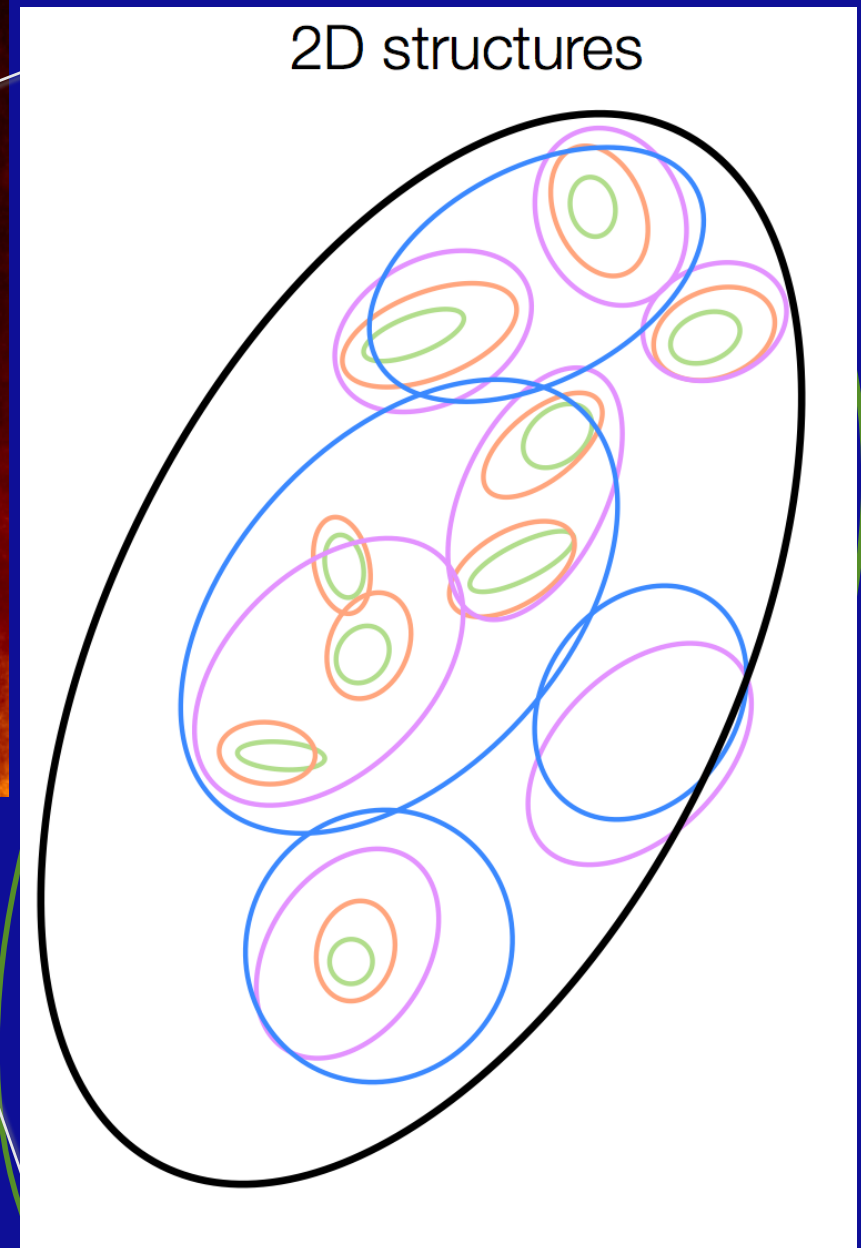


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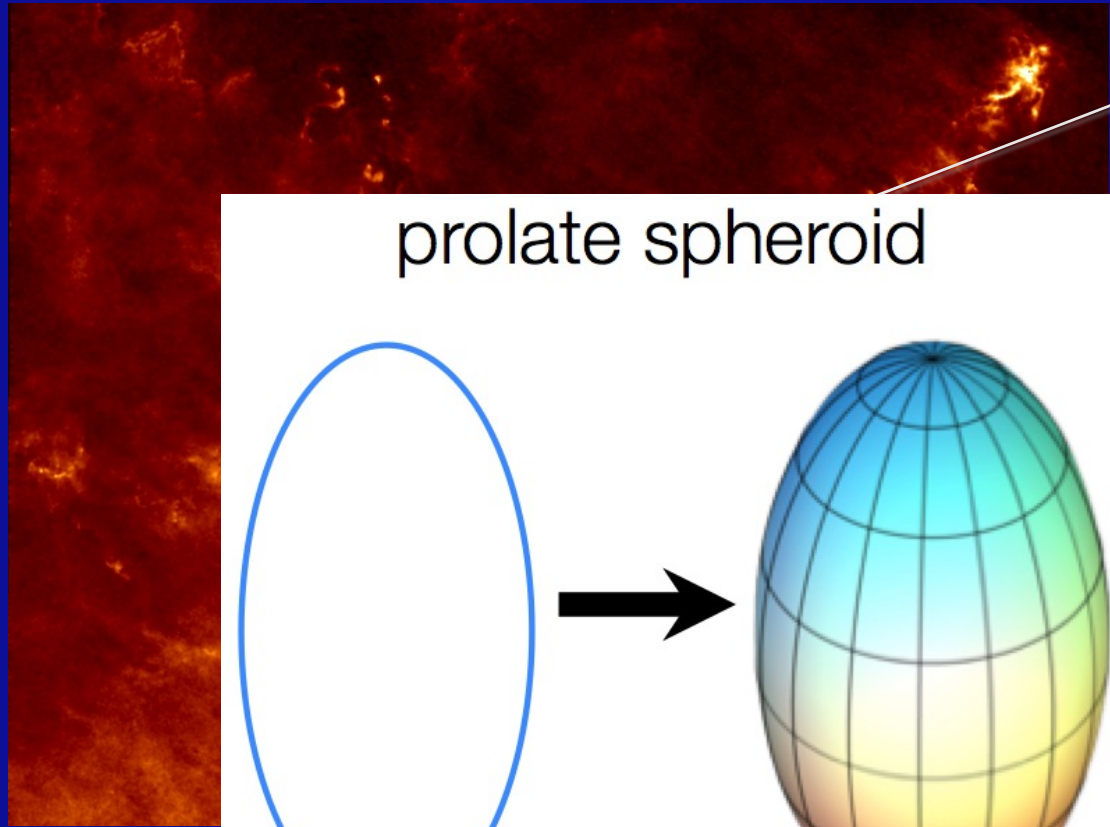


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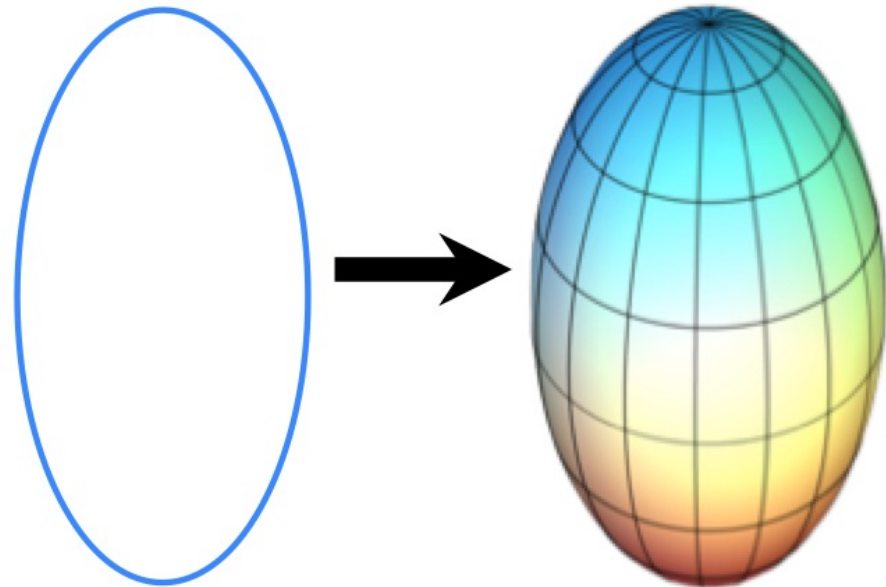


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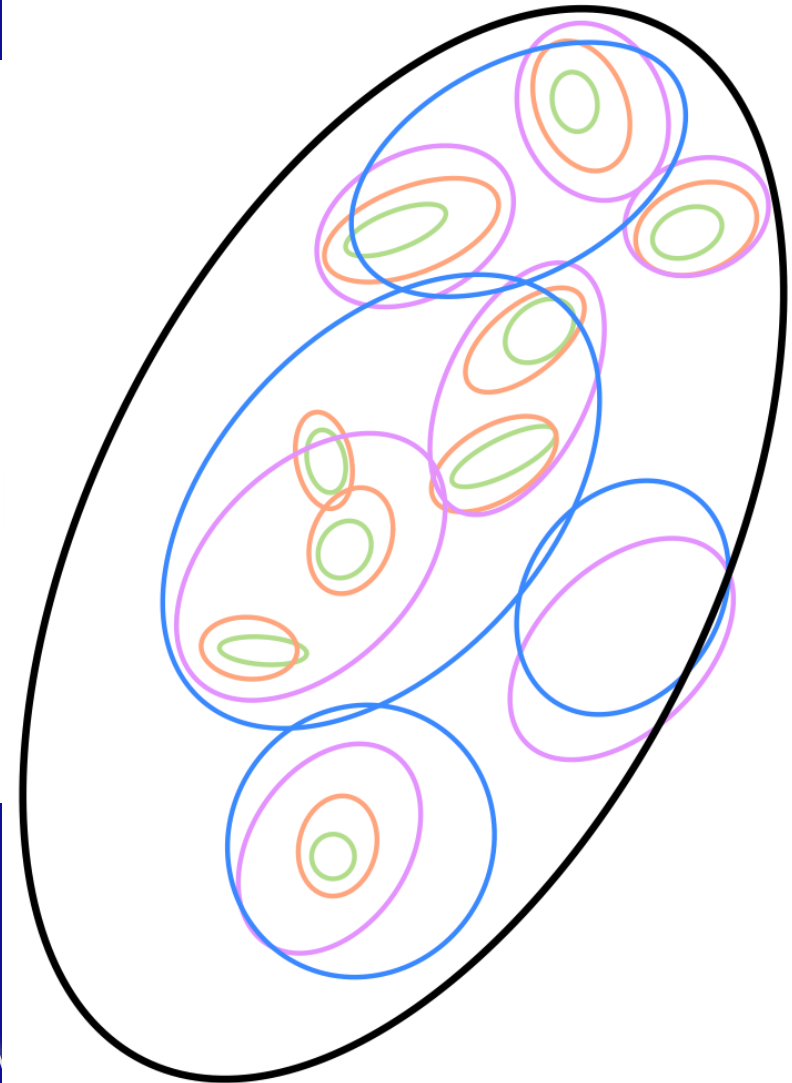


Column

prolate spheroid



2D structures



Scale decomposition
and object recognition

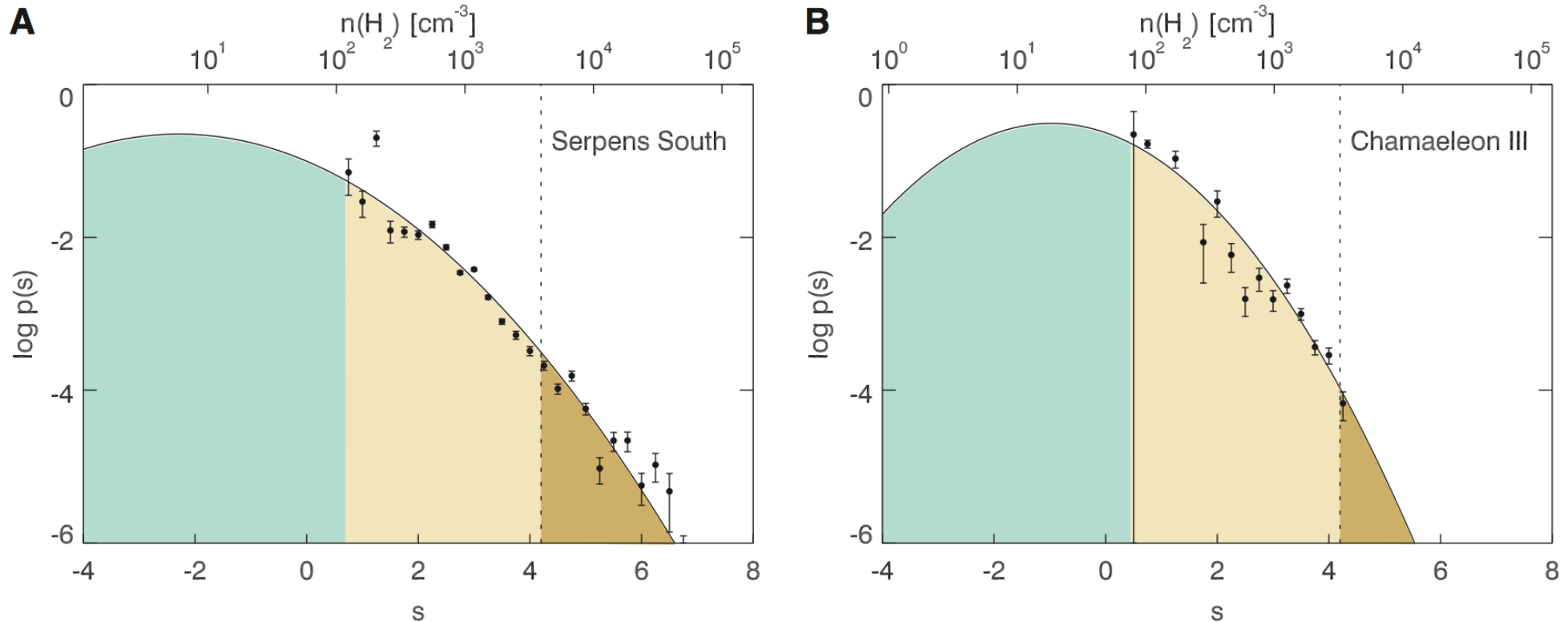
→ Density PDFs

Based on sample of Gould-Belt clouds

Dark brown: star-forming gas

light brown: structures enveloping star-forming gas

Green: non-structured gas



- Direct comparison with theory
- Star formation density threshold $\rightarrow 5 \times 10^3 \text{cm}^{-3}$

Topics today

- Physical distributions (cont.)
- Components of the interstellar medium
- General characteristics of molecular clouds
- **Important cloud relations**
- Cloud fragmentation

A GMC in viral equilibrium

Shortest version of virial theorem (next week): $2T = -W$

(T kinetic energy, W gravitational energy)

$$2T = 2 * (1/2 m \Delta v^2) = -W = Gm^2/r$$

$$\rightarrow \text{virial velocity: } v_{\text{vir}} = (Gm/r)^{1/2}$$

$$\rightarrow \text{or virial mass: } m_{\text{vir}} = v^2 r / G$$

Luminosity-mass relation

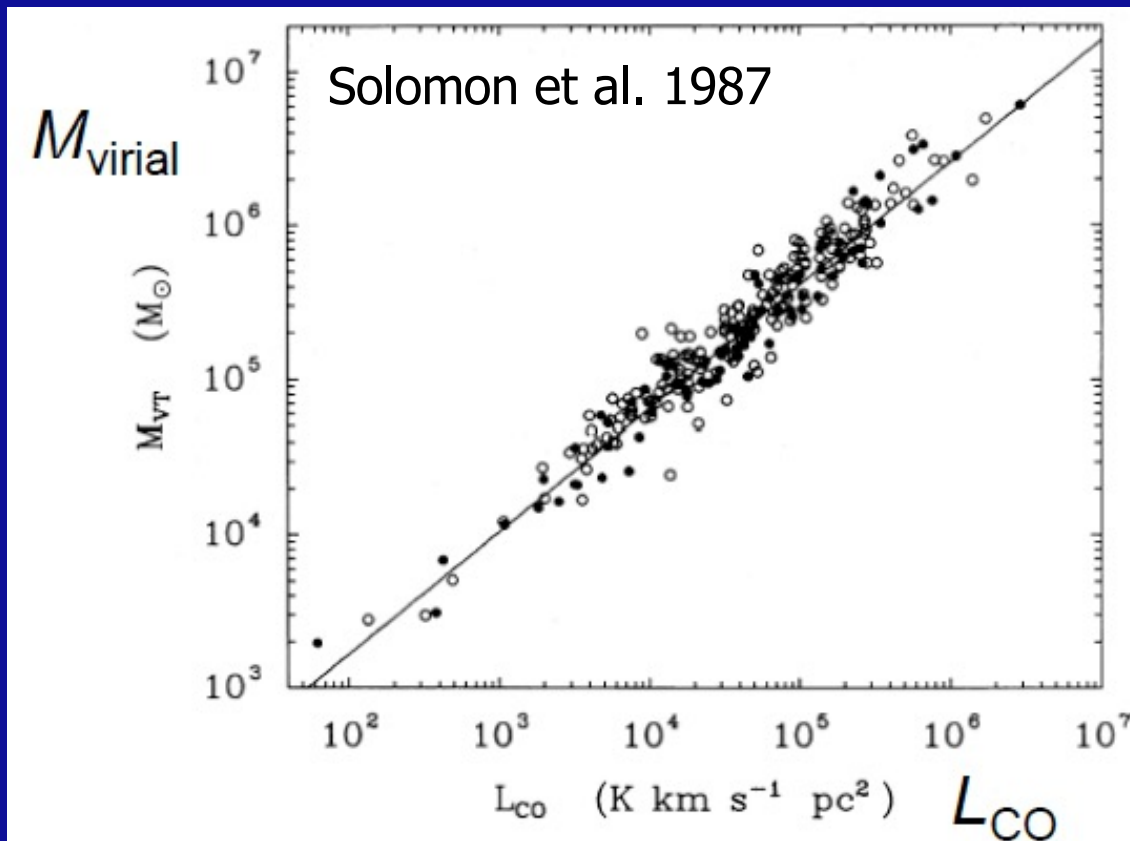
Integrated CO intensity: $I_{\text{CO}} = \int T(v)dv$

CO luminosity $L_{\text{CO}} = T\Delta v \pi r^2$

(T brightness temperature, Δv linewidth, r cloud radius)

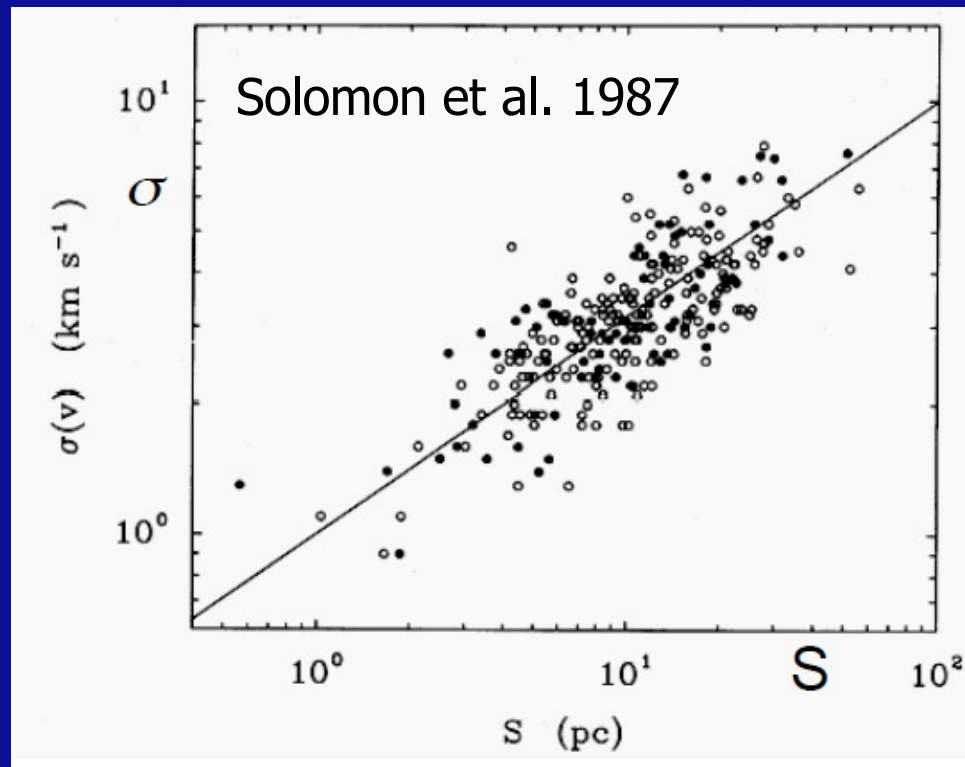
Substituting $v = (Gm/r)^{1/2}$ and mass $m = 4/3\pi r^3 \rho$

$$\rightarrow L_{\text{CO}} = (3\pi G/(4\rho))^{1/2} T m$$



GMCs are roughly in virial equilibrium.

The linewidth-size relation



- Linewidth-size relation first found by Larson 1981 (Thermal CO linewidth at 20K only ≈ 0.1 km/s)
- Approximate relation: linewidth $\approx \sqrt{\text{size}}$
- Extends over many orders of magnitude in size but not down to cores
- Implies strong turbulent contribution to the ISM

Additional relations

- Linewidth-size relation: $dv \approx r^{1/2}$
- Virial equilibrium: $dv \approx (Gm/r)^{1/2} \rightarrow m = dv^2r/G$

This leads to other relations:

$\rightarrow m = r^2/G \rightarrow m/r^2 = \text{constant} \rightarrow$ approximate constant column density N in GMCs

$\rightarrow \rho \approx m/r^3 \approx dv^2/G * 1/r^2$

Some average empirical values for GMCs:

$$N \approx 1.5 \times 10^{22} \text{ cm}^{-2}$$

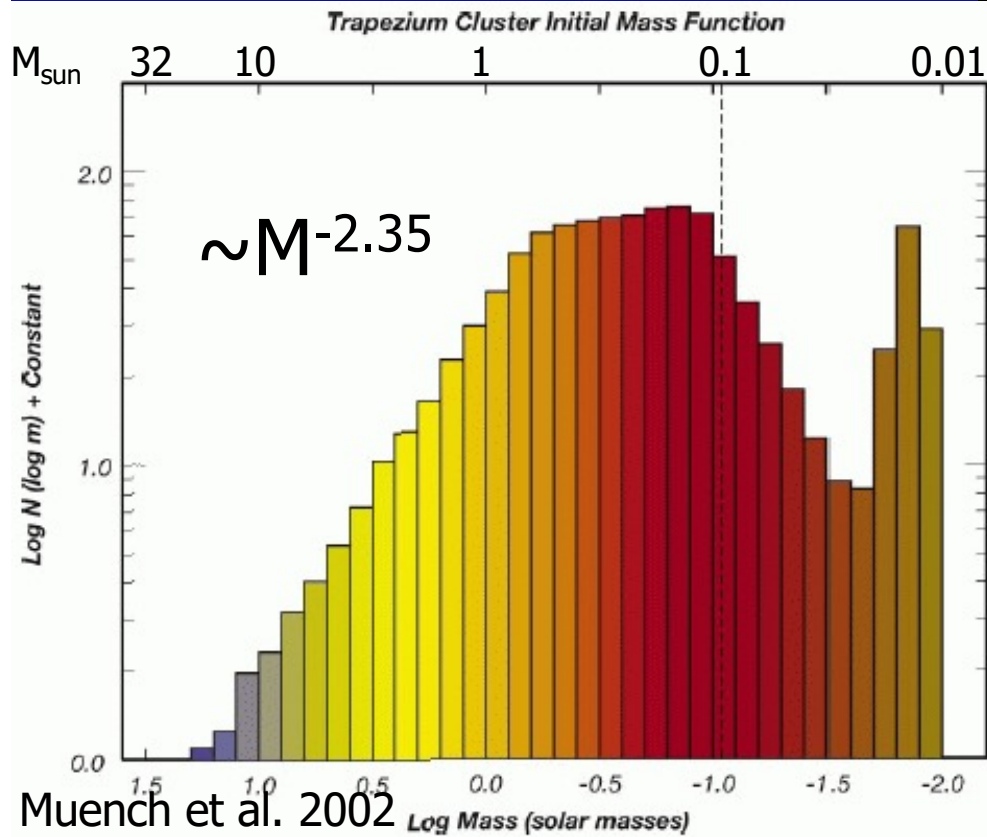
$$A_V \approx 10 \text{ mag}$$

$$\Sigma \approx 150 M_{\text{sun}} \text{ pc}^{-2}$$

Topics today

- Physical distributions (cont.)
- Components of the interstellar medium
- General characteristics of molecular clouds
- Important cloud relations
- **Cloud fragmentation**

Clusters and the Initial Mass Function (IMF)



Orion Nebula

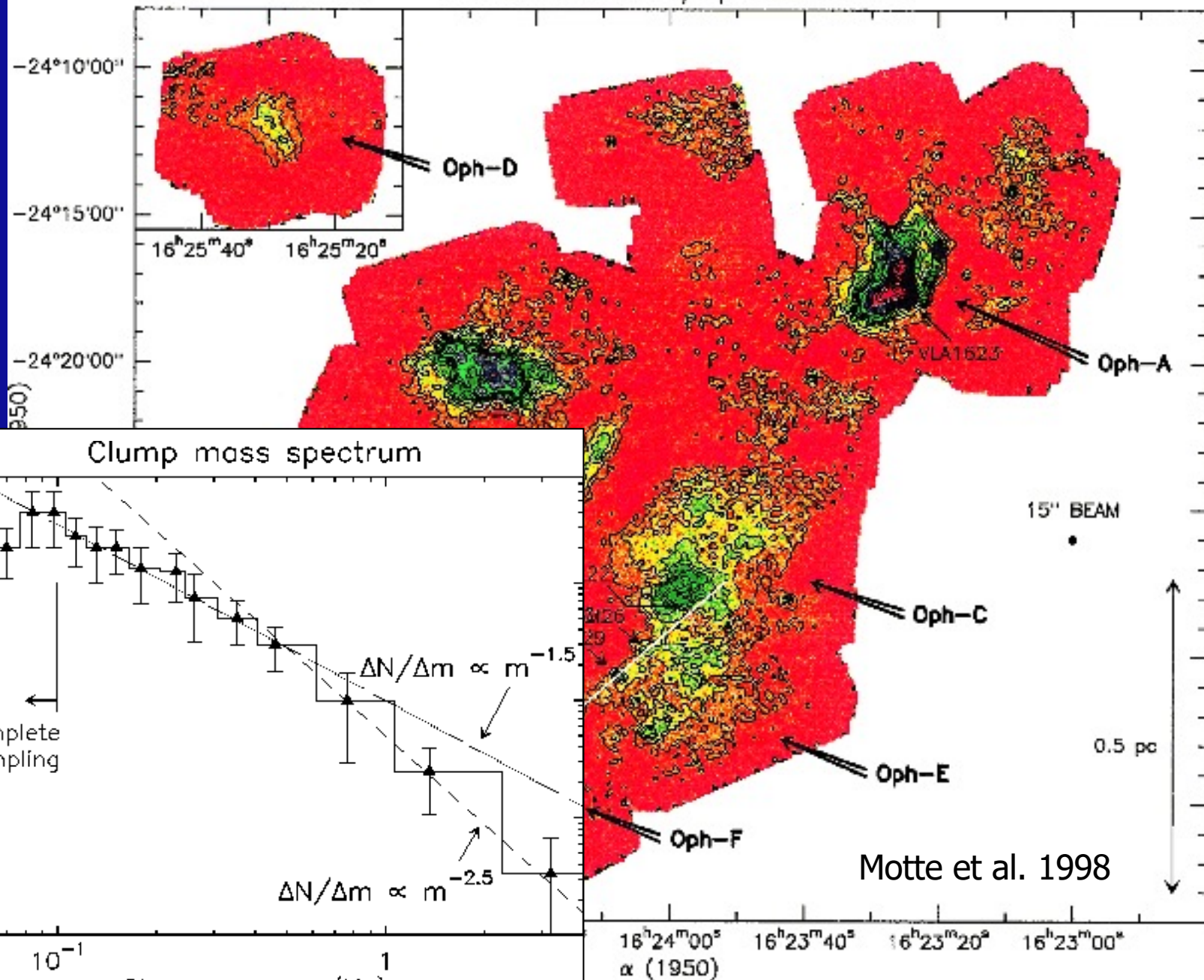
Subaru Telescope, National Astronomical Observatory of Japan

CISCO (J, K' & H₂ ($v=1-0$ S(1)))

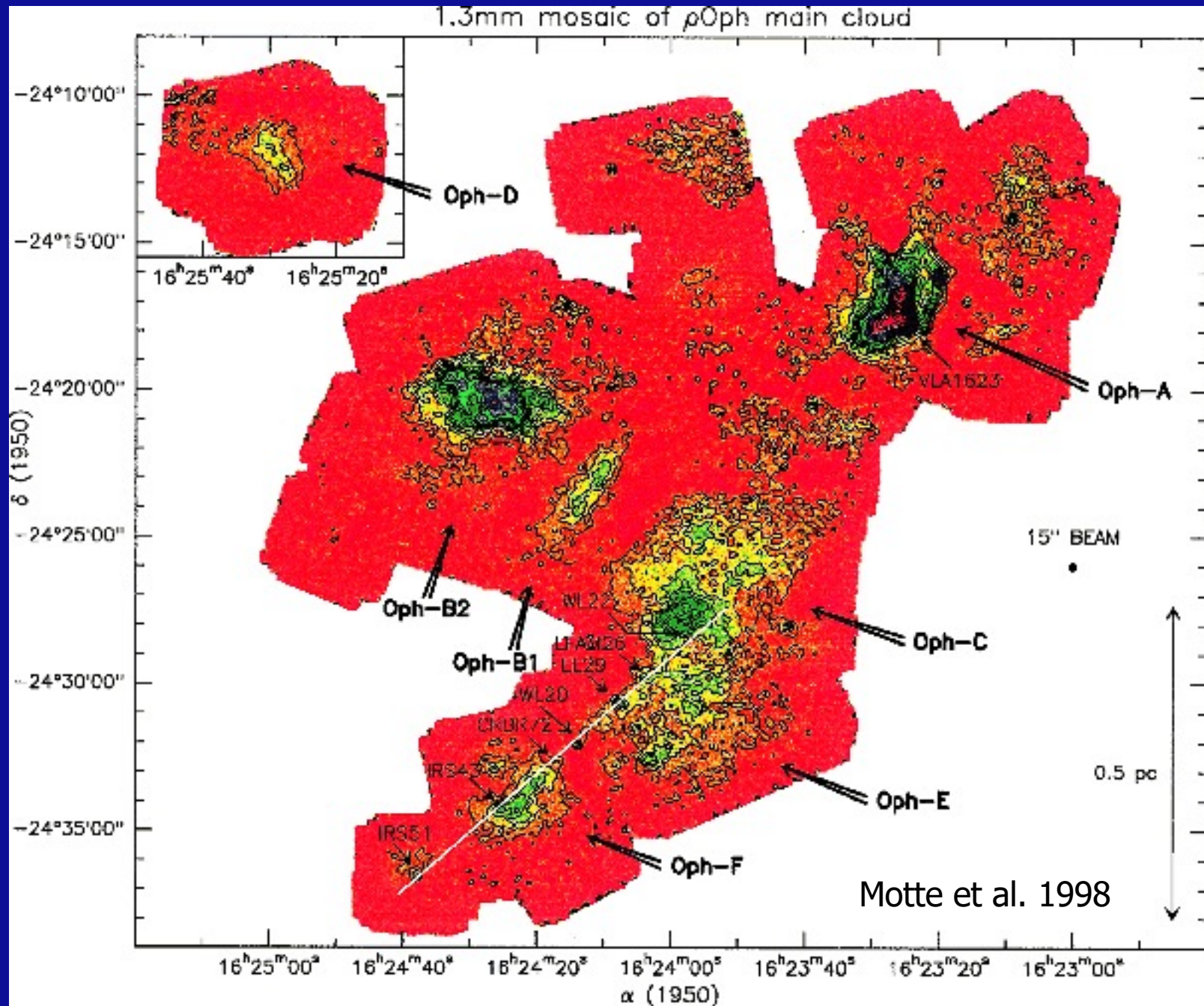
January 28, 1999

Pre-stellar core mass functions

1.3mm mosaic of ρ Oph main cloud



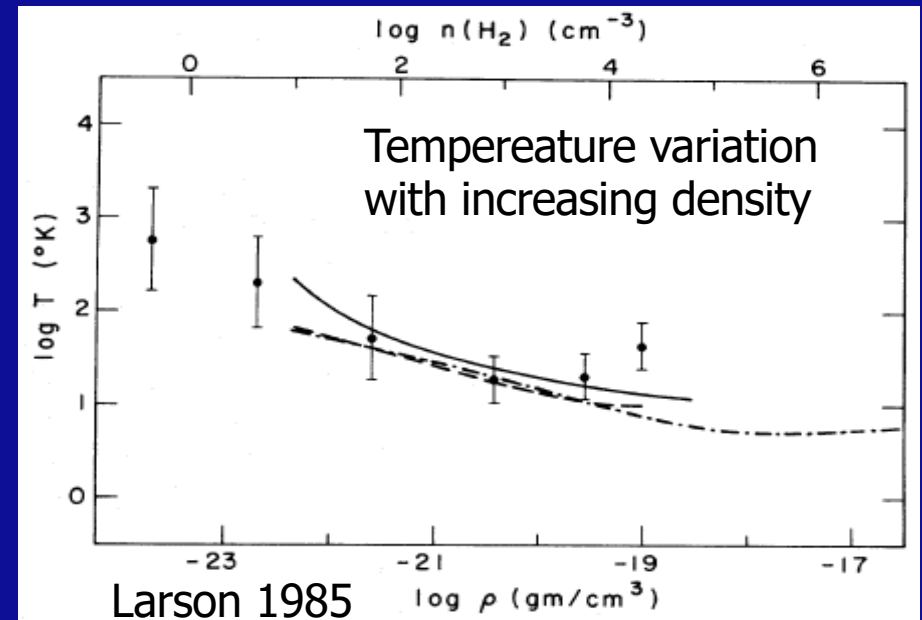
Pre-stellar core mass functions



Characteristic mass defined by thermal physics

- Jeans mass depends on T:

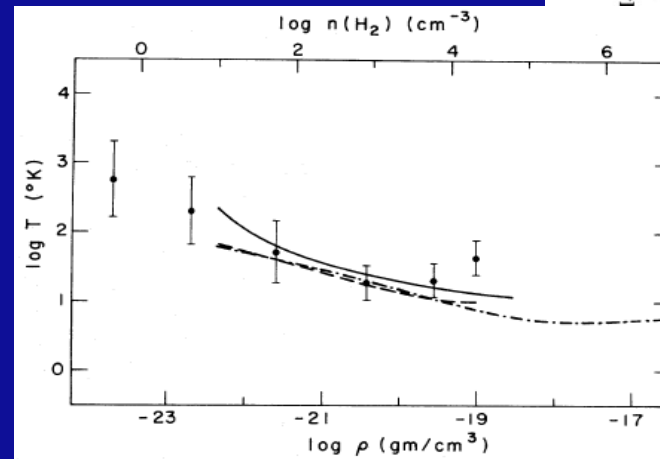
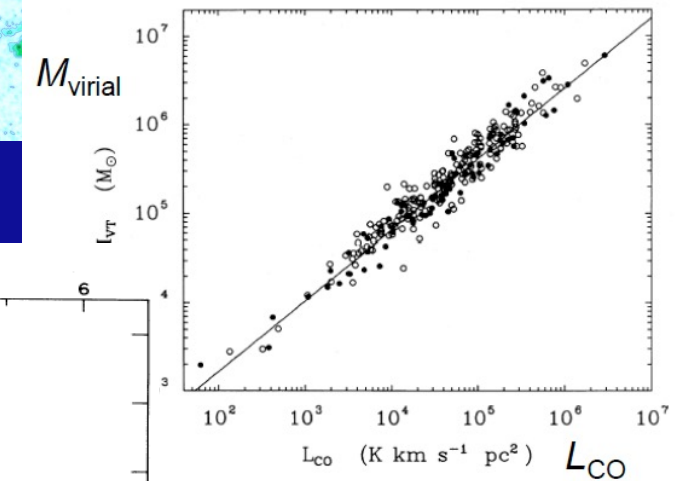
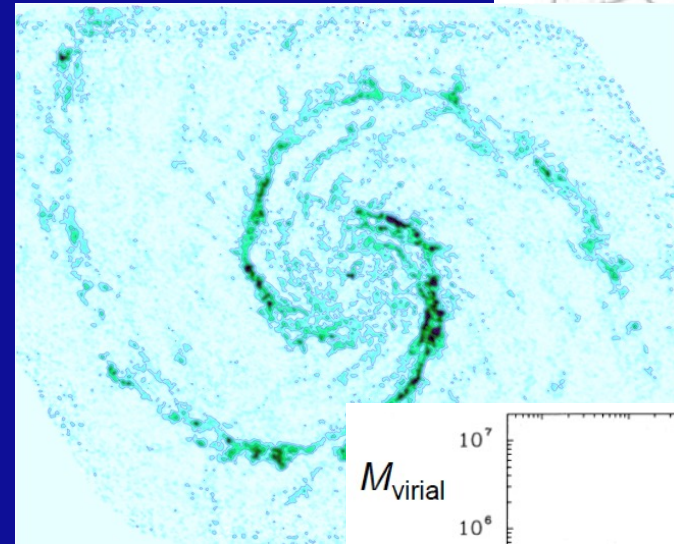
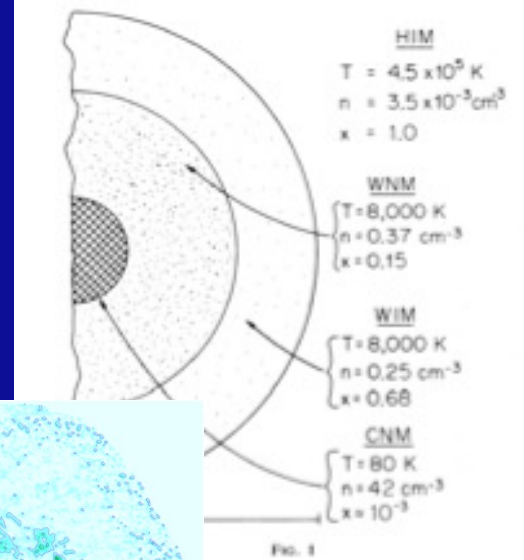
$$M_J \sim a_t^3 / \rho_0^{1/2} \sim T^{3/2} / \rho^{1/2}$$



- Low densities \rightarrow T decreases with increasing $\rho \rightarrow$ regions cool efficiently
 \rightarrow decreasing M_J suggests that fragmentation is favoured
- Further increasing $\rho \rightarrow$ gas thermally couples to dust and clouds, and become partially optically thick \rightarrow Cannot cool well enough anymore
 \rightarrow temperature slightly increases again.
 \rightarrow M_J decreases slower, inhibiting much further fragmentation.
- \rightarrow Regime with lowest T should correspond to preferred fragmentation scale
- \rightarrow The Jeans mass at this point is about $0.5 M_{\text{sun}}$.

Summary

- Physical distribution
- Different components of ISM
- Basic characteristics
- Important cloud relations
- Cloud fragmentation



Sternentstehung - Star Formation

Winter term 2024/2025

Henrik Beuther, Thomas Henning & Caroline Gieser

- 15.10 *Today: Introduction & Overview* (Beuther)
22.10 *Physical processes I* (Beuther)
29.10 --
05.11 *Physical processes II* (Beuther)
12.11 *Molecular clouds as birth places of stars* (Beuther)
19.11 Molecular clouds (cont.), Jeans Analysis (Henning)
26.11 Collapse models I (Beuther)
03.12 Collapse models II (Beuther)
10.12 Protostellar evolution (Gieser)
17.12 Pre-main sequence evolution & outflows/jets (Henning)
07.01 Accretion disks I (Henning)
14.01 Accretion disks II (Henning)
21.01 High-mass star formation, clusters and the IMF (Gieser)
28.01 Extragalactic star formation (Henning)
04.02 Planetarium@HdA, outlook, questions
11.02 Examination week, no star formation lecture (Beuther, Gieser, Henning)

Book: Stahler & Palla: The Formation of Stars, Wileys

More Information and the current lecture files: http://www.mpia.de/homes/beuther/lecture_ws2425.html
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Heidelberg Joint Astronomical Colloquium Winter Semester 2024/25

Tuesday November 12th Main Lecture Theatre, Philosophenweg 12, 16:30 CEST

Elena Pancino

(INAF – Osservatorio Astrofisico di Arcetri):

Stardance: The non-canonical evolution of stars in clusters

<https://www.physik.uni-heidelberg.de/hephysto/>

Host: Michela Mapelli (mapelli@uni-heidelberg.de)

Image Credit: Mark A. Garlick / University of Warwick